



Measurement of the cumulants of
net-proton multiplicity distribution in
Au+Au collisions at
 $\sqrt{s_{NN}} = 7.7 - 200 \text{ GeV}$

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Fluctuation of Conserved Quantities

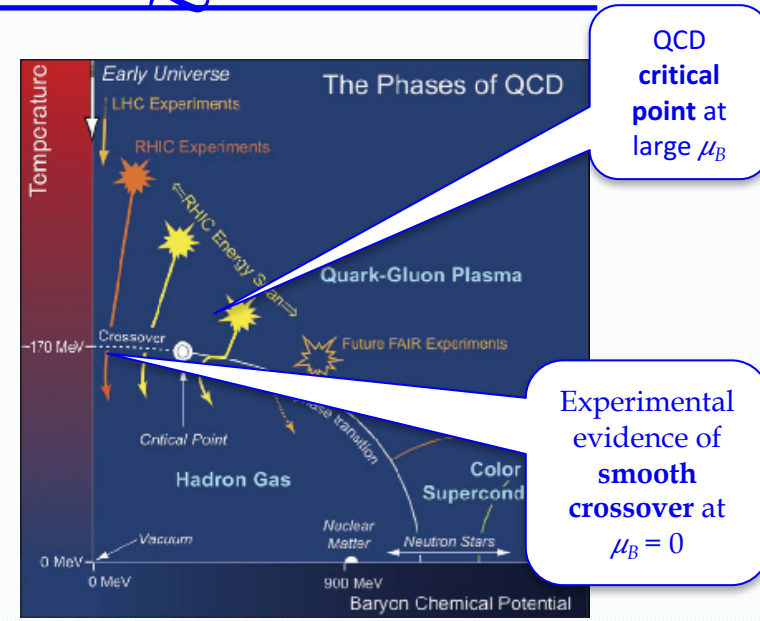
➤ Connection to the susceptibility of the system (χ)

$$\chi_q^{(n)} = \frac{1}{VT^3} \times C_{n,q} = \frac{\partial^n (p/T^4)}{\partial (\mu_q/T)^n} \quad q = B, Q, S$$

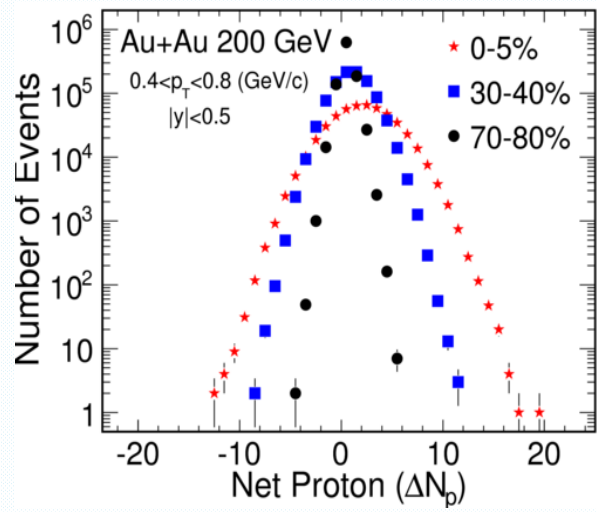
$$\frac{\chi_q^{(4)}}{\chi_q^{(2)}} = \frac{C_{4,q}}{C_{2,q}} \quad \frac{\chi_q^{(6)}}{\chi_q^{(2)}} = \frac{C_{6,q}}{C_{2,q}}$$

➤ Higher order cumulants are more sensitive to the signatures of QCD phase transition

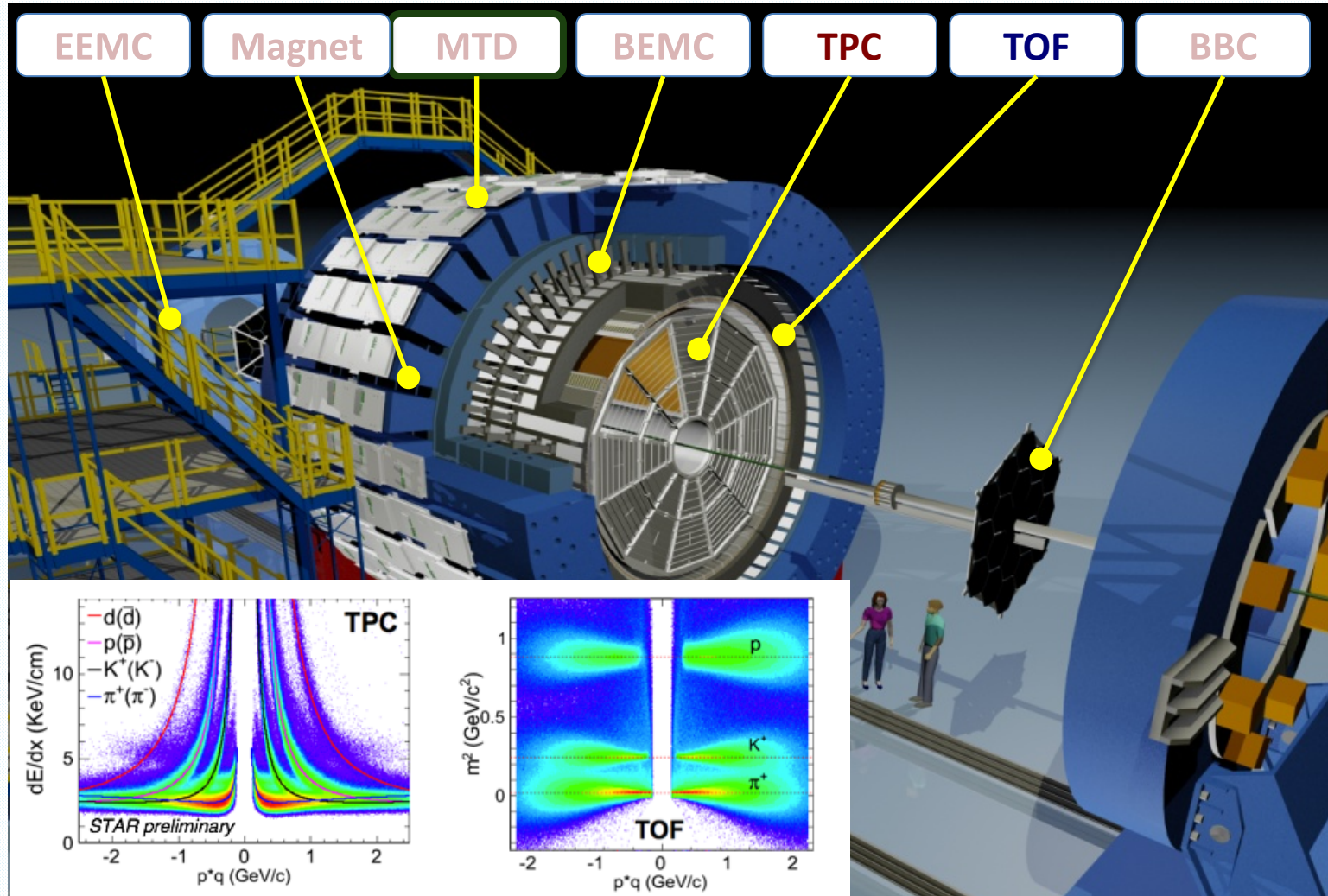
M. A. Stephanov, *Phys. Rev. Lett.* 102, 032301 (2009).
 M. Asakawa, S. Ejiri and M. Kitazawa, *Phys. Rev. Lett.* 103, 262301 (2009).
 M. A. Stephanov, *Phys. Rev. Lett.* 107, 052301 (2011).



STAR Collaboration, *Phys.Rev.Lett.* 105 (2010) 022302



STAR Detector



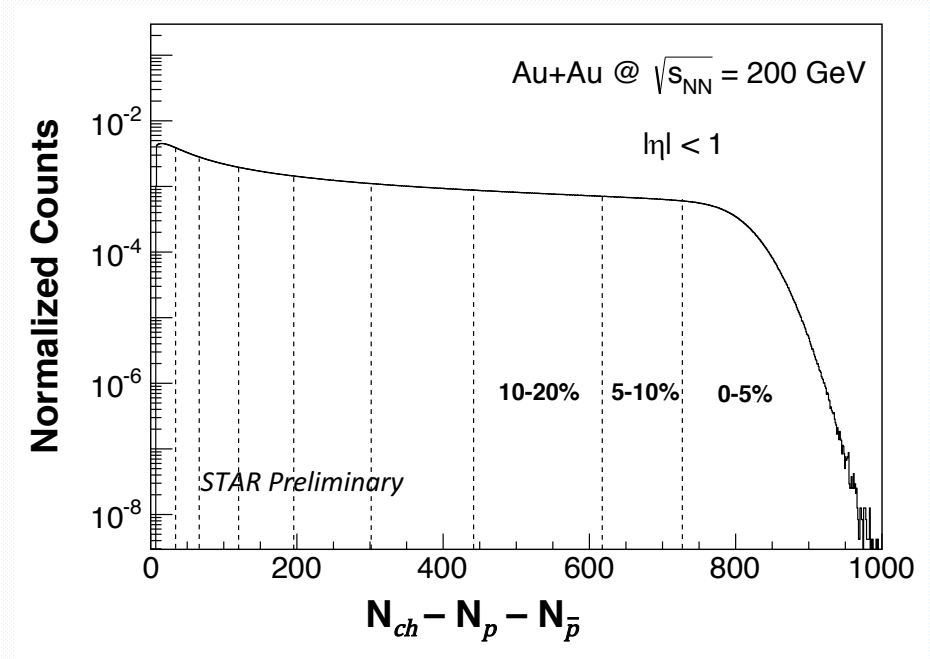
Analysis Technique

➤ Centrality determination

Use charged particles within $|\eta| < 1$, other than protons and anti-protons, to avoid auto-correlations

➤ Centrality bin width correction

Evaluate cumulants for each centrality bin to suppress volume fluctuations



X. Luo and N. Xu, arXiv:1701.02105

STAR Collaboration, Phys.Rev.Lett. 105 (2010) 022302.

STAR Collaboration, Phys.Rev.Lett. 113 (2014) 092301 .

Analysis Technique

➤ Error estimation

Statistical errors are based on **Bootstrap technique** or the **Delta Theorem**.

$$Error (C_r) \propto \frac{\sigma^r}{\sqrt{n}}$$

$$Error (C_r/C_2) \propto \frac{\sigma^{r-2}}{\sqrt{n}}$$

σ : width of the distribution

n : Number of events

*B. Efron et al. An Introduction to Bootstrap, Chapman & Hill (1993).
X. Luo, J. Xu, B. Mohanty, N. Xu, J. Phys. G 40, 105104 (2013)*

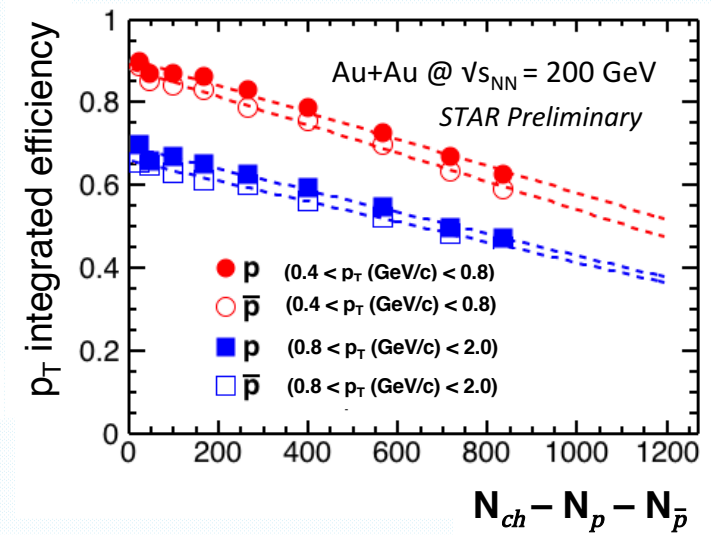
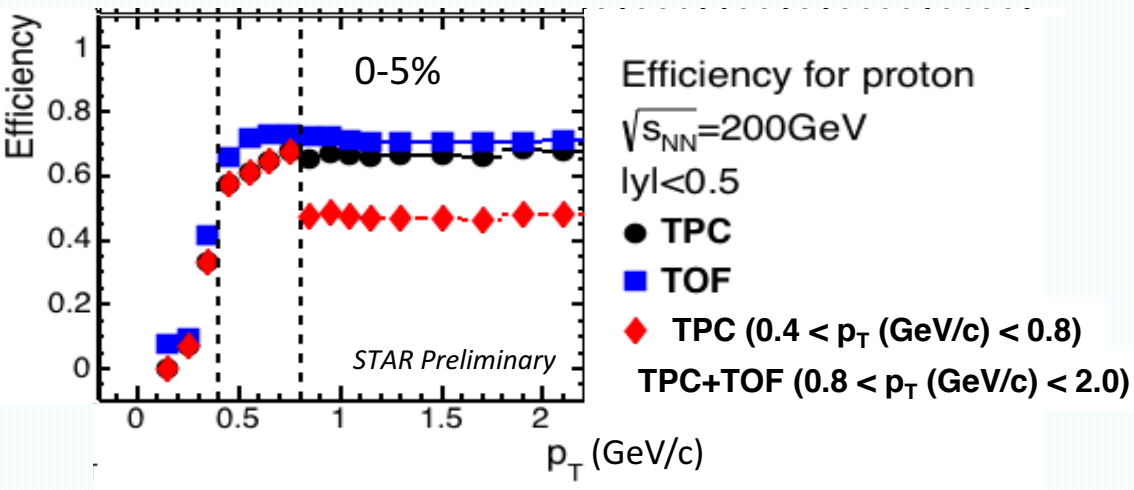
➤ Efficiency correction

Express the cumulants in terms of the factorial moments or factorial cumulants, which can be easily efficiency corrected by assuming **binomial response function for efficiency**.

Based on factorial cumulants: T. Nonaka, M. Kitazawa and S. Esumi, In preparation.

Based on factorial moments: A. Bzdak and V. Koch, PRC91, 027901 (2015). X. Luo, PRC91, 034907(2015). X. Luo and N. Xu, 1701.02105

Detector Efficiency



$$\langle \epsilon \rangle = \frac{\int_{p_{T1}}^{p_{T2}} \epsilon(p_T) f(p_T) dp_T}{\int_{p_{T1}}^{p_{T2}} f(p_T) dp_T}$$

- p_T - integrated efficiency is calculated as a function of multiplicity
- Efficiency correction is applied at each multiplicity bin



Search for the QCD Critical Point

Cumulants and Correlation Function

- Higher order cumulants are more sensitive to the correlation lengths

$$C_2 = \langle (\delta N)^2 \rangle \sim \xi^2; \quad C_3 = \langle (\delta N)^3 \rangle \sim \xi^{4.5}; \quad C_4 = \langle (\delta N)^4 \rangle \sim \xi^7 \quad \text{with} \quad \delta N = N - \langle N \rangle$$

*M. A. Stephanov, Phys. Rev. Lett. 102, 032301 (2009); M. A. Stephanov, Phys. Rev. Lett. 107, 052301 (2011).
M. Asakawa, S. Ejiri and M. Kitazawa, Phys. Rev. Lett. 103, 262301 (2009); Y. Hatta, M. Stephanov, Phys. Rev. Lett. 91, 102003 (2003).*

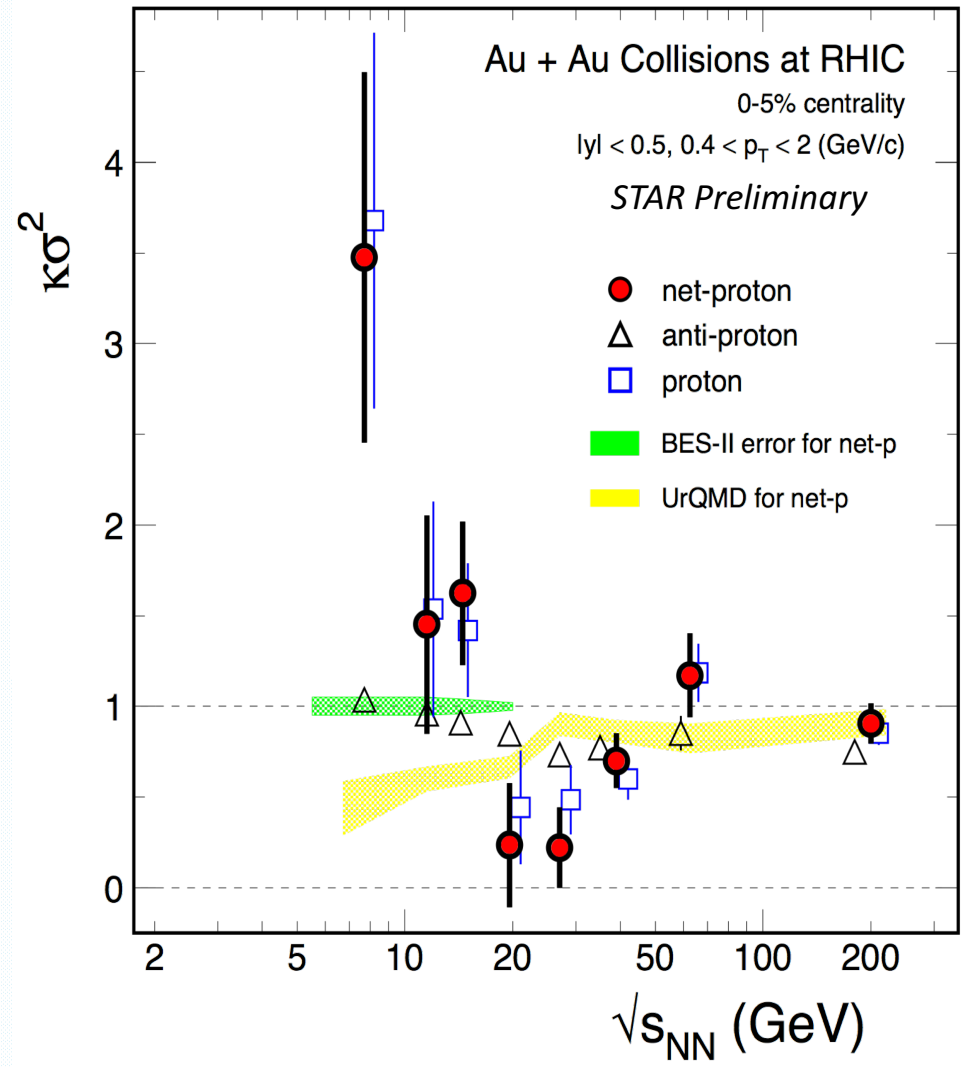
- Relation between cumulants (C_n) and correlation functions (\hat{K}_n)

$\hat{K}_1 = C_1$	$C_1 = \langle N \rangle$
$\hat{K}_2 = C_2 - C_1^2$	$C_2 = \langle N \rangle^2 + \hat{K}_2$
$\hat{K}_3 = C_3 - 3C_2C_1 + 2C_1^3$	$C_3 = \langle N \rangle^3 + 3\hat{K}_2\langle N \rangle + \hat{K}_3$
$\hat{K}_4 = C_4 - 6C_3C_1 + 11C_2^2 - 6C_2C_1^2 + 3C_1^4$	$C_4 = \langle N \rangle^4 + 7\hat{K}_2\langle N \rangle^2 + 6\hat{K}_3\langle N \rangle + \hat{K}_4$

$$\hat{K}_2 \propto \xi^2, \quad \hat{K}_3 \propto \xi^{4.5}, \quad \hat{K}_4 \propto \xi^7$$

*B. Ling, M. Stephanov, Phys. Rev. C 93, 034915 (2016); A. Bzdak, V. Koch, N. Strodthoff, arXiv:1607.07375;
A. Bzdak, V. Koch, V. Skokov, arXiv:1612.05128*

Fourth Order Net Proton Fluctuation

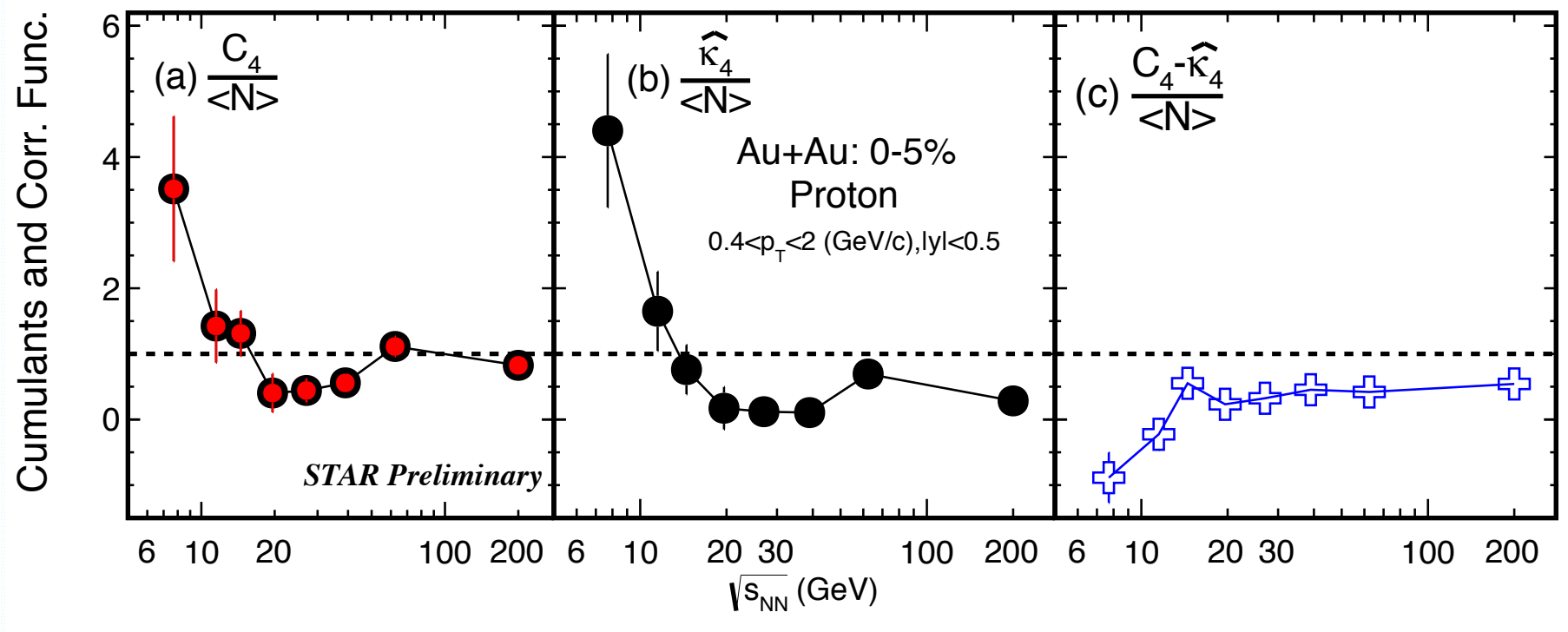


➤ Non-monotonic energy dependence is observed for 4th order net-proton, proton fluctuations in most central Au+Au collisions.

$$\kappa\sigma^2 = \frac{C_4}{C_2}$$

➤ UrQMD results show monotonic decrease with decreasing collision energy.

Contributions from Four-Particle Correlations

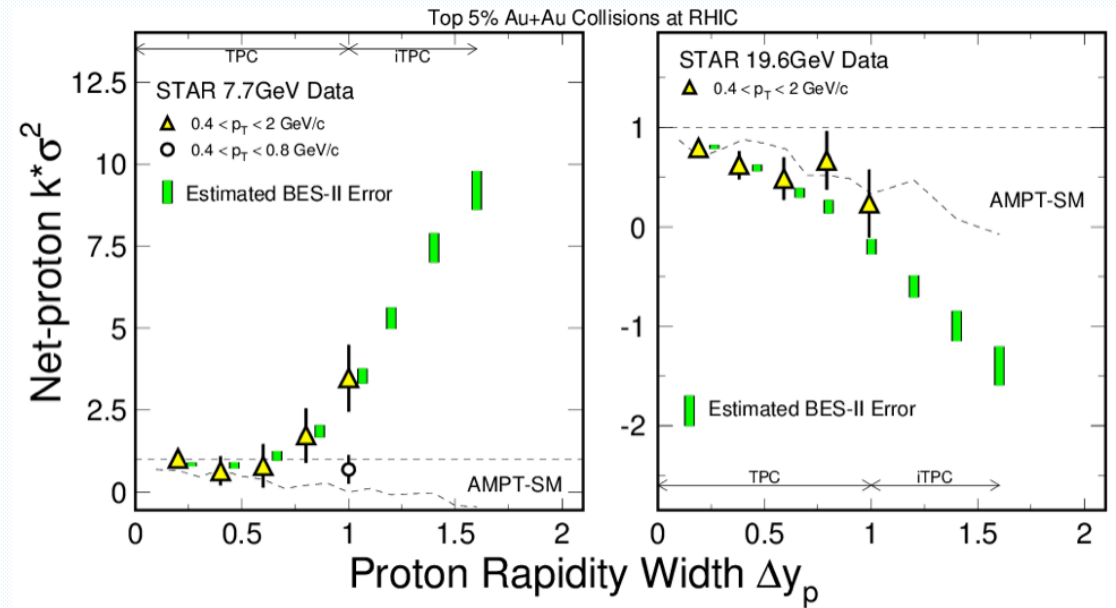


- Four-particle correlations contribute dominantly to the observed non-monotonicity.

Summary - I

- Non-monotonic energy dependence of net-proton and proton C_4/C_2 is observed for 0–5% central Au+Au collisions.
- Four-particle correlations contribute dominantly to the observed non-monotonicity.

➤ More data will be collected in BES-II at $\sqrt{s_{NN}} = 7.7 - 19.6$ GeV in 2019–2020 with detector upgrades.



STAR Collaboration, <https://drupal.star.bnl.gov/STAR/starnotes/public/sn0619>



*The measurement of net-proton
sixth-order cumulant at small μ_B
and its comparison to Lattice QCD*

T. Nonaka
Poster ID 239

Connections with Lattice QCD

- LQCD predicts a “crossover” for $\mu = 0$

Y. Aoki, Nature 443, 675(2006)

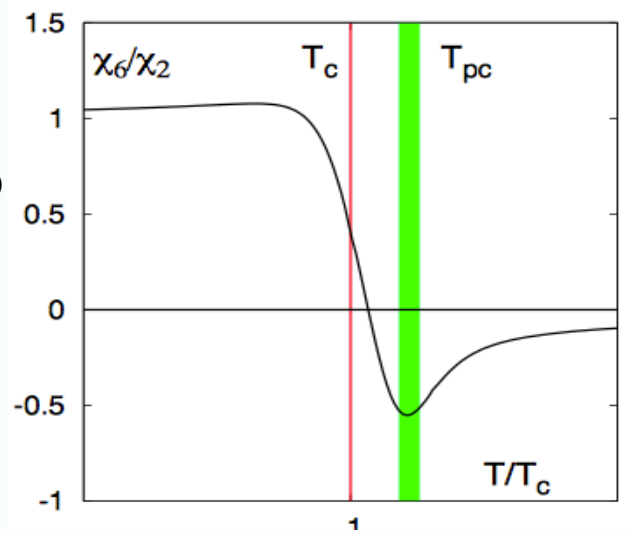
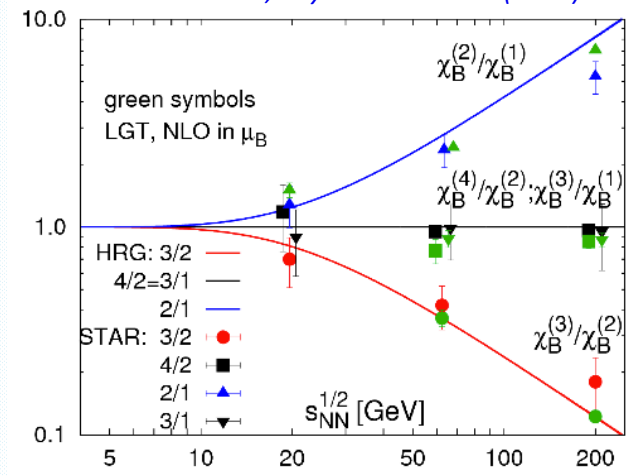
- No established approach to do QCD calculations at finite baryon chemical potential

- By putting $\mu_Q = \mu_S = 0$ and using Taylor expansion, the equation of state for finite μ_B :

$$\frac{P(T, \mu_B) - P(T, 0)}{T^4} = \frac{1}{2} \chi_2^B(T) \left(\frac{\mu_B}{T}\right)^2 \times \left[1 + \frac{1}{4} \frac{\chi_4^B(T)}{\chi_2^B(T)} \left(\frac{\mu_B}{T}\right)^2 + \frac{1}{360} \frac{\chi_6^B(T)}{\chi_2^B(T)} \left(\frac{\mu_B}{T}\right)^4 \right] + O(\mu_B^8)$$

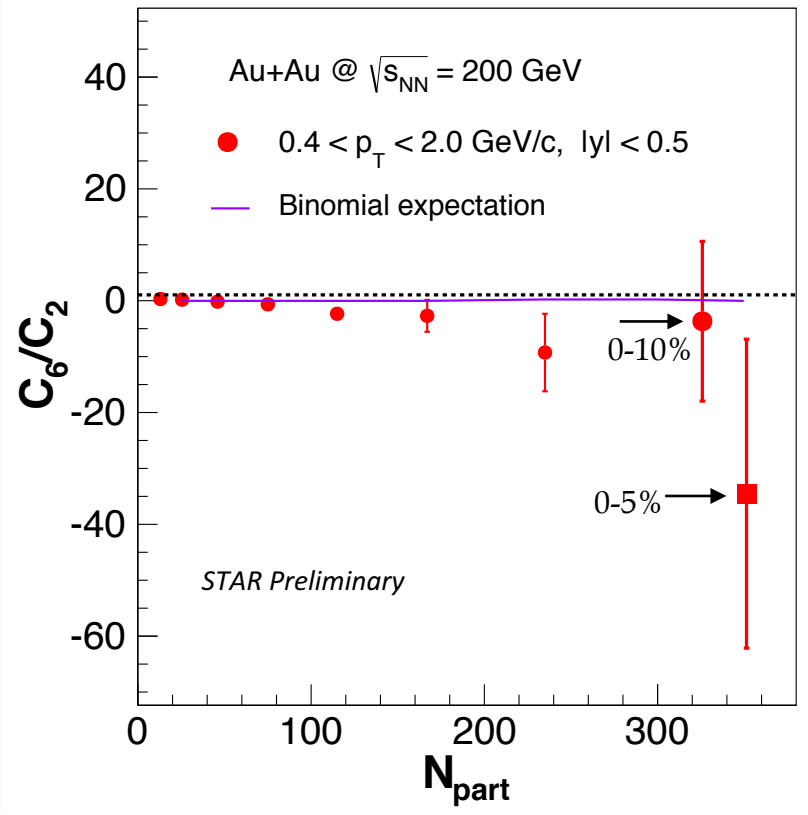
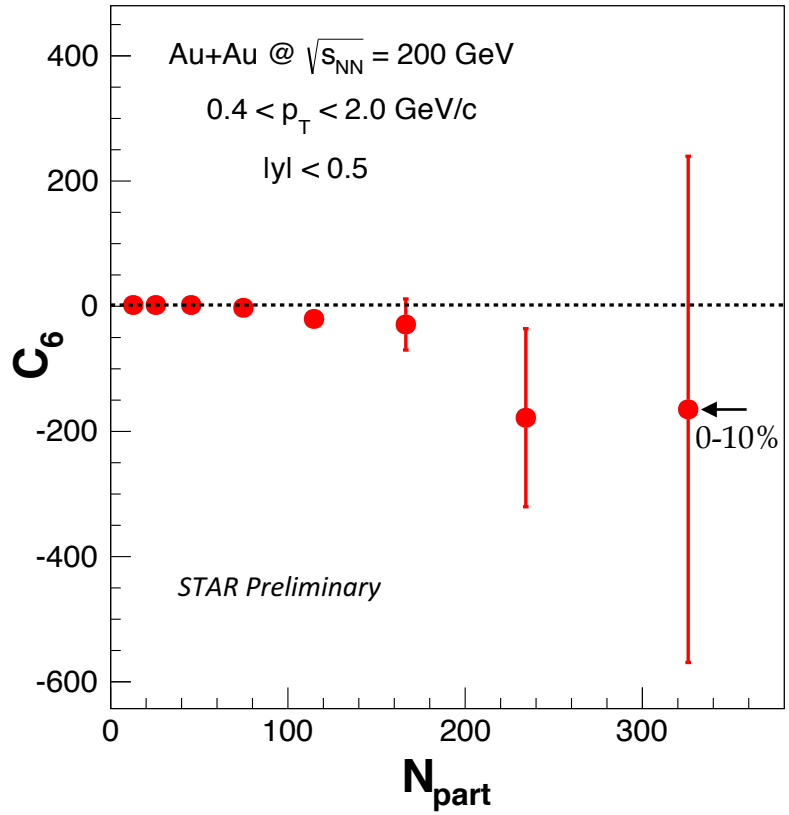
- The sixth order cumulants of baryon number is expected to be negative at chiral transition temperature

F. Karsch and K. Redlich, Phys. Lett. B 695, 136 (2011)
STAR Collaboration, Phys.Rev.Lett. 112 (2014) 032302.



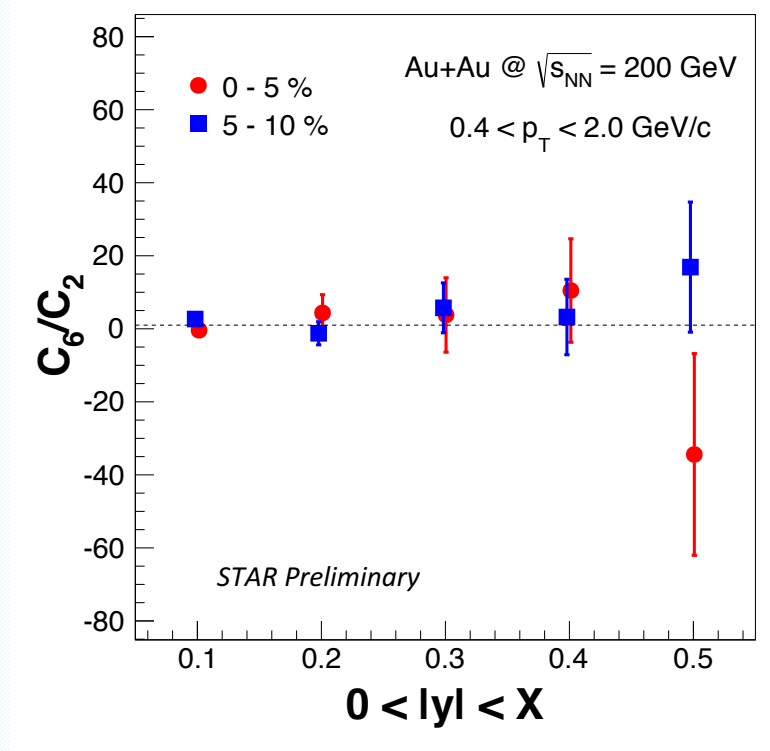
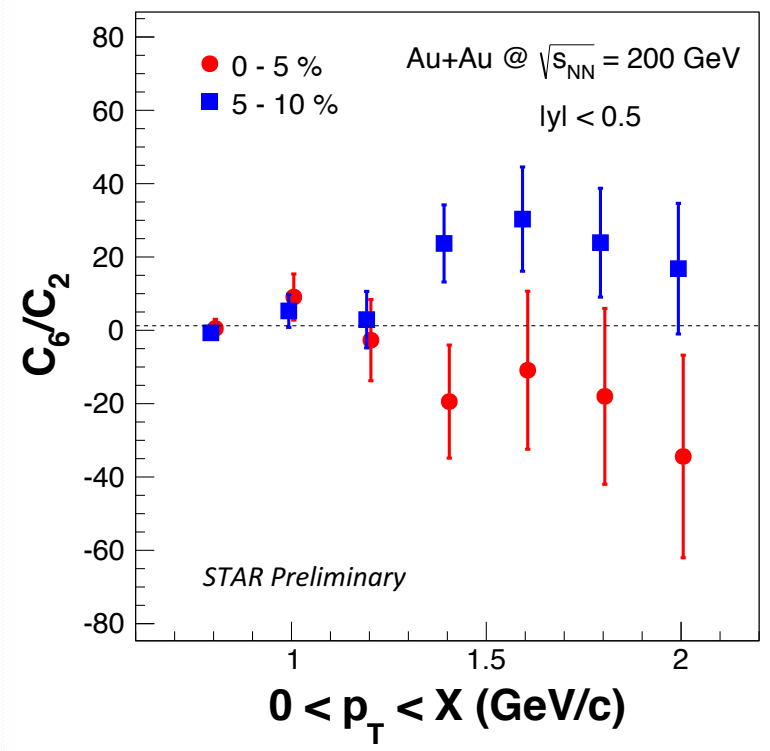
B. Friman, F. Karsch, K. Redlich, V. Skokov Eur. Phys. J. C 71, 1694 (2011)

Net-Proton Sixth Cumulant



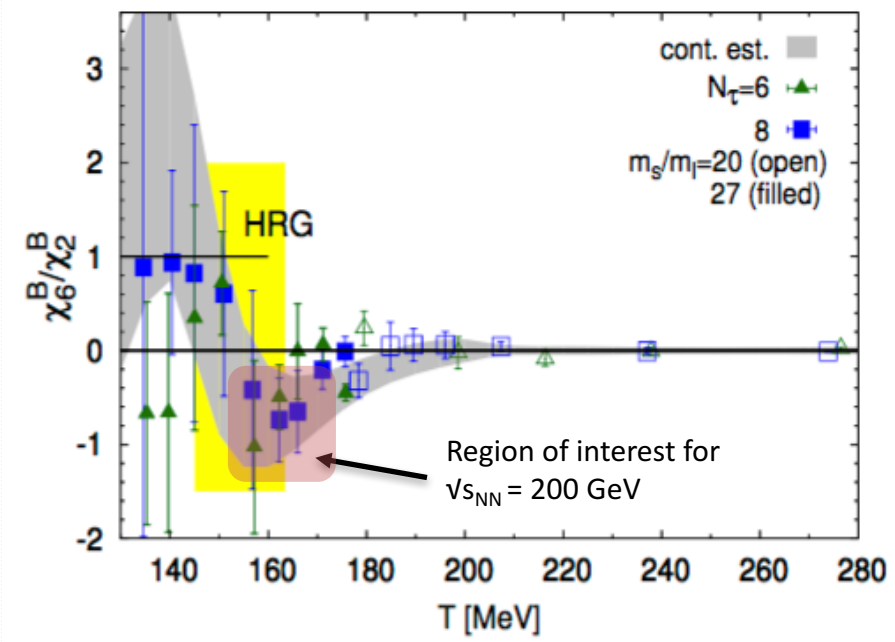
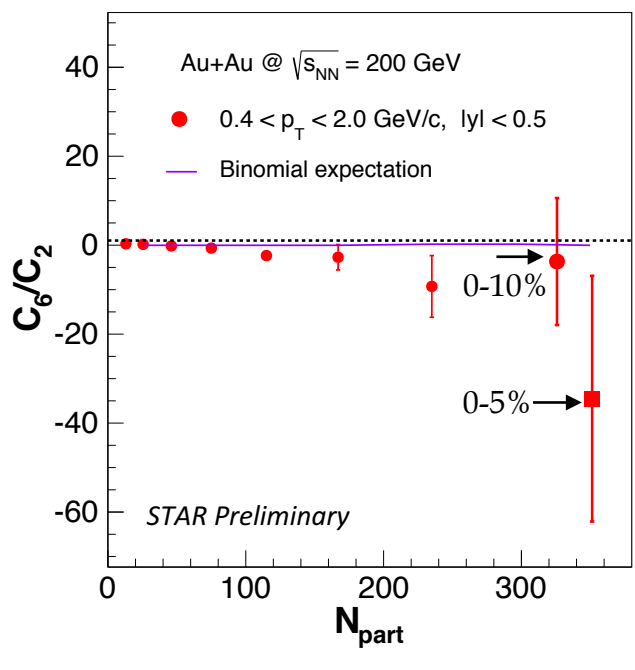
- Combined data of Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV from year 2010 and 2011: 200 M (0 - 10 %) and 650 M (10 - 80 %) events
- The C_6 and C_6/C_2 is negative for central collisions with large uncertainties

Acceptance Dependence



- Around 160 M events are analyzed for 0-10% central collisions for Au + Au at $\sqrt{s_{NN}} = 200$ GeV with central trigger from year 2010

Comparison with LQCD



A. Bazavov et al, arXiv:1701.04325.

- C_6/C_2 for most central collisions is negative with large uncertainties
- Some differences between experiment and LQCD measurements:
 1. Net-proton is not equivalent to net-baryon
 2. Limited phase space
 3. $\mu_B \neq 0$ ($\mu_B \sim 20$ MeV at 200 GeV)



Summary - II

- We report the efficiency corrected sixth order cumulant of the net-proton multiplicity distribution for Au+Au collisions at $\sqrt{s_{\text{NN}}} = 200$ GeV.
- Centrality, transverse momentum and rapidity dependence of the ratio C_6/C_2 are presented.
- C_6 and C_6/C_2 are negative for central collisions with large statistical uncertainty.
- Assessment of systematic errors is underway.
- Combining the data taken in year 2014 and 2016, with more than 2 billion events, we can get a better control on the statistical errors for C_6/C_2 .



Thank you !



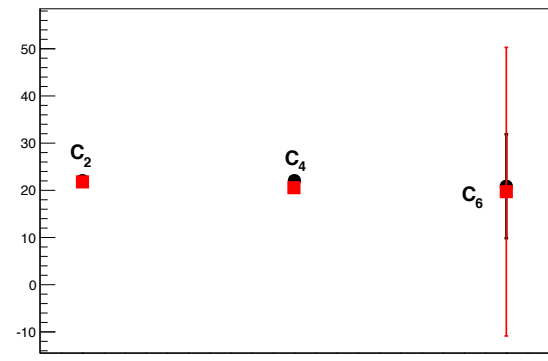
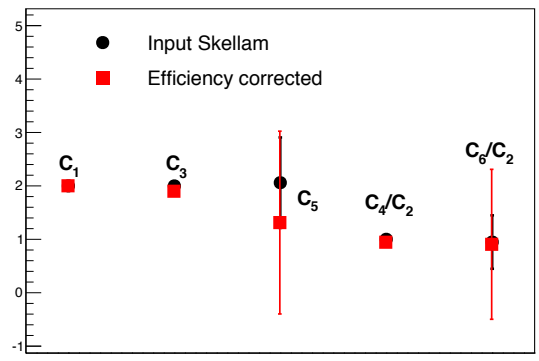
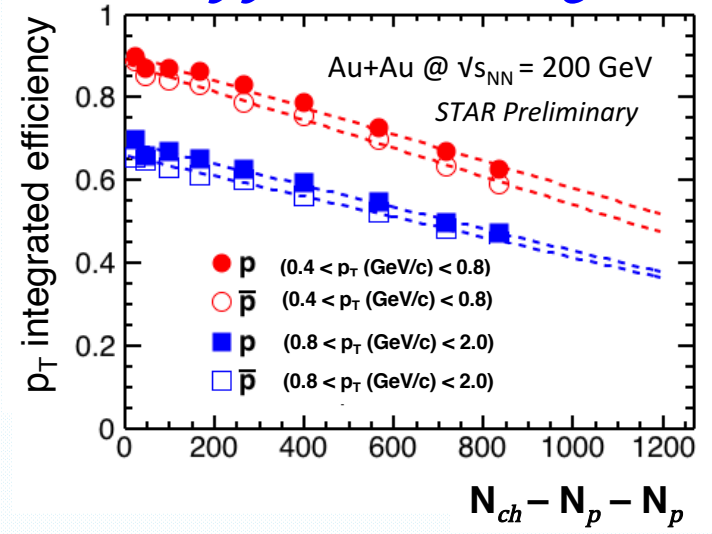
Back-up

Multiplicity Dependent Efficiency

$$\varepsilon(N) = \varepsilon_0 + \varepsilon' (N - \langle N \rangle)$$

↑ averaged
slope
event by event
↑ mean of Poisson

	ε_0	ε'	$\langle N \rangle$
proton	0.7	-0.0003	12
antiproton	0.68	-0.0003	10



- Analyzed for 1B events
- Protons and anti-protons are sampled from Poisson distribution with Binomial efficiency
- Deviations for efficiency corrected higher order cumulants from the input Skellam are within statistical errors

A. Bzdak et al Phys.Rev. C94, 064907 (2016).