

1 **Semi-inclusive charged jet production dependence on**  
2 **event activity at high backward-rapidity in**  
3  **$\sqrt{s_{NN}} = 200$  GeV  $p$ +Au collisions at STAR**

---

4 **David Stewart (for the STAR Collaboration)<sup>a,\*</sup>**

5 <sup>a</sup>*Yale University,*

6 *Wright Laboratory, New Haven, CT 06520*

7 *E-mail: [david.j.stewart@yale.edu](mailto:david.j.stewart@yale.edu)*

8 These proceedings present measurements of mid-rapidity semi-inclusive charged jet spectra and acoplanarity in relation to event activity (EA) at backward-rapidity (Au-going) in  $\sqrt{s_{NN}} = 200$  GeV  $p$ +Au collisions at STAR. Small system jet measurements have been motivated by apparent flow signals in these systems, suggestive of QGP formation. These specific small system jet measurements are particularly timely because they are the first semi-inclusive jet spectrum measurements in small systems at RHIC energy. As such, they complement results from PHENIX, ATLAS, and ALICE, which collectively are not yet well understood. The new STAR semi-inclusive jet spectra are dependent on EA; they are distinctly anti-correlated. Evidence is presented, using both the new measurements and simulation, that this anti-correlation results from constraints in the phase space of the initial collisions.

*HardProbes2020*

*1-6 June 2020*

*Austin, Texas*

---

\*Speaker

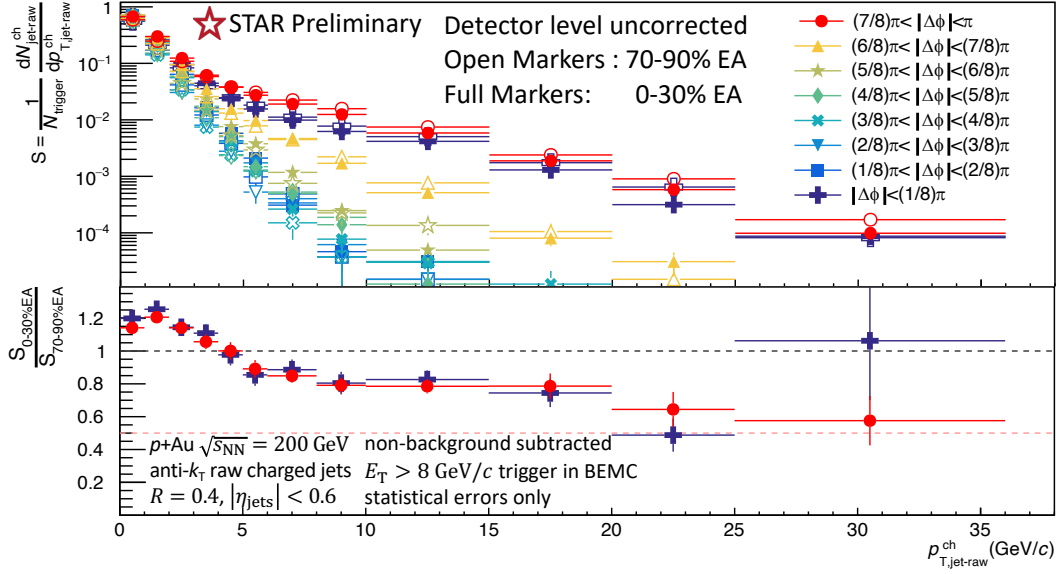
9 The observation of flow-like signals in  $pp$  collisions about a decade ago [1] has precipitated  
 10 a resurgence of interest in the heavy-ion community in studying small systems ( $pp$ ,  $d+A$ ,  $\text{He}^3+A$ ,  
 11 or  $p+A$  collisions) for results attributed to final-state, quark gluon plasma (QGP), effects in A+A  
 12 collisions (such as collectivity). Due to their early formation times, jets have proven to be valuable  
 13 probes of hot nuclear matter effects in A+A collisions, and a natural venue of investigation for  
 14 possible analogous hot nuclear matter effects in small systems.

15 Several methods of normalization are used for jet modification measurements. Inclusive jet  
 16 production per binary nucleon-nucleon collision ( $N_{\text{coll}}$ ) can be compared to  $pp$  results ( $R_{AA}$ ,  $R_{pA}$ ,  
 17  $R_{dA}$ ,  $R_{\text{He}^3A}$ ); they may be compared semi-inclusively as production per-trigger; or the shape  
 18 of their per-jet distributions may be compared, such as dijet momentum balance distributions.  
 19 Primary jet quenching observables include transverse momentum ( $p_T$ ) spectra, acoplanarity (jet  
 20 azimuthal distribution relative to a trigger), momentum imbalances with a recoil partner, and  
 21 possible substructure modifications.

22 When not binned by event activity (EA), all reported inclusive measurements of  $R_{pA}$  and  
 23  $R_{dA}$  ( $R_{p/dA}$ ) have been consistent with scaled  $pp$  results [2–5]. However, when binned by high-  
 24 rapidity EA, both ATLAS and PHENIX reported enhancement (suppression) of  $R_{p/dA}$  for peripheral  
 25 (central) events. Intriguingly, ATLAS showed that the  $p$ -going jet spectra modification appeared  
 26 to scale with Bjorken- $x$ . When measuring recoil jet spectra per trigger (a high- $p_T$  hadron) at a  
 27 lower Bjorken- $x$ , ALICE found no modification [6], setting an out-of-jet cone energy transport  
 28 limit. The jet modification in the ATLAS measurement, if resulting only from out-of-cone energy  
 29 transport, violate the limit set at ALICE. The question still under study is if the modifications  
 30 observed by ATLAS and PHENIX are indeed dependent on Bjorken- $x$  and hence not seen at the  
 31 kinematics probed by ALICE, or if they result from misclassifying centrality (and therefore  $N_{\text{coll}}$ )  
 32 due to phase space constraints and/or other mechanisms correlating the multiple scatterings of the  
 33 proton [8]. In parallel efforts, measurements of dijet momentum balance, acoplanarity, and jet mass  
 34 distributions (shape measurements normalization per dijet or per jet) in  $p+Pb$  at the LHC have  
 35 found no modification relative to  $pp$  collisions [9–11].

36 The STAR measurements reported here are for mid-rapidity charged jets whose spectra is  
 37 normalized by number of triggers of high energy hits at mid-rapidity in the Barrel Electromagnetic  
 38 Calorimeter. Event activity (EA) is measured in the pseudorapidity ( $\eta$ ) range of  $[-5, -3.4]$ . These  
 39 jets access similar Bjorken- $x$  ranges as the ATLAS and PHENIX results while avoiding  $N_{\text{coll}}$  scaling  
 40 uncertainties. The jet spectra (for  $p_{T,\text{jet}}^{\text{ch}} > 5 \text{ GeV}/c$ ) per trigger are clearly suppressed in high EA  
 41 events compared to low EA events, as shown in Fig. 1. This is also qualitatively present in PYTHIA 8  
 42 simulations, which has no mechanisms modeling jet modification from hot or cold nuclear matter  
 43 effects.

44 The collision triggers, in data and simulation, are provided by jets (or parts thereof) at mid-  
 45 rapidity (mid- $\eta$ ). These jets are accompanied by dijet partners which are, at leading order, exactly  
 46 opposite in azimuth and contribute to the spectra-per-trigger at  $|\eta| < 0.6$ . A correlation between jet  
 47 spectra and EA may arise generally from three scenarios. First, there could be an actual correlation  
 48 between the cross sections for hard scattering and the particles that contribute to EA. Second, an  
 49 apparent correlation between cross sections could result in, due to methodology, auto-correlations  
 50 among EA, jets, and/or triggers. A primary example of this for the STAR spectra would be a trivial  
 51 autocorrelation from when the recoil dijet partner is not at mid-rapidity but instead fragments within



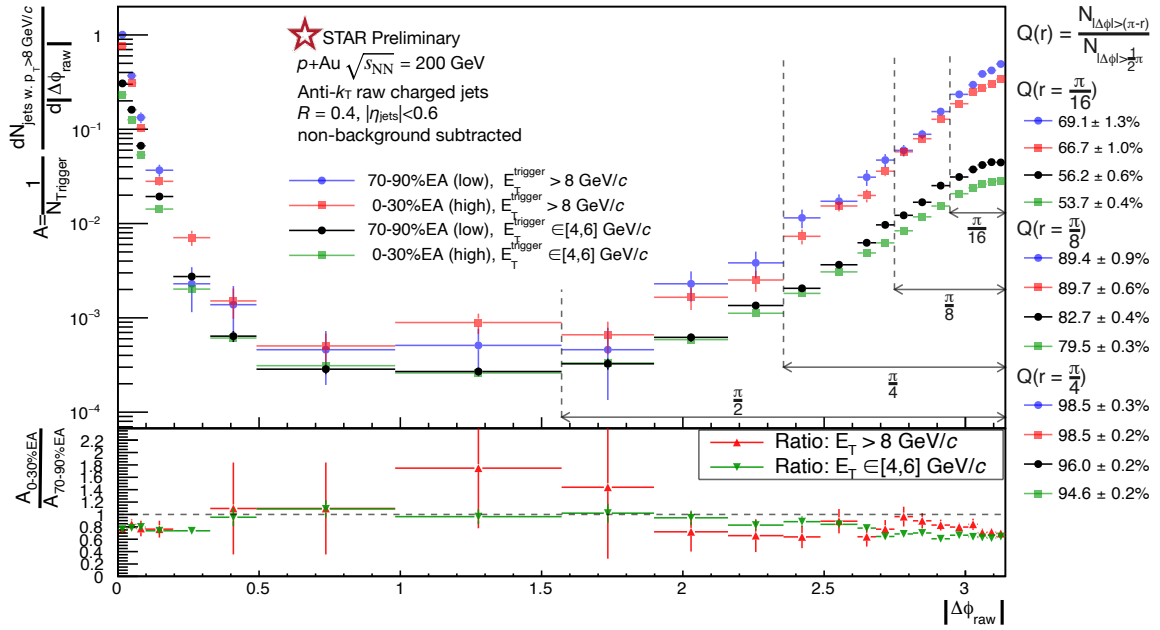
**Figure 1:** Top: raw, uncorrected, jet spectra per trigger in bins of high (0-30%) and low (70-90%) EA, sub-divided into bins of  $|\phi_{\text{jet}} - \phi_{\text{trigger}}|$ . Bottom: ratios of high-to-low EA jet spectra. EA is measured in the Au-going direction at  $\eta \in [-5, -3.4]$ .

52 the EA acceptance. These events would thereby preferentially lower the jet spectra per trigger while  
 53 raising the measured EA. This has already been demonstrated not to be the case for this STAR  
 54 measurement, as reported in [12]. Third, a correlation would result from actual jet modification  
 55 occurring preferentially in high EA events relative to low EA events; for example, jet quenching in  
 56 a QGP.

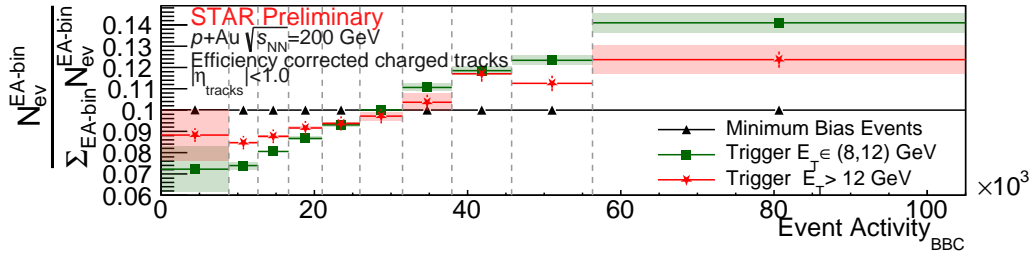
57 There are several empirical indications that the jet spectra suppression in high EA p+Au  
 58 collisions measured by STAR results from phase space constraints and not from actual modification  
 59 of generated jets. First, the comparable level of suppression for near-side and away-side jets in  
 60 Fig. 1 is indicative of a common phase space constraint on the hard scattering itself and EA. When  
 61 modified by a QGP, at least in A+A collisions, path-length differences modify spectra differently  
 62 for near- and away-side jets.

63 Figure 2 presents jet acoplanarity, the correlation between  $|\Delta\phi| = |\phi_{\text{jet}} - \phi_{\text{trigger}}|$ . It shows the  
 64 overall suppression of high-EA jet production. However, the recoil shape does not broaden as would  
 65 be expected from jet quenching. Such a broadening has been reported by ALICE in  $\sqrt{s} = 13$  TeV  
 66 high multiplicity  $pp$  collisions. However, it was reported at this conference that a principle cause  
 67 may be a selection bias towards tri-jet events [13]; whereas tri-jet events have negligible impact on  
 68 the STAR results.

69 Figure 3 presents a more direct empirical indication that the STAR jet suppression is due  
 70 to phase space constraints. Raising the transverse energy requirement on triggers from  $E_T \in$   
 71  $(8, 12)$  GeV to  $E_T > 12$  GeV corresponds to a  $\sim 20\%$  drop in the the number of events in the highest  
 72 EA bin in proportion to those in the lowest EA bin. This suppression of  $E_T > 12$  GeV (for this EA  
 73 selection) is comparable in magnitude to the suppression of jet spectra in Fig. 1. We also note new  
 74 jet mass measurements in high and low EA p+Au collisions, which are normalized per jet rather  
 75 than per trigger, also revealed no modification relative to  $pp$  collisions [14].



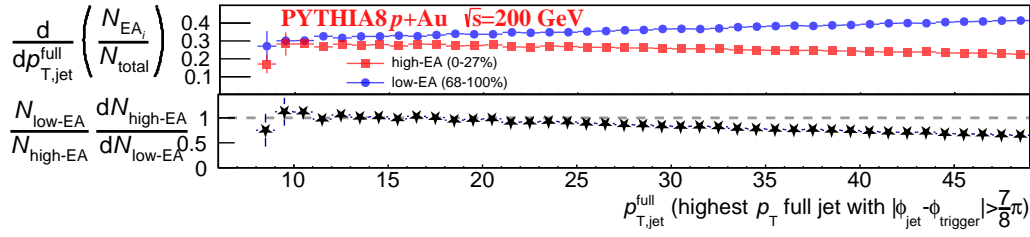
**Figure 2:** Acoplanarity:  $|\phi_{jet} - \phi_{trig}|$  for all raw, anti- $k_T$ ,  $R = 0.4$ ,  $p_T > 8$  GeV/c, jets, within  $|\eta_{jet}| < 0.6$ . Jets are uncorrected; however, the correction in  $\phi$  for jets is expected to be minimal.



**Figure 3:** Fraction in each EA bin for MB events and events containing high- $E_T$  triggers. Decile boundaries are defined to contain 10% of the minimum bias events in each bin.

76 Figure 4 shows correlations, from a PYTHIA 8 Agantyr model simulation, between EA and  
 77 hardness of leading recoil jets at mid-rapidity for  $\sqrt{s_{NN}} = 200$  GeV  $p+Au$  collisions. Binned by  
 78 jet  $p_T$ , the top panel shows the percentage of events in the high (low) EA bins, which decreases  
 79 (increases) with increasing jet  $p_T$ . The bottom panel shows the corresponding anti-correlation  
 80 of EA to jet  $p_T$ . Further investigations show that this behavior is not the result of a trivial  
 81 autocorrelation caused by dijet fragments in the EA acceptance. Additionally, the simulations do  
 82 not have jet quenching mechanisms; as such, the correlation modeled is between cross sections of  
 83 jet production and EA production.

84 Investigations into the cause of the modeled correlations are ongoing. However, we note  
 85 here that the Agantyr model does correlate energy and momentum conservation among all partonic  
 86 scatterings in  $p+Au$  collisions [15]. This correlates a high- $p_T$  jet production from one hard scattering  
 87 with decreased EA elsewhere from other scatterings. However, we have not isolated the strength  
 88 of this effect and PYTHIA 8  $pp$  simulations with multi-parton interactions turned off (which does  
 89 not have this effect) also demonstrate an anti-correlation between EA and mid-rapidity jet  $p_T$ . In



**Figure 4:** PYTHIA 8 Agantyr simulation of events triggered by a neutral particle with  $E_T > 8$  GeV in  $|\eta| < 1$ . Events are binned by EA (the sum of particles within  $\eta \in [-5, -3.5]$ ) and the leading recoil full jet  $p_T$  at mid-rapidity. Bottom panel: Ratio of the shapes of the high-EA to low-EA event distributions.

90 addition to the Agantyr model, other theories, using color fluctuations in the proton which may result  
 91 in a “shrinking proton” at high Bjorken- $x$  [16], have been proposed to explain this anti-correlation  
 92 in the context of the ATLAS and PHENIX  $R_{p/dA}$  results. One model calculation found that energy  
 93 conservation alone is sufficient to account for the  $R_{p/dA}$  modification [8], although it could not be  
 94 determined if color transparency-like effects are also indicated.

95 In conclusion, semi-inclusive jet spectra at STAR from  $p$ +Au collisions at  $\sqrt{s_{NN}} = 200$  GeV are  
 96 distinctly suppressed in high EA events relative to low EA events. These results confirm modification  
 97 in a similar Bjorken- $x$  range as the  $R_{p/dA}$  modification measured by PHENIX and ATLAS [2, 3].  
 98 However, acoplanarity measurements, PYTHIA-8 studies, and the spectra themselves, indicate that  
 99 the new STAR results originate primarily from phase space constraints on the hard scattering cross  
 100 section relative to generating high EA, rather than from jet quenching effects. These results motivate  
 101 further studies to determine if this mechanism of phase space constraints are sufficient to explain  
 102 all the reported jet measurements in  $p/d$ +A collisions at RHIC and the LHC.

## 103 References

- 104 [1] CMS, V. Khachatryan *et al.*, JHEP **09**, 091 (2010), 1009.4122.  
 105 [2] PHENIX, A. Adare *et al.*, Phys. Rev. Lett. **116**, 122301 (2016), 1509.04657.  
 106 [3] ATLAS, G. Aad *et al.*, Phys. Lett. **B748**, 392 (2015), 1412.4092.  
 107 [4] ALICE, J. Adam *et al.*, Phys. Lett. **B749**, 68 (2015), 1503.00681.  
 108 [5] CMS, V. Khachatryan *et al.*, Eur. Phys. J. **C76**, 372 (2016), 1601.02001.  
 109 [6] ALICE, J. Adam *et al.*, Eur. Phys. J. **C76**, 271 (2016), 1603.03402.  
 110 [7] ALICE, S. Acharya *et al.*, Phys. Lett. **B783**, 95 (2018), 1712.05603.  
 111 [8] M. Kordell and A. Majumder, Phys. Rev. **C97**, 054904 (2018), 1601.02595.  
 112 [9] CMS, S. Chatrchyan *et al.*, Eur. Phys. J. C **74**, 2951 (2014), 1401.4433.  
 113 [10] ALICE, J. Adam *et al.*, Phys. Lett. B **746**, 385 (2015), 1503.03050.  
 114 [11] ALICE, S. Acharya *et al.*, Phys. Lett. B **776**, 249 (2018), 1702.00804.  
 115 [12] STAR, D. Stewart, Quark Matter Proceedings (2020), 2002.06133.  
 116 [13] ALICE, F. Krizek, These Proceedings.  
 117 [14] STAR, I. Mooney, These Proceedings.  
 118 [15] H. Shah, MDPI Proc. **10**, 18 (2019).  
 119 [16] D. McGlinchey, J. L. Nagle, and D. V. Perepelitsa, Phys. Rev. **C94**, 024915 (2016), 1603.06607.