Anisotropic Flow of Identified Particles in Au + Au Collisions at $\sqrt{s_{NN}}$ = 3 - 3.9 GeV at RHIC

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5	Abstract. In these proceedings, we present transverse momentum dependence
6	of the mid-rapidity slope of directed flow $(dv_1/dy _{y=0})$ for π^+ and K_S^0 in Au + Au
7	collisions at $\sqrt{s_{NN}} = 3.0, 3.2, 3.5, \text{ and } 3.9 \text{ GeV}$. Both π^+ and K_S^0 show negative
8	v_1 slope at low p_T ($p_T < 0.6 \text{ GeV}/c$). Collision energy dependence of v_1 slope
9	and p_T -integrated v_2 for π^{\pm} , K_S^0 , and Λ are also presented. A comparison to JAM
10	model calculations indicates that spectator shadowing can lead to anti-flow at
11	low p_T . In addition, a breaking of the Number of Constitute Quark (NCQ)
12	scaling of elliptic flow (v_2) is observed at $\sqrt{s_{NN}} = 3.2$ GeV, which implies the
13	dominance of hadronic degrees of freedom occurs in collisions at $\sqrt{s_{NN}} = 3.2$
14	GeV and below.

15 1 Introduction

The goals of Beam Energy Scan (BES) program at Relativistic Heavy Ion Collider (RHIC) are 16 searching for the possible QCD critical point and locating the first order phase boundary [1]. 17 The energy dependence of net-proton v_1 slope [2] shows possible minimum at $\sqrt{s_{NN}} \approx 10$ -18 20 GeV, implies that the softest point of Equation of State (EoS) may exist within this range 19 of collision energy. The existence of partonic collectivity is observed through NCQ scaling 20 of v_2 at higher BES energies ($\sqrt{s_{NN}} > 7.7$ GeV) [3], while the break NCQ scaling of v_2 at 21 $\sqrt{s_{NN}}$ = 3.0 GeV [4] indicates the partonic collectivity is disappeared at this energy. In this 22 contribution, we present the most recent measurements of directed flow (v_1) and elliptic flow 23 (v₂) of identified particles (π^{\pm} , K^{\pm} , K^0_S , p, and A) at $\sqrt{s_{NN}} = 3.0 - 3.9$ GeV in Fixed Target 24 Au + Au collisions. 25

26 2 Experiment Setup

For identification of π^{\pm} , K^{\pm} , protons and anti-protons, a combination of Time Projection 27 Chamber (TPC) [5] and Time of Flight (TOF) [6] is used. The left panel of Figure 1 illustrates 28 the rigidity (p/q): particle momentum divided by charge) dependence of ionization energy 29 loss (dE/dx) in the TPC. The dashed line represents the theoretical ionization energy loss 30 curve for particle passing through the TPC. Particle identification by TOF is based on particle 31 mass square (m^2) distribution, which can be obtained from particle velocity (β). Moreover, 32 the Kalman Filter (KF) particle package [7], where the covariance matrix of reconstructed 33 tracks is taken into account, is employed to reconstruct weak decay particles (K_s^0 and Λ). An 34 example of reconstructing the invariant mass using the KF particle package is demonstrated 35 with the K_S^0 meson in the right panel of Figure 1. 36

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Figure 1. Left: Rigidity dependence of particle ionization energy loss in TPC. Right: Invariant mass distribution of K_S^0 in Au + Au collisions at $\sqrt{s_{NN}} = 3.5$ GeV.

37 3 Results

38 3.1 Anti-flow of Kaon

The anti-flow of kaon was first observed by E895 Collaboration at 6 A GeV [8]. It was at-39 tributed to the repulsive potential associated with the strange quark in K_{s}^{0} . We have observed 40 anti-flow behavior in kaons and pions for $p_T < 0.6 \text{ GeV}/c$ in mid-central Au + Au collisions 41 at $\sqrt{s_{NN}}$ = 3.0, 3.2, 3.5, and 3.9 GeV using the fixed target data from STAR. Figure 2 shows 42 transverse momentum (p_T) dependence of v_1 slope $(dv_1/dy|_{y=0})$ for π^+ and K_s^0 from STAR. 43 The hadronic transport model JAM [9] calculations are compared with experimental data at 44 3.9 GeV. The JAM model in hadronic cascade mode (blue band) can successfully capture the 45 anti-flow pattern at low p_T for π^+ and K_S^0 , even without the inclusion of a kaon potential [8]. 46 However, the JAM model with baryonic mean field (red band), tends to overestimate the v_1 47 slope for π^+ and K_s^0 . Additionally, the JAM mean field without spectator contribution (black 48 band) exhibits a larger v_1 slope compared to the one with spectators. The data-model com-49 parisons suggest that the shadowing effect [10] from the spectator may also play a significant 50

⁵¹ role in generating anti-flow at low p_T .



Figure 2. v_1 slope of π^+ (left) and K_S^0 (right) as function of transverse momentum and a comparison with JAM calculation at $\sqrt{s_{NN}} = 3.9$ GeV.

52 3.2 NCQ Scaling of v_2

The left and right panels of Figure 3 illustrate the number of constituent quarks (n_a) scaled 53 elliptic flow (v_2/n_q) as a function of transverse kinetic energy $((m_T - m_0)/n_q)$ for particles 54 $(\pi^+, K^+, K_S^0, p, \text{ and } \Lambda)$ and the corresponding anti-particles $(\pi^-, K^-, \text{ and } K_S^0)$, respectively 55 for Au + Au collisions at $\sqrt{s_{NN}}$ = 3.2 GeV. The NCQ scaling of v_2 is broken completely 56 for particles and anti-particles at $\sqrt{s_{NN}} = 3.2$ GeV. The existence of partonic collectivity is 57 observed through NCQ scaling of v_2 at higher BES energies ($\sqrt{s_{NN}} > 7.7$ GeV) [3]. The 58 disappearing of NCQ scaling in v_2 at $\sqrt{s_{NN}} = 3.2$ GeV implies that hadronic interactions 59 play an important role at this energy and below [4, 11]. 60



Figure 3. Number of Constituent Quark scaling of v_2 as a function of the scaled transverse kinetic energy for particles (left) and the corresponding anti-particles (right) at $\sqrt{s_{NN}} = 3.2$ GeV.

⁶¹ 3.3 Energy Dependence of v_1 and v_2

The top panel of Figure 4 shows collision energy dependence of v_1 slope in 10-40% mid-62 central Au + Au collisions at $\sqrt{s_{NN}} = 3.0 - 3.9$ GeV. The v_1 slopes of π^+ (solid square) are 63 negative at 3.0 - 3.9 GeV, while v_1 slopes of π^- (open square) are positive. The difference 64 between π^+ and π^- may be explained by Coulomb effect [12]. Furthermore, v_1 slopes of K_S^0 65 (solid triangle) are greater than π^+ , v_1 slopes of Λ (solid circle) are largest among these four 66 particle species. The v_1 slopes of all particles (π^{\pm}, K_S^0 , and Λ) decrease in magnitude as colli-67 sion energy increases. The lower panel in Figure 4 depicts the collision energy dependence of 68 transverse momentum (p_T) integrated v_2 . It is observed that the sign of v_2 changes from neg-69 ative to positive for all particles (π^{\pm} , K_{S}^{0} , and Λ) within the collision energy range of $\sqrt{s_{NN}}$ = 70 3.0 - 3.9 GeV. This shift signifies the transition from out-of-plane to in-plane expansion [13], 71 occurring specifically within the mentioned energy range. 72

The JAM calculations for Λ are represented by colored bands, with blue, red, and black bands corresponding to cascade, baryonic mean field, and mean field without spectators modes, respectively. The comparison of Λv_1 between the data and model calculations suggests the presence of a strong baryon mean field [4] in the high baryon density region. Furthermore, the comparison between the measured p_T -integrated v_2 and model calculations indicates the significant influence of spectator shadowing in the energy range of $\sqrt{s_{NN}} = 3.0$ - 3.9 GeV.

Summary

In summary, we present directed flow (v_1) and elliptic flow (v_2) measurements for identified 81 particles $(\pi^{\pm}, K^{\pm}, K^0_S, p, \text{ and } \Lambda)$ in Au + Au collisions at $\sqrt{s_{NN}} = 3.0, 3.2, 3.5, \text{ and } 3.9 \text{ GeV}$. 82 The measurements for π^+ and K_S^0 show negative v_1 slope $(dv_1/dy|_{y=0})$ at low p_T ($p_T < 0.6$ 83 GeV/c). The transport model JAM reproduces anti-flow at low p_T without incorporating kaon 84 potential, and indicates shadowing effect from spectator can lead to anti-flow. Secondly, NCQ 85 scaling of v_2 is broken completely for particles $(\pi^+, K^+, K^0_S, p, and \Lambda)$ and anti-particles $(\pi^-, K^+, K^0_S, p, and \Lambda)$ 86 K^{-}) at $\sqrt{s_{NN}} = 3.2$ GeV, implying that the hadronic interactions are dominant at $\sqrt{s_{NN}} =$ 87 3.2 GeV and below. At last, collision energy dependence of v_1 slope $(dv_1/dy|_{u=0})$ and p_T -88 integrated v_2 at $\sqrt{s_{NN}} = 3.0 - 3.9$ GeV are presented. The v_1 slopes of all particles (π^{\pm} , 89 $K_{\rm S}^0$, and Λ) decrease in magnitude as collision energy increases. And the sign change in v_2 90 indicates that the change of out-of-plane to in-plane expansion happens between $\sqrt{s_{NN}} = 3.0$ 91 - 3.9 GeV. 92

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Figure 4. v_1 slope (top) and p_T -integrated v_2 (bottom) as a function of collision energy and compared with JAM calculation for Λ . Note that p_T windows for π^{\pm} , K_s^0 , and Λ are $0.2 < p_T < 1.6$ GeV/*c*, $0.4 < p_T < 1.6$ GeV/*c*, and $0.4 < p_T < 2.0$ GeV/*c*, respectively. And the rapidity window is -0.5 < y < 0 for p_T -integrated v_2 .

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