Supported in part by





# **Measurement of** $\pi^+\pi^-$ **Azimuthal Correlation** Asymmetry Using $p^{\uparrow}p$ Data At $\sqrt{s} = 200$ GeV At STAR

**Babu Pokhrel** for the STAR collaboration **June 7 - 10, 2022** 

2022 RHIC/AGS ANNUAL USERS' MEETING

From RHIC to EIC At the QCD Frontiers





#### College of Science and Technology





#### **Motivation:** Transversity $(h_1(x))$

+ Transversity,  $h_1(x)$ , is a leading order parton distribution  $\overrightarrow{P}$  = Nucleon polarization function (PDF). It describes transverse polarization of quark  $\vec{p}$  = Nucleon momentum in transversely polarized nucleon, which is least known from  $h_1(x) \approx$ = Quark polarization experiments.

- + It is chiral-odd. It can be accessed only by coupling with another chiral-odd function, such as Interference Fragmentation Function, IFF  $(H_1^{\triangleleft})$ .
- + In polarized proton-proton collisions  $(p^{\uparrow}p)$ , the  $\pi^+\pi^-$  azimuthal correlation asymmetry,  $A_{UT}$ , gives rise to the sensitivity to  $h_1(x)$  coupled with  $H_1^{\triangleleft}$ .

 $\blacklozenge$  Nucleon tensor charge,  $g_T$ , is essential to characterize the nucleon spin structure.



 $g_T = \delta u - \delta d,$  $dx(h_1^u(x) - h_1^{\bar{u}}(x)),$  $\delta u =$  $dx(h_1^d(x) - h_1^{\bar{d}}(x))$  $\delta d =$ 

Babu Pokhrel, 2022 RHIC AGS ANNUAL USERS' MEETING





### **STAR Experiment, Datasets, And Kinematics**



✦ This analysis uses 2015 data, which provides the most precise  $A_{UT}$ measurement in the range 0.1 < x < 0.4.

Collision	
Polarization	
Year	2006
$\sqrt{s}$ (GeV)	200
Lumi. (pb^-1)	~1.8
Avg. Polz.(%)	~60

 $\bullet$  STAR covers much higher  $Q^2$  than HERMES and COMPASS and probes  $h_1(x)$  in the valence quark region (0.1 < x < 0.4).



• Relativistic Heavy Ion Collider is the first polarized proton-proton collider in the world. + It is capable of colliding polarized proton beams up to a center-of-mass energy,  $\sqrt{s}$ , of 510 GeV. ✦ STAR's particle identification relies on the measured ionization energy loss (dE/dx) by the Time Projection Chamber (TPC).





- The PID capability based on the TPC is limited for p > 1 GeV/c, resulting in large systematic uncertainties. Time of Flight (TOF) detector helps
- to improve PID for p > 1 GeV/c, in conjunction with the TPC.

Х





### **Extraction of** $\pi^+\pi^-$ **Azimuthal Correlation Asymmetry (** $A_{UT}$ **)**

**+ Reaction channel**:  $p^{\uparrow} + p \rightarrow h^+h^- + X$ 

**+** Detailed subprocess and the final state:



Di-hadron azimuthal correlation asymmetry:

$$A_{UT} = \frac{d\sigma^{\uparrow} - d\sigma^{\downarrow}}{d\sigma^{\uparrow} + d\sigma^{\downarrow}} \propto h_1 H_1^{\triangleleft}$$

Bacchetta et.al. (Phys.Rev.D 70 (2004) 094032)



modulation of di-hadron azimuthal angle

$$\phi_{RS}(=\phi_s-\phi_R).$$

 $A_{UT}$  can be extracted using cross-ratio formula,



$$A_{UT} \cdot sin(\phi_{RS}) = \frac{1}{P} \frac{\sqrt{N^{\uparrow}(\phi_{RS})N^{\downarrow}(\phi_{RS} + \pi)} - \sqrt{N^{\downarrow}(\phi_{RS})N^{\uparrow}(\phi_{RS} + \pi)}}{\sqrt{N^{\uparrow}(\phi_{RS})N^{\downarrow}(\phi_{RS} + \pi)} + \sqrt{N^{\downarrow}(\phi_{RS})N^{\uparrow}(\phi_{RS} + \pi)}}$$

where,  $N^{\uparrow(\downarrow)}$  is the number of  $\pi^+\pi^-$  in respective  $\phi_{RS}$  bin when the beam polarization is  $\uparrow (\downarrow)$ .

 $\blacklozenge$  It is free from detector effects, which leads to the reduced systematic uncertainties.







## **Preliminary Results**

**Fig. 5(a):**  $A_{UT}$  (top panel) and fractional momentum of proton carried by scattered quark (x) and fractional energy of quark carried by outgoing  $\pi^+\pi^-(z)$ A<sup>Sin(</sup>UT (bottom panel) vs  $\pi^+\pi^-$  pseudorapidity ( $\eta^{\pi^+\pi^-}$ ). ▶ x and z are estimated from simulation. x increases linearly from ~ 0.1 to 0.22, but z doesn't show a clear dependence (  $\langle z \rangle \sim 0.46$ ) over the measured  $\eta^{\pi^+\pi^-}$  range. A<sub>UT</sub> increases linearly in  $\eta^{\pi^+\pi^-} > 0$  region, corresponding to high-x, where  $h_1(x)$  is expected to be sizeable.



**Fig. 5(b)i:**  $A_{UT}$  vs di-pion invariant mass  $(M_{inv}^{\pi^+\pi^-})$  in  $\eta^{\pi^+\pi^-} > 0$  region integrated over the di-pion transverse momentum  $(p_T^{\pi^+\pi^-})$  compared with the theory calculation and the previous **STAR** STAR result.

**Fig. 5(b)ii:**  $A_{UT}$  vs  $M_{inv}^{\pi^+\pi^-}$  in different  $p_T^{\pi^+\pi^-}$  bins. A resonance peak at  $M_{inv}^{\pi^+\pi^-} (\approx M_{\rho}) \sim 0.8 \ GeV/c^2$ is observed, which is expected due to the interference between the  $\pi^+\pi^-$  produced from different production channel.

0.03

0.02

0.01

•  $A_{UT}$  evolves with the  $p_T^{\pi^+\pi^-}$  with the prominent resonance peak at the highest  $p_T^{\pi^+\pi^-}$  bin in  $\eta^{\pi^+\pi^-} > 0$ . In  $\eta^{\pi^+\pi^-} < 0$ ,  $A_{IIT}$ is small as expected.

The systematic uncertainty includes those from the trigger bias and PID with the dominance from the PID. The PID systematic effect is well understood, and expected to be reduced by more than 50% after including TOF.









