

1 **Highlights from the STAR experiment***

2 PRITHWISH TRIBEDY (FOR THE STAR COLLABORATION)

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5 Despite the challenges of pandemic the years 2020-21 were quite suc-
6 cessful for STAR. We completed the Beam Energy Scan program phase 2
7 and installed the forward upgrade with which STAR finished data-taking
8 for polarized p+p collisions at 510 GeV. In this contribution, we discuss
9 STAR results on five different topics that were presented in twenty one
10 parallel talks, forty-seven posters, and two flash talks at the Quark Matter
11 2022 conference.

12 **1. Isobar collisions and strong field effects**

13 An important program at STAR since the previous Quark Matter has
14 been the blind analysis of isobar data at $\sqrt{s_{NN}} = 200$ GeV [1] to search for
15 the chiral magnetic effect (CME). Isobars were collided to utilize the fact
16 that Ru+Ru collisions produce larger magnetic fields than the Zr+Zr colli-
17 sions. Therefore, the ratio of the CME-sensitive observables in Ru+Ru over
18 Zr+Zr has to be greater than unity in the presence of CME. The run was
19 specially designed to reduce the systematics in the ratio by alternating two
20 species. This provides the best possible control of signal and background
21 compared to all previous experiments for the search for CME. A precision
22 in the measurements down to 0.4% was achieved in the blind analysis of the
23 isobar data, and no predefined signatures of CME were observed. At this
24 Quark Matter, we present important progress toward estimating the back-
25 ground expectations by incorporating the multiplicity difference between
26 the two isobars and the non-flow effects [2].

27 In STAR we continue to search for the electromagnetic (EM) field driven
28 effects. The consequence of the Faraday and Hall effects have been predicted
29 to lead to differences in the slope of directed flow ($d\Delta v_1/dy$) with rapidity
30 between positive and negative particles. Our measurements in Au+Au and
31 isobar collisions for protons and anti-protons show a difference in slope that
32 changes sign from central to peripheral events. In another measurement

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33 the v_1 slope difference is studied with respect to various values of electric
 34 charge difference between combinations of hadrons that are made up of
 35 only produced quarks. We see significant splitting with stronger strength
 36 in Au+Au collisions at $\sqrt{s_{NN}} = 27$ GeV compared to 200 GeV. These
 37 observations of splitting of v_1 slope cannot be explained by baryon transport.
 38 Model studies to interpret these results in the context of EM-field driven
 39 effects are highly anticipated [3].

40 Similar to strong EM fields, the presence of strong vortical fields is stud-
 41 ied via measurements of the global polarization of Λ hyperons as a function
 42 of collision energy. Recently, we have extended our measurements to fixed
 43 target 3 GeV for Λ , fixed target 7.2 GeV for Λ , and to 19.6 GeV for $\Lambda + \bar{\Lambda}$.
 44 Our new measurements with high statistics data at 200 GeV for Au+Au and
 45 isobar collisions for Λ and $\bar{\Lambda}$ enable us to test three following predictions.
 46 We do not see a splitting between Λ and $\bar{\Lambda}$ polarization due the opposite
 47 sign of their magnetic moment. For either Λ or $\bar{\Lambda}$ we do not see a difference
 48 in the magnitude of global polarization due to the B-field difference between
 49 Ru+Ru and Zr+Zr collisions. Finally, we do not see any significant system
 50 size dependence while going from the larger size of the Au+Au to isobars
 51 at a given centrality [4].

52 A similar kind of effect is the spin alignment of vector mesons. The
 53 quantity of interest is the spin alignment coefficient (ρ_{00}) as a function of
 54 collision energy. For the ϕ -mesons we see almost 8.4σ deviation from the
 55 baseline of $1/3$; the K^{*0} results are consistent with $1/3$ in Au+Au collisions
 56 in the range of $\sqrt{s} = 7.7 - 200$ GeV. An outstanding question is what causes
 57 this vector meson spin alignment. For our measurement, we see a model
 58 that includes a very strong vector meson field of the order of m_π^2 can provide
 59 an explanation for ϕ -mesons. What about the K^{*0} ? For that we perform
 60 new measurements of charged and neutral K^{*0} in isobar collisions to gain
 61 more insights [5].

62 As an another way to study the strong field effects, we try to test if the
 63 very early EM-field difference between the two isobars is reflected in the
 64 low p_T photon induced processes. The first process of interest is $\gamma + \gamma \rightarrow$
 65 $e^+ + e^-$, also known as the Breit-Wheeler process. The cross section of this
 66 process is expected to go as $\sigma \sim Z^4$, where Z is the atomic number of the
 67 nucleus emitting the soft photon. So, if we measure it in Ru+Ru over Zr+Zr
 68 collisions we expect to see a $(44/40)^4$ scaling. We perform measurements
 69 of the ratio of the $e^+ + e^-$ yields in Ru+Ru over Zr+Zr collisions as a
 70 function of p_T in 40 – 80% centrality. At the lowest p_T we see above 3σ
 71 deviation of the yield ratio from unity and a value that is close to the
 72 expectation of $(44/40)^4$ scaling and QED predictions. A similar process is
 73 the photoproduction of J/ψ , which is expected to follow a Z^2 scaling. For
 74 that measurement we also see some hints of deviations for the yield ratio

75 from unity in the direction of the Z^2 scaling. Therefore, isobar data suggest
 76 that low p_T photon induced processes follow “Z” scaling, consistent with
 77 the EM-field difference between the isobar collision systems [6].

78 The photon induced processes can also be used to constrain the nuclear
 79 charge and the mass radius in a novel way. The commonly used Woods-
 80 Saxon parameterization to characterize the spatial distribution of nucleons
 81 inside colliding nuclei includes parameters such as the charge radius and
 82 the skin depth. In modelling of heavy-ion collisions we normally use fixed
 83 values of such parameters extracted from low energy electron scattering
 84 data. Measuring the $\gamma + \gamma \rightarrow e^+ + e^-$ process, we have a novel way to
 85 constrain the combination of skin depth and the charge radius. On the other
 86 hand, using the diffractive vector meson production through the $\gamma + A \rightarrow \rho^0$
 87 process, it is possible to constrain the strong interaction or the gluonic radius
 88 of the gold and uranium nuclei [6].

89 2. New insights on collective effects

90 The nuclear structure leaves imprints in the final observables through
 91 collective effects. Nuclear parameters, particularly the neutron skin and de-
 92 formation, can be constrained through measurements of the ratio of various
 93 quantities in isobar collisions. Our measurements of the ratio of multiplicity
 94 distribution as well as the net-charge density at midrapidity in Ru+Ru over
 95 Zr+Zr collisions indicate that Zr nucleus has a thicker neutron skin than
 96 that of Ru nucleus. The measurements of the ratio of v_2 , v_3 , and trans-
 97 verse momentum fluctuations deviate from unity in central collisions in a
 98 significant way. These measurements indicate that Ru has a larger quadru-
 99 ple deformation (β_2) than Zr, while Zr has larger octuple deformation (β_3)
 100 than Ru. These studies in isobar collisions have pioneered new ways of
 101 constraining nuclear deformation parameters [7].

102 Nuclear deformation can also be constrained by another observable,
 103 which is known as the Pearson coefficient between flow harmonics v_n and
 104 mean transverse momentum [p_T]: $\rho(v_n^2, [p_T])$. This observable measures the
 105 correlation between shape and size of the fireball. The measurements of
 106 $\rho(v_2^2, [p_T])$ as a function of centrality in U+U collisions change sign from
 107 positive to negative values in central collisions indicating a highly deformed
 108 shape of uranium nucleus. In Au+Au collisions no sign change is observed
 109 indicating very little deformation. The measurements of $\rho(v_3^2, [p_T])$ also
 110 stay positive at all centralities and serve as a data-driven baseline. In order
 111 to investigate if the observable $\rho(v_2^2, [p_T])$ is sensitive to the lifetime and
 112 the nature of the hydrodynamic evolution of the system, we extend our
 113 measurements with Beam Energy Scan phase two (BES-II) data. Another
 114 measurement we perform with the BES-II data is the de-correlation of the

115 event planes with pseudorapidity. The measurement indicates that the third
 116 order event plane decorrelates a lot more (40%) compared to the second or-
 117 der event plane (10%) over one unit of pseudorapidity around midrapidity.
 118 These measurements provide important insights on the longitudinal dynam-
 119 ics and three dimensional modeling of heavy-ion collisions [8].

120 At this Quark Matter, we present the observation of a new phenomenon
 121 which indicates that the triangular flow can drive local polarization of hyper-
 122 ons. Previously, the polarization of hyperons in the longitudinal direction
 123 due to the elliptic flow was observed. The origin of such a phenomenon led
 124 to puzzles in terms of explaining the periodic sign change of polarization
 125 with respect to the reaction plane. Our observation of a similar effect of
 126 local polarization driven by the triangular flow might shed some light on
 127 this puzzle and bring more insights into our understanding of the thermal
 128 vorticity [4].

129 3. Prerequisites for phase transitions and freeze-out

130 In STAR we perform measurements to gain insights on what happens
 131 before and after the QCD phase transitions in the medium formed in rela-
 132 tivistic heavy-ion collisions. The mechanism of baryon stopping provides the
 133 necessary prerequisite for QCD phase transition and enables us to scan the
 134 phase diagram in varying baryon chemical potential. To understand baryon
 135 stopping we perform measurements of proton density with respect to the
 136 center of mass rapidity in Au+Au collisions in fixed target (FXT) mode at
 137 $\sqrt{s_{NN}} = 3$ GeV. An interesting observation is that the shape of the distribu-
 138 tion changes while going from central to peripheral events. Since previous
 139 SPS measurements in this kinematics were performed only in central events,
 140 STAR's measurements provide an opportunity to study the mechanism of
 141 stopping with centrality in the largely unexplored regime of high baryon
 142 density [9].

143 We also study baryon stopping in photonuclear processes in which one
 144 of the colliding object is baryon-free. We use peripheral Au+Au collisions
 145 as a baseline. An interesting observation is that the double ratio of anti-
 146 proton over proton yield in photonuclear over peripheral events is below
 147 unity and has a very strong rapidity dependence. The baryon stopping
 148 measured by the double ratio increases towards the rapidity direction of the
 149 target ion. These results can not be reproduced by PYTHIA simulations
 150 and help gain insights on the microscopic origin of baryon stopping. It
 151 has also the potential to shed light on fundamental questions such as what
 152 exactly carries the baryon number, is it quarks or non-perturbative objects
 153 like baryon junctions [10].

154 Another conserved quantity of importance in the context of QCD tran-

155 sition is strangeness. We perform measurements of the yields of ϕ meson
 156 and compare them to non-resonance particles such as K^- and Ξ in Au+Au
 157 collisions at $\sqrt{s_{NN}} = 3$ GeV. These results can constrain the strangeness
 158 correlation length in a canonical ensemble. In order to understand how
 159 strange hadrons survive freeze out we perform measurement of the yields
 160 of K^{*0} relative to K mesons using the BES-I data. We compare the re-
 161 sults to the ratio of yields for ϕ meson over K as a baseline. When the
 162 yield ratios are plotted against $N_{\text{ch}}^{1/3}$, a proxy for volume, K^{*0}/K decreases
 163 exponentially towards the central event. But for the ϕ/K no such trend
 164 is observed. This observation could be understood by the fact that K^{*0}
 165 may be lost in the medium due to the re-scattering because of its shorter
 166 lifetime, but longer lifetime of ϕ mesons can keep them unaffected. It turns
 167 out these results can be utilized to extract the lower limit of the hadronic
 168 phase lifetime [11].

169 To know more about the late time dynamics of freeze-out and inter-
 170 actions between nucleons and hyperons we perform measurements of the
 171 yields of hyper-nuclei and nuclei. High statistics BES-II data allow pre-
 172 cision hyper-nuclei yield and lifetime measurements. For the first time we
 173 measure the lifetime of hyper-helium-4 in heavy-ion collisions. Our measure-
 174 ments of the relative yields of hyper-triton and hyper-hydrogen-4 indicate
 175 the possible formation of excited hyper-nuclei states in heavy-ion collisions.
 176 At this conference we report the first observation of anti-hyper-hydrogen-4.
 177 This particular hyper-nucleus is made of two anti-neutrons, one proton, and
 178 one anti-lambda; it decays to an anti-helium-4 and a π^+ [12].

179 We perform measurement of the yields of proton and light nuclei in
 180 Au+Au collisions at $\sqrt{s_{NN}} = 3$ GeV. We fit the measurement of p_T spec-
 181 tra using a cylindrical blast-wave model and extract the kinetic freeze-out
 182 parameters such as the effective kinetic freeze-out temperature T_{kin} and the
 183 collective velocity β_T . An important observation is that deuterons freeze
 184 out at a higher effective kinetic temperature than protons [13].

185 To understand more about deuteron production we perform femtoscopic
 186 measurements. We study deuteron-deuteron correlation functions for dif-
 187 ferent centralities, which can be well explained by a coalescence model.
 188 To investigate whether a deuteron is really formed through coalescence of
 189 a proton and a neutron, we perform measurements of the Pearson coeffi-
 190 cients between the number of protons and deuterons as a function of the
 191 collision energy. We see a negative Pearson coefficient. The measurement
 192 is explained by models that include baryon number conservation and co-
 193 alescence. A general conclusion is that a coalescence between a proton
 194 and a neutron provides a consistent explanation of deuteron formation at
 195 RHIC [14, 15].

196 4. Critical phenomena and mapping of QCD phase diagram

197 We continue the search for the QCD critical point (CP) by studying
 198 the net-proton higher order cumulants as a function of collision energy.
 199 The measurements with the BES-I data have established a non-monotonic
 200 trend of kurtosis times variance with collision energy. The most recent
 201 addition to extend the CP search is the measurements in Au+Au collisions
 202 at $\sqrt{s_{NN}} = 3$ GeV. The results of net-proton kurtosis times variance at
 203 this energy indicate that the measurement is dominated by baryon number
 204 conservation. We also perform the measurements of the higher moments of
 205 deuteron number fluctuations as a function of collision energy. The kurtosis
 206 times variance for deuterons is found to be below unity over the energy
 207 range of $\sqrt{s_{NN}} = 7.7 - 200$ GeV, but no non-monotonicity is observed. The
 208 outstanding question is why there is a difference between the measurements
 209 of proton and deuteron fluctuations. In this context, the connection to the
 210 smaller yields of deuterons and different freeze-out parameters that were
 211 seen from the other measurements are being investigated [15,16].

212 Another topic of prime interest is the search for chiral crossover that
 213 is predicted to happen at low baryon chemical potential (μ_B). For this
 214 we perform measurements of fifth (C_5) and sixth (C_6) order cumulants of
 215 net-protons and the ratios such as C_5/C_1 and C_6/C_2 at the top RHIC
 216 energy. Our measurements performed as a function of hadron multiplicity
 217 at mid-rapidity in $p + p$, isobar, and Au+Au collisions show that these
 218 ratios decrease with increasing multiplicity and eventually approach the
 219 predictions from lattice QCD, which also predicts a smooth crossover at
 220 $\mu_B = 0$ [17].

221 Another measurement from STAR that is compared to lattice QCD pre-
 222 dictions uses the dilepton as a thermometer of the medium. Using the data
 223 from Au+Au collisions at $\sqrt{s_{NN}} = 27$ and 54.4 GeV we measure the excess
 224 yield of dileptons over the cocktail as a function of the invariant mass. We fit
 225 the data at the low-mass ($m_{ll} < 1$ GeV) and intermediate-mass ($1 < m_{ll} < 3$
 226 GeV) regions to extract the medium temperature. When shown on a tem-
 227 perature (T) versus μ_B plot, the values of effective temperature extracted
 228 by the low-mass region T_{LMR} are very close to the chiral crossover band that
 229 is predicted by the lattice QCD. This is indicative of ρ mediated dilepton
 230 emissions dominating near the chiral crossover transition. The extracted
 231 T_{IMR} values (≈ 300 MeV) are much higher than T_{LMR} . This indicates the
 232 intermediate-mass dilepton spectrum probes the temperature of the QGP
 233 medium. This is the first blue-shift free measurement of QGP temperature
 234 at RHIC.

5. Hard probes

235

236 Another well known indicator of the QGP temperature, categorized as
 237 a hard probe, is quarkonium. In this context, we perform measurements of
 238 nuclear modification factor R_{AA} of J/ψ in isobar collisions at $\sqrt{s_{NN}} = 200$
 239 GeV, and also in Au+Au collisions at $\sqrt{s_{NN}} = 54.4$ GeV. At a given number
 240 of participants (N_{part}), no significant difference is observed between different
 241 systems and energies. The values of R_{AA} at both energies, in central events,
 242 are significantly below unity, unlike what was observed at the LHC. This
 243 indicates that the RHIC measurements are consistent with the dominance of
 244 J/ψ dissociation. We perform the measurement of R_{AA} of charged hadrons
 245 at high $p_T > 5.1$ GeV/ c in isobar collisions, which is also considered as a
 246 hard probe. We find a suppression in central events driven by mechanism
 247 that also leads to the phenomenon of jet quenching. In peripheral events,
 248 our measurements are affected by centrality bias which is investigated using
 249 PYTHIA combined with Monte-Carlo Glauber model simulations [18].

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Using the Heavy Flavor Tracker of STAR, we have performed the first
 251 measurements of D^0 tagged jets at RHIC. We reconstruct jets that con-
 252 tain a D^0 with $p_{T,D^0} > 5$ GeV/ c with resolution parameter of $R = 0.4$ in
 253 Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. Our measurements are unfolded
 254 using response matrix obtained from the PYTHIA-8 Detroit tune and the
 255 STAR GEANT simulation. We measure the nuclear modification factor (in
 256 reference to peripheral events) R_{CP} in central and mid-central events with
 257 the transverse momentum of the reconstructed jets $p_{T,\text{jet}}$. We see that R_{CP}
 258 increases against $p_{T,\text{jet}}$ to approach unity around $p_{T,\text{jet}} = 12$ GeV/ c . To
 259 study the radial profile of D^0 , we vary “r”, the distance of the D^0 from the
 260 jet axis in the range of 0-0.2. For a given “r” we study the ratio of the
 261 yields between central or mid-central events to peripheral events. The ratio
 262 is found to be consistent with unity indicating no modification of D^0 radial
 263 profile with centrality, within measurement uncertainties. These measure-
 264 ments can constrain theoretical models on heavy quark diffusion and energy
 265 loss at RHIC [19].

266

To understand how the vacuum parton shower in $p+p$ collisions gets
 267 modified in heavy-ion collisions due to in-medium gluon radiations, we per-
 268 form measurements of the ratio of the recoil jet yield as a function of jet p_T
 269 with $R = 0.2$ over $R = 0.5$ in $p+p$, and Au+Au collisions at $\sqrt{s_{NN}} = 200$
 270 GeV. We use π^0 and direct photon (γ^{dir}) triggered jets. We observe that the
 271 ratios in Au+Au collisions are significantly lower than that in $p+p$ indicat-
 272 ing medium-induced broadening of jet showers at RHIC. We also perform
 273 the first measurements of acoplanarity for both π^0 and γ -triggered jets. We
 274 observe medium-induced jet acoplanarity in heavy-ion collisions compared
 275 to the PYTHIA baseline for larger values of jet radius [20].

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6. Upgrades and future program

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In 2021, STAR has successfully installed the forward upgrades. Three different subsystems, the forward silicon tracker, the small strip thin gap chamber, and the forward calorimeter that includes electromagnetic and hadronic layers, have been fully installed and participated in data taking in $p+p$ collisions at 510 GeV. These systems are installed in one of the forward directions at STAR, and will allow important measurements during the anticipated Au+Au runs in 2023 and 2025, and in polarized $p+p$ and $p+Au$ runs in 2024 [21]. An important plan is to perform measurements that can be repeated at the Electron-Ion Collider (EIC). In this direction we perform an exploratory measurement of di-hadron correlations in photonuclear events to search for signatures of collectivity. Although no such signature is observed now, this will be revisited in future measurements. STAR forward upgrade program will open paths to study the microstructure of the QGP and enable measurements informative towards EIC science.

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