

Hard Probes 2013, Cape Town

Inclusive Spectrum of Fully Reconstructed Charged Jets in Central Au+Au Collisions at $\sqrt{s_{NN}}=200$ GeV by the STAR Collaboration



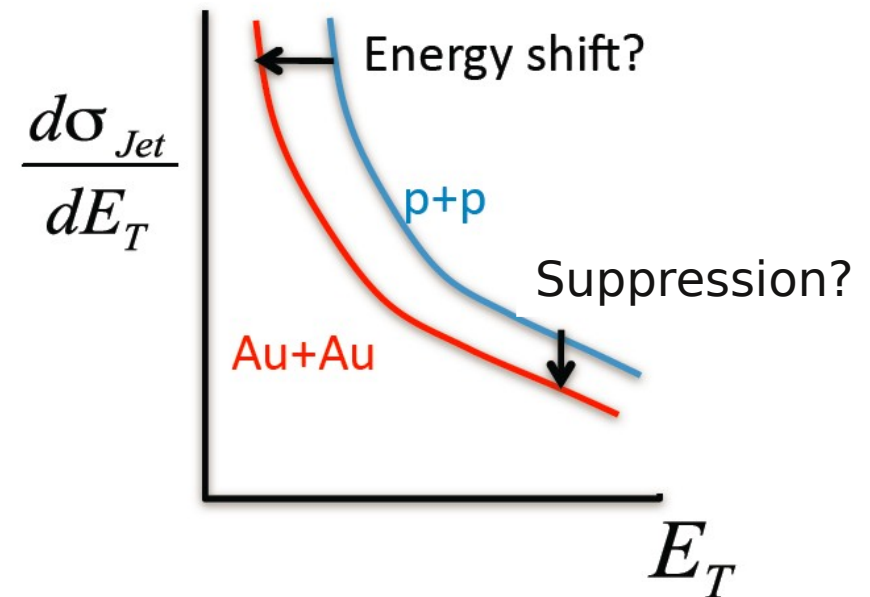
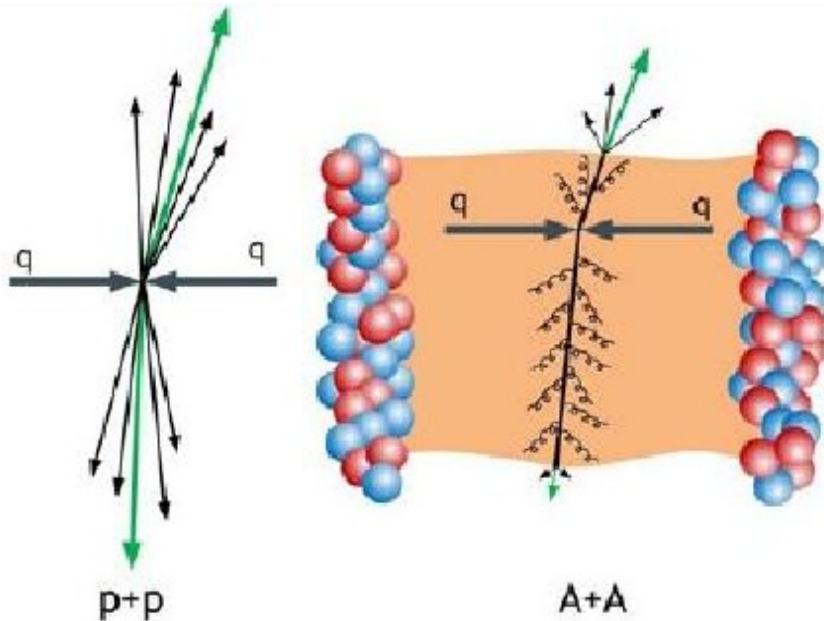
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(for the STAR Collaboration)

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Motivation: jets as a probe of nuclear matter

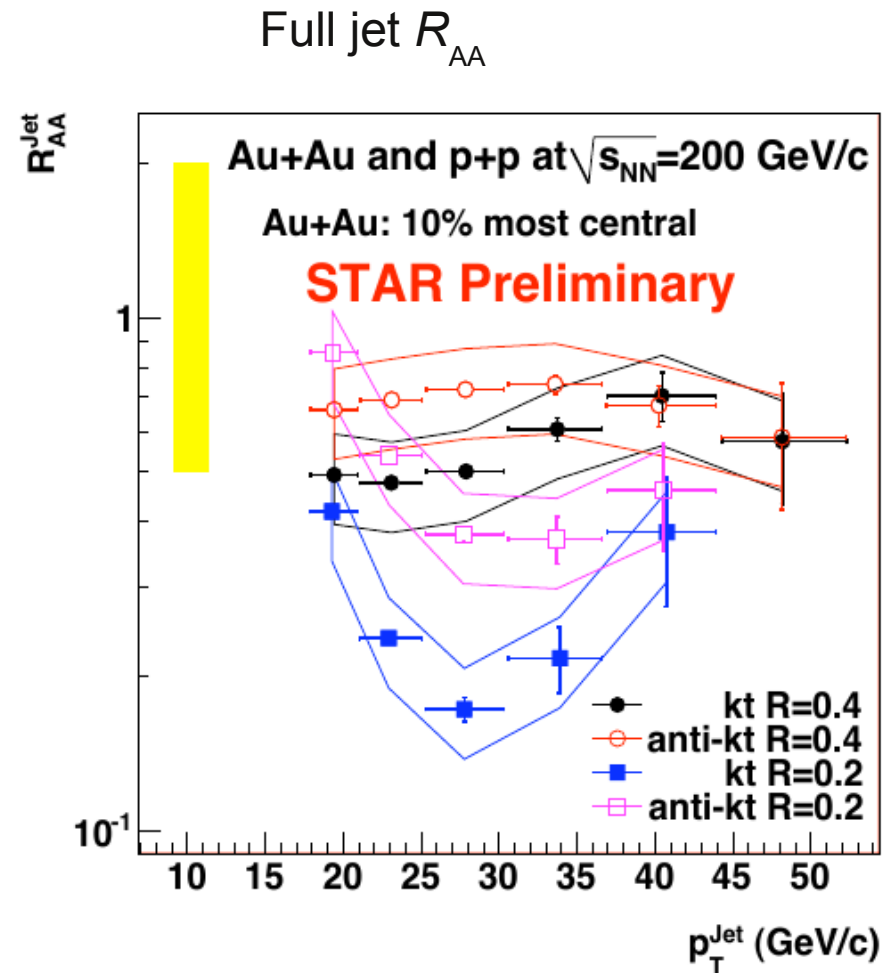
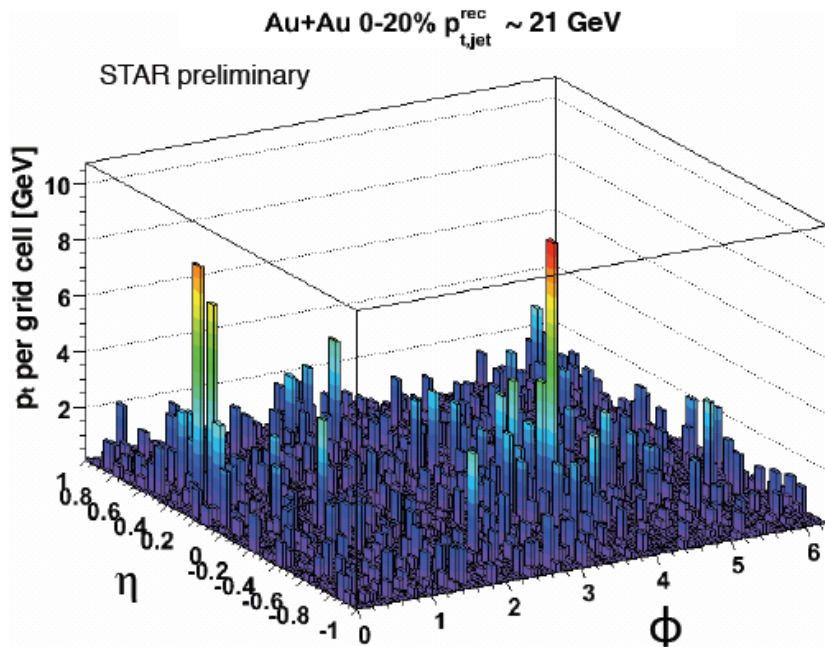
- jets are collimated sprays of hadrons created by fragmentation and hadronization of hard-scattered partons
- in heavy-ion collisions: probe of hot and dense nuclear matter
jet quenching – modification of fragmentation



However: Jet reconstruction is an extremely challenging task in the high multiplicity environment and large and fluctuating background ...

Full jet reconstruction in Au+Au collisions at STAR

Early Quark Matter '09 results on Run 7 Au+Au data at 200 GeV
- limited statistics, new methods developed since then



M. Ploskon, Nucl.Phys.A830:255c-258c,2009

Jet tomography in A+A collisions: requirements

- well defined and transparent selection of jet population
 - *understand what biases we impose*
- direct comparison to theory
 - *jet distributions corrected for background and instrumental effects to particle level*
- same approach and algorithms at RHIC and LHC over the full jet kinematic range: $p_T^{\text{jet}} > \sim 15 \text{ GeV}/c$
 - *to achieve a well controlled energy evolution of quenching*
- collinear safety – as close as possible, low infrared (IR) cutoff:
 - *no background suppression via cuts on jet constituents*

Caveat! Detector response also imposes infrared cutoffs



Jet Reconstruction

reconstructed jet (jet candidate) = output of the jet reconstruction algorithm

FASTJET [Eur.Phys.J. C72 (2012) 1896]

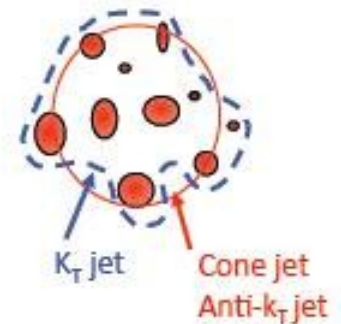
- implements fast IR and collinear safe algorithms
- sequential clustering algorithms: kT, anti-kT
 - cluster hadrons into jet objects over a range specified by parameter R
 - in our analysis: $R=0.2, 0.3, 0.4$

Clustering:

$$d_{ij} = \min(p_{T,i}^a, p_{T,j}^a) \frac{(\Delta\varphi)^2 + (\Delta\eta)^2}{R^2} \quad \left. \vphantom{d_{ij}} \right\} \begin{array}{l} \rightarrow \min(d_{iB}, d_{ij}) \\ \rightarrow \text{add to jet} \\ \rightarrow \text{final jet} \end{array}$$

$d_{iB} = \frac{p_{T,i}^a}{R}$

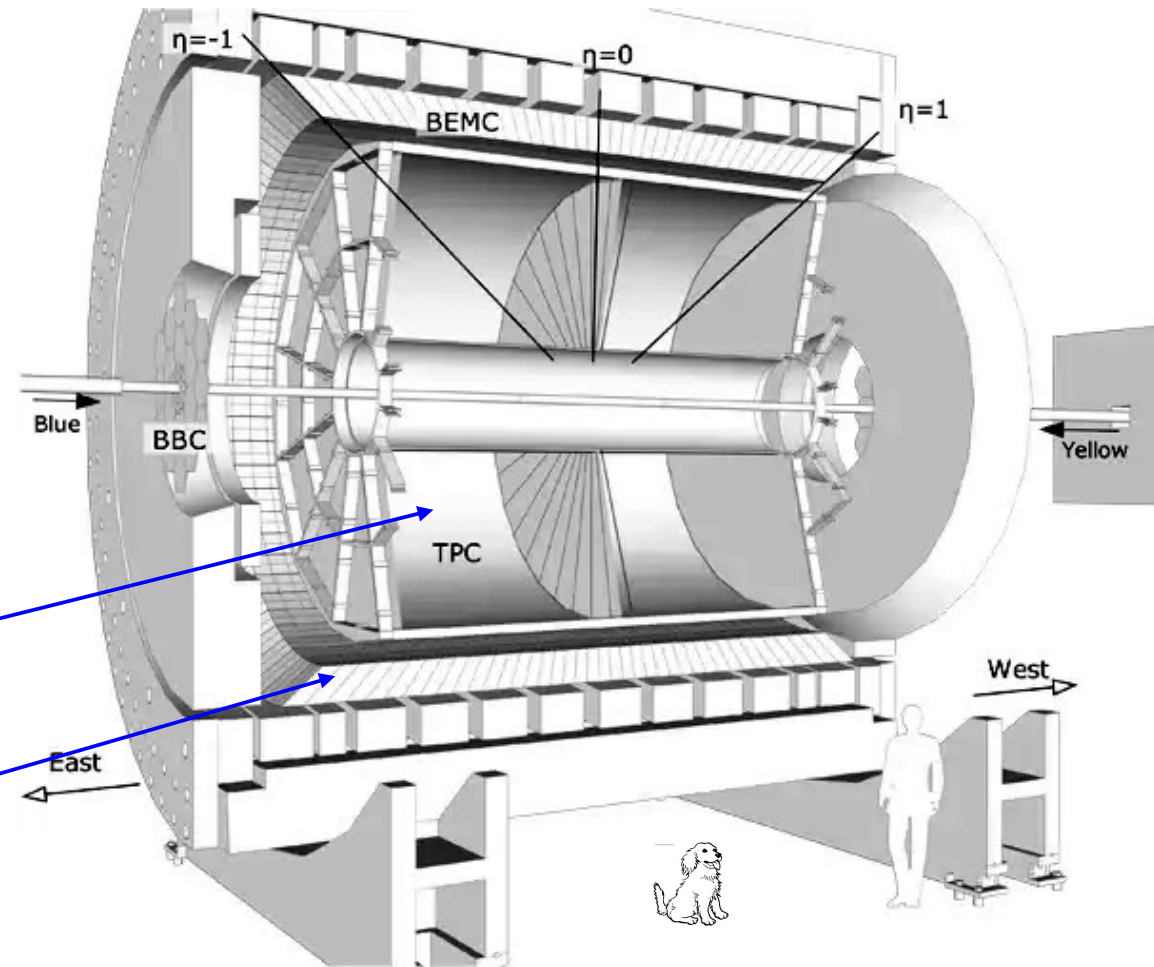
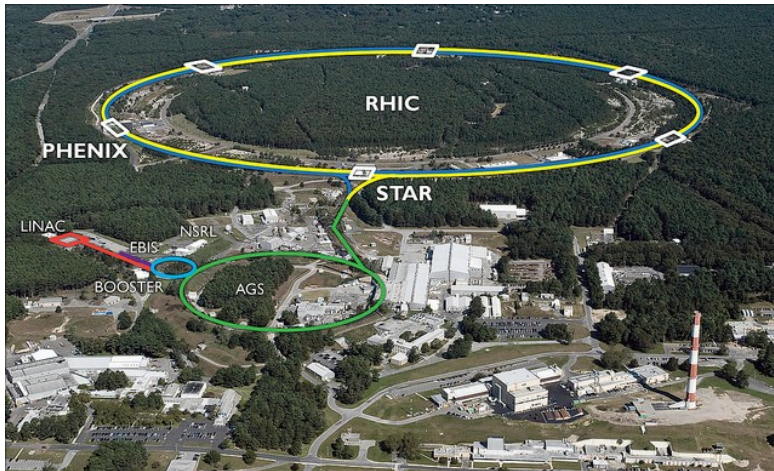
Where $a = 2$ for kT and -2 for anti-kT



STAR Detector

Relativistic Heavy Ion Collider (RHIC)

Solenoidal Tracker at RHIC (STAR)



Time Projection Chamber

Barrel Electromagnetic Calorimeter

full azimuthal coverage
pseudorapidity coverage: $-1 < \eta < 1$

Data Sample and Charged Jet Analysis

Data set:

- RHIC Run11 data
- Au+Au @ $\sqrt{s_{NN}}=200$ GeV
- 10% most central Au+Au collisions from minimum bias (MB) trigger
- Integrated luminosity $\sim 60\mu\text{b}^{-1}$ (15% of total production of Run11 data)

Event selection:

- $|z(\text{vertex})| < 30\text{cm}$

Charged track selection:

- reconstructed in TPC
- DCA $< 1\text{cm}$
- Number of fit hits > 20
- # of fit hits/# of max. available > 0.55

Jet reconstruction:

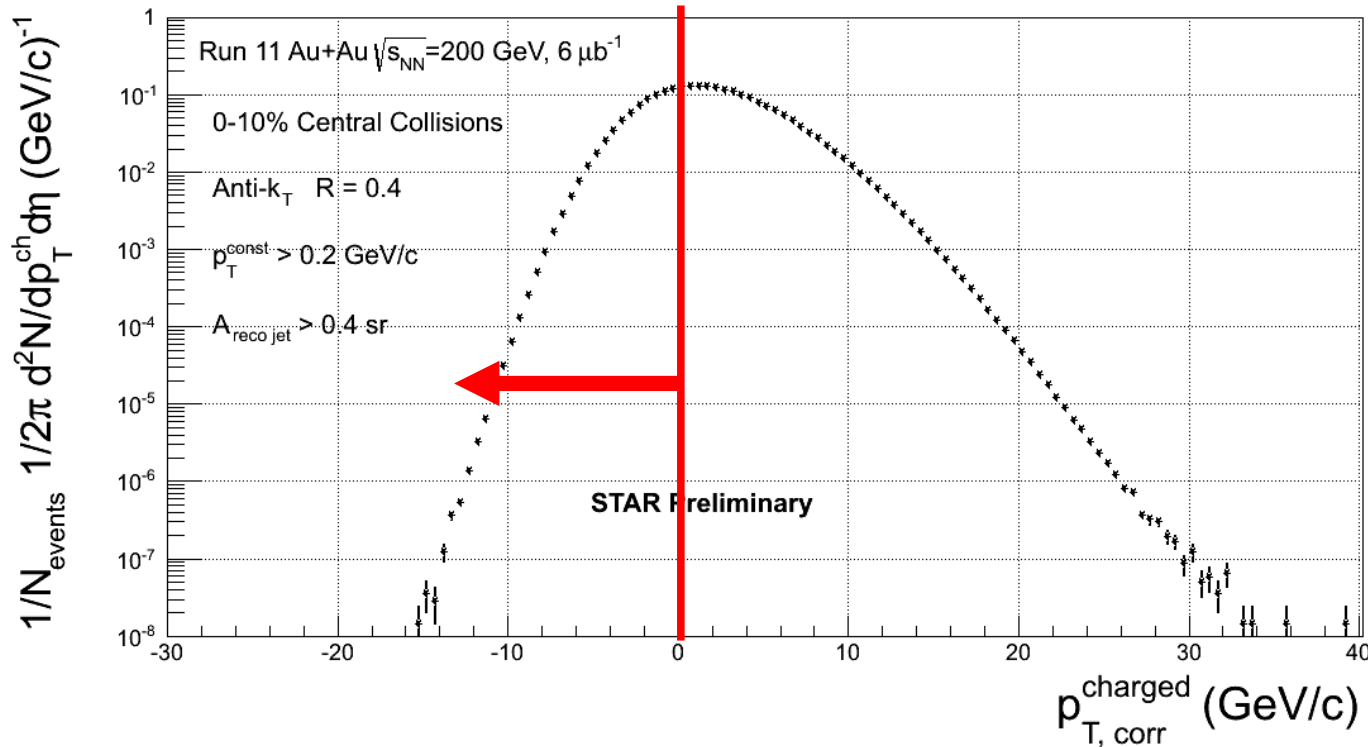
- Charged jets
- anti- k_T jet reconstruction algorithm
resolution parameter: $R=0.2, 0.3, 0.4$
- very low IR cutoff:
 $p_T(\text{constituent}) > 0.2 \text{ GeV}/c$
- jet active area: $A > 0.4 \text{ Sr}$ ($R=0.4$)
 $A > 0.2 \text{ Sr}$ ($R=0.3$)
 $A > 0.09 \text{ Sr}$ ($R=0.2$)
- fiducial jet acceptance: $|\eta| < 1-R$

Heavy Ion Jet Reconstruction

- take all jets in acceptance
- Jet candidates: reconstructed using anti- k_T algorithm
- jet area A : Fastjet definition
- bckd. energy density – calculated event-wise (kT):
(hard jets not discarded for the calculation)
- distribution corrected for bckd. energy density:

$$\rho = \text{med} \left\{ \frac{p_{T,i}}{A_i} \right\}$$

$$p_{T,corr} = p_T - A_{jet} \times \rho$$



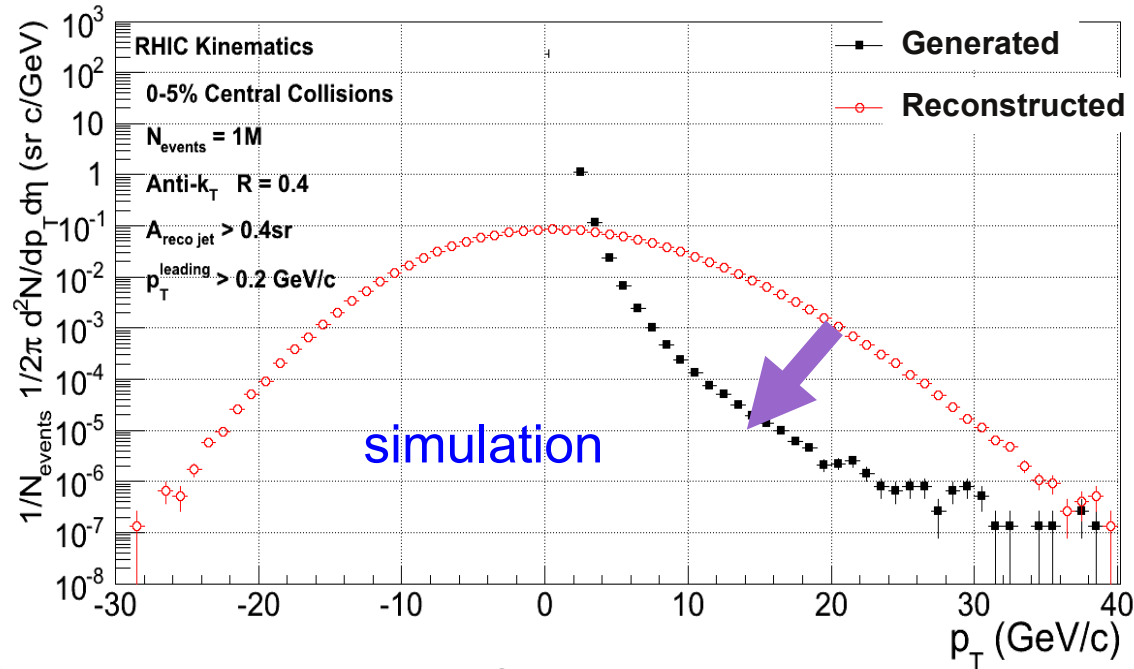
~half of the jet candidates have negative p_T^{corr}

we don't discard them (for now)

contain crucial information about background

Unfolding of Measured Spectra

- Undo the effects of smearing on hard jet spectrum
- In order to compare measured data directly with theory
- Correction for BG fluctuations
- Correction for detector effects



- “Inversion” of response matrix => unfolding matrix
- We use iterative method based on Bayes' theorem [[G. D'Agostini, arXiv:1010.0632](#)]
- Singular Value Decomposition (SVD) unfolding used to validate Bayesian for background fluctuations. Full unfolding with SVD in progress [[Nucl.Inst.Meth.A372:469-481,1996](#)]

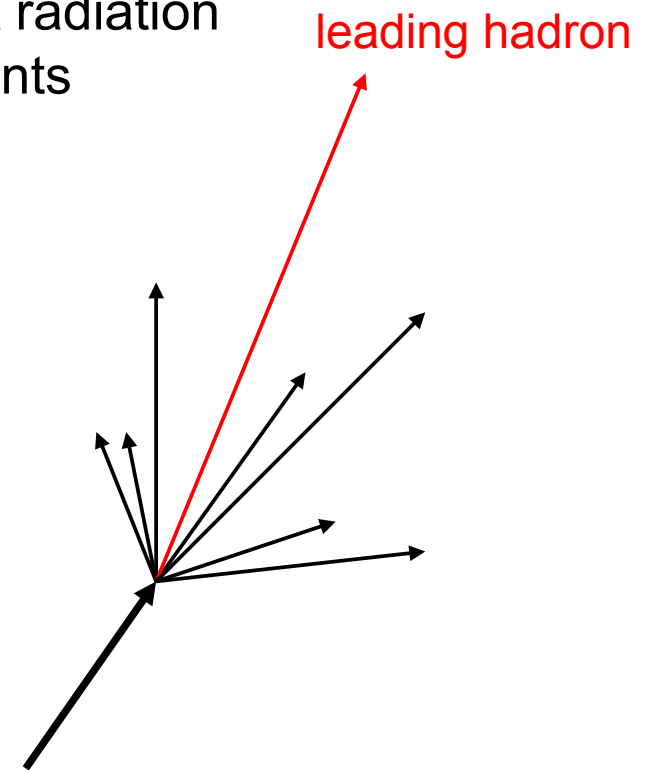
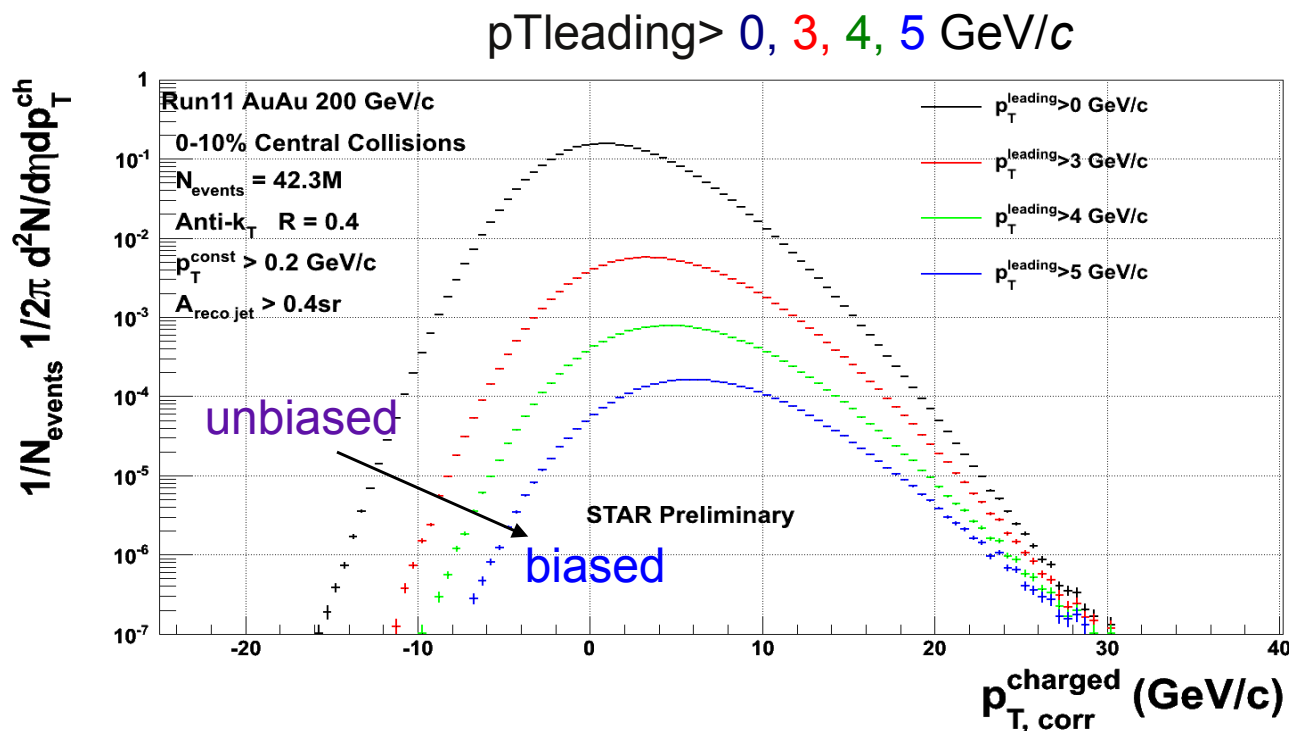
Toward the Inclusive Jet Spectrum

- Stable unfolding of inclusive jet spectrum:

need to **reduce the combinatorial background prior to unfolding**

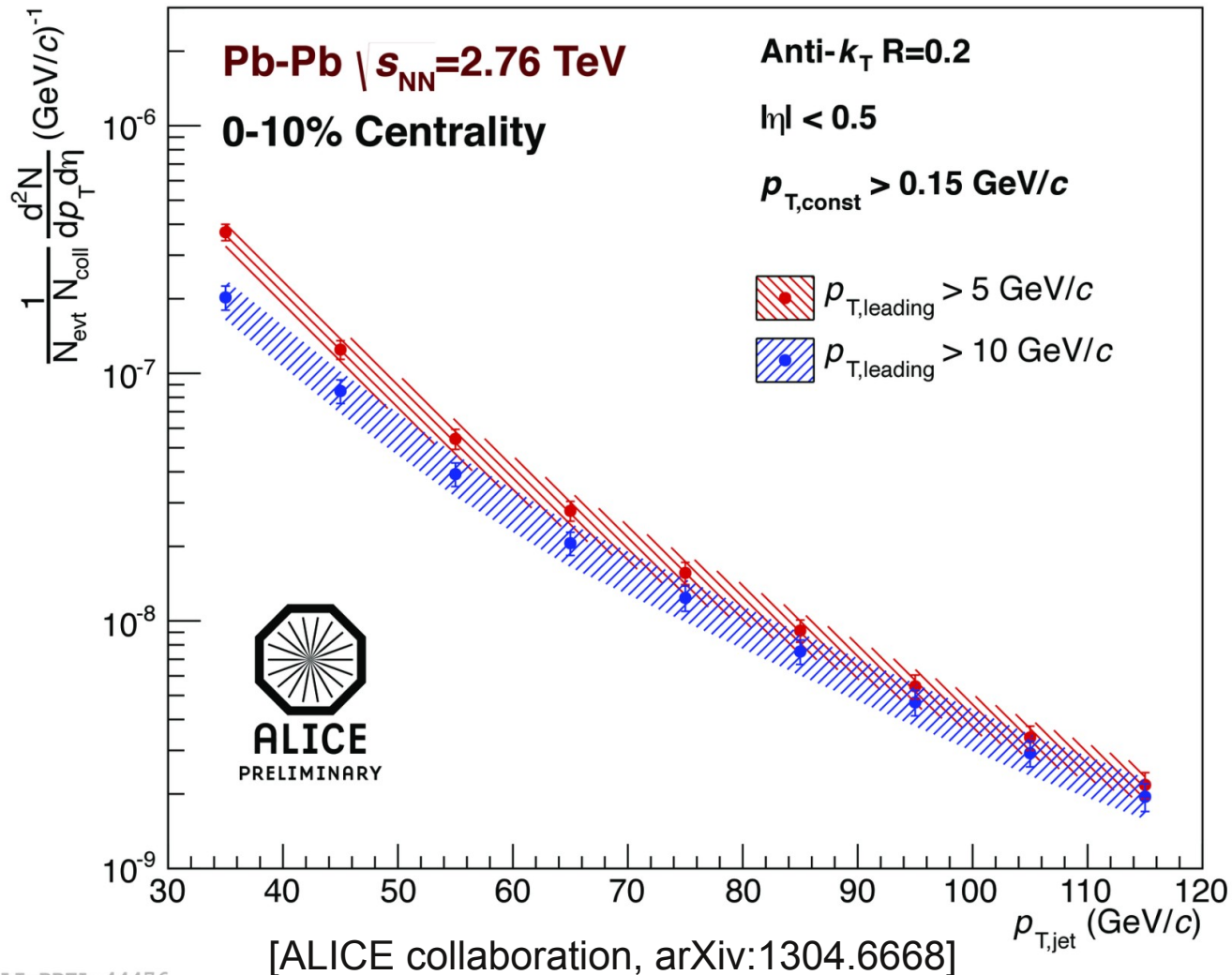
[G. de Barros et al, Hard Probes 2012, arXiv:1208.1518v2]

- Combinatorial background is reduced by p_T cut on jet's leading hadron
- Large fraction of jet energy can be carried by soft radiation
 - May be important for jet quenching measurements



Methodology

- Method has been successfully used by ALICE collaboration
- We have tested the method on toymodel and now we are applying the same method to STAR data



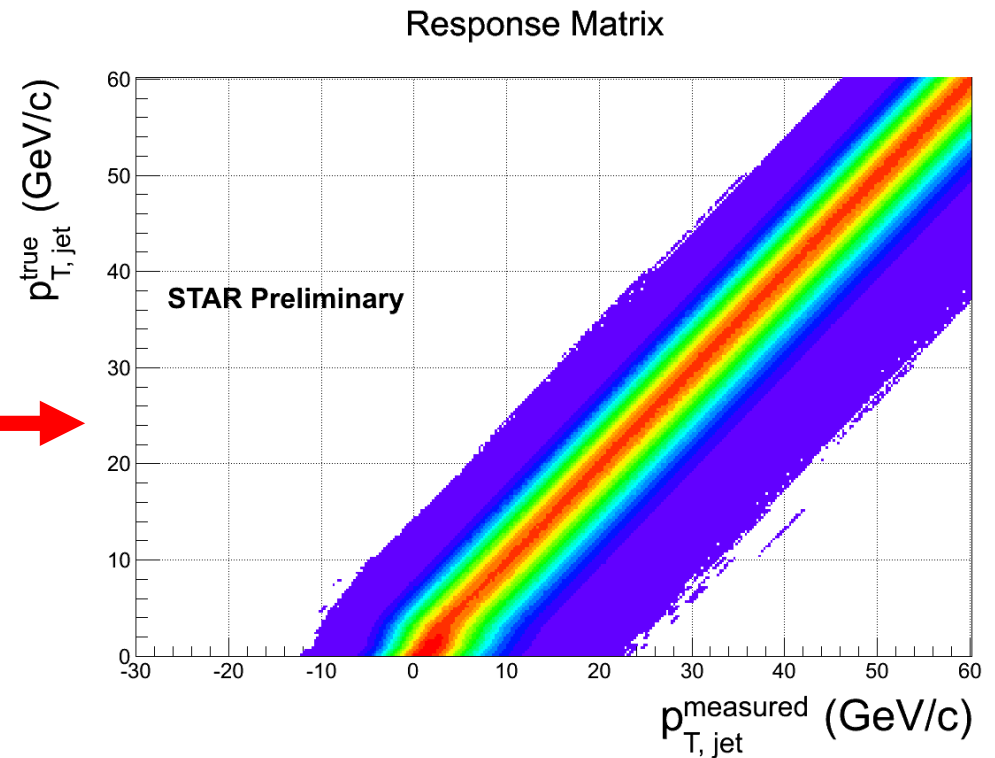
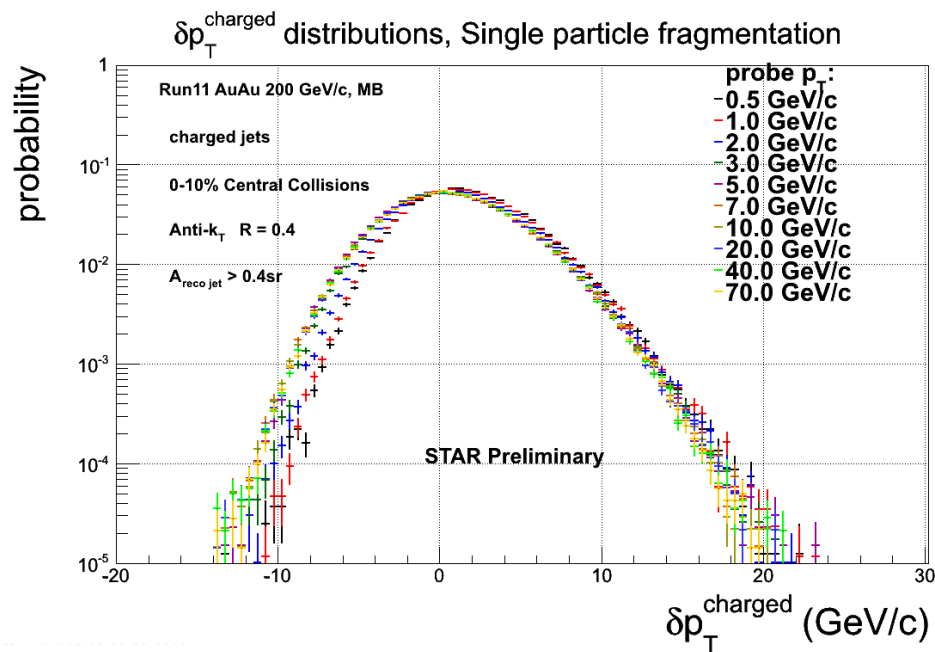
Response of jet to soft background: $\delta(p_T)$ distribution

- embedding of a simulated jet into an event => jet reconstruction

$$\delta p_T = p_{T,corr} - p_{T,emb} = p_T - A_{jet} \times \rho - p_{T,emb}$$

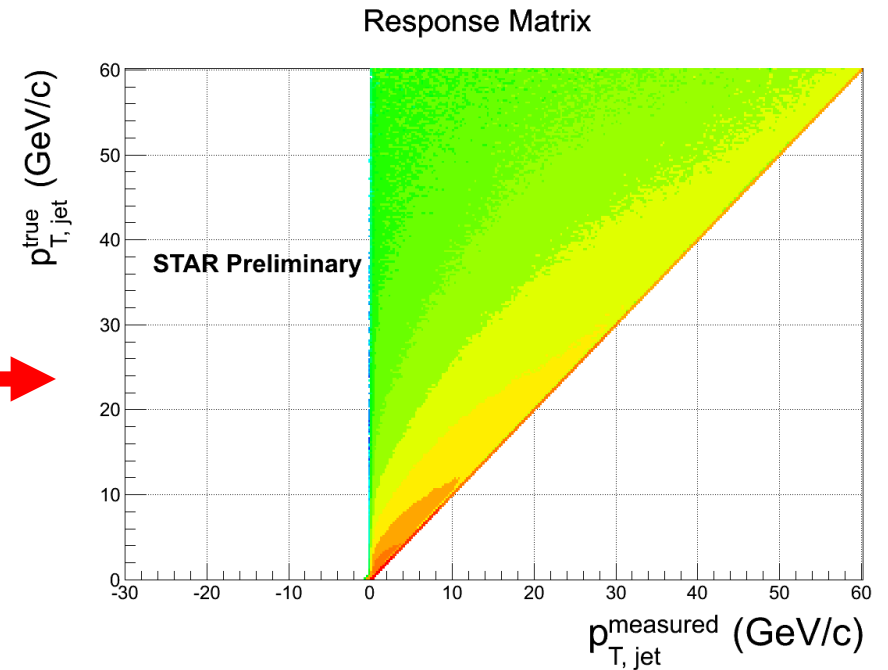
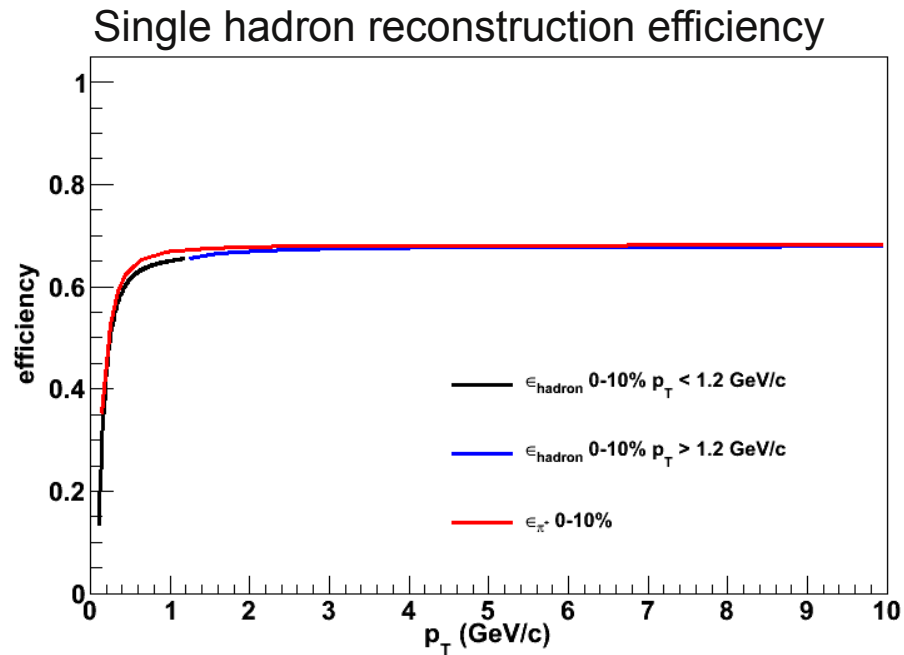
$$\rho \sim 29 \text{ GeV/Sr}, A_{R=0.4} \sim 0.5 \text{ Sr}$$

- ensemble-averaged $\delta(p_T)$ distribution => determination of the response matrix



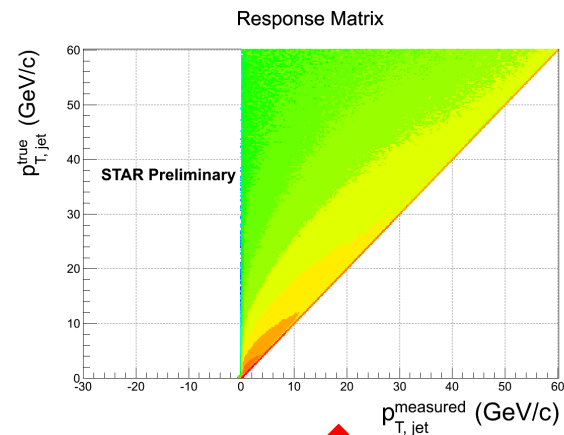
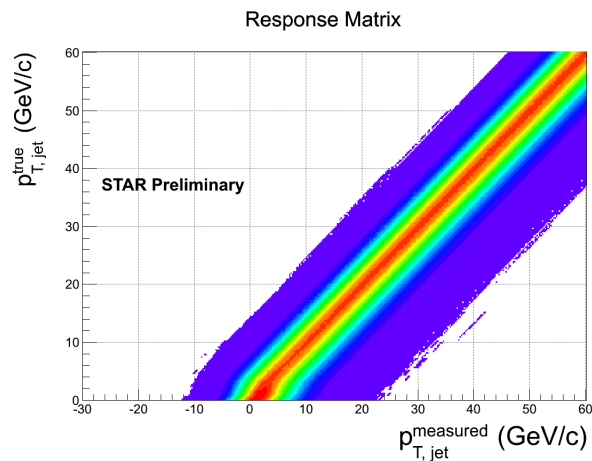
Response of jet due to detector efficiency

- detector response: dominated by tracking efficiency
- we utilize PYTHIA jets (to do: other fragmentation models)

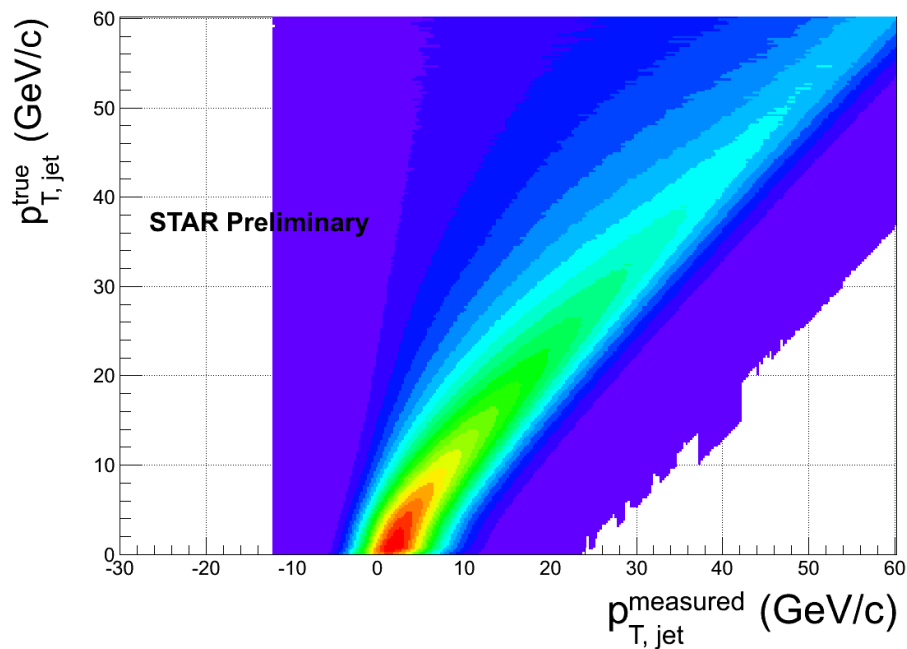


- tracking efficiency: $\sim 5\%$ uncertainty

Response of jet to soft background and detector efficiency



Response Matrix



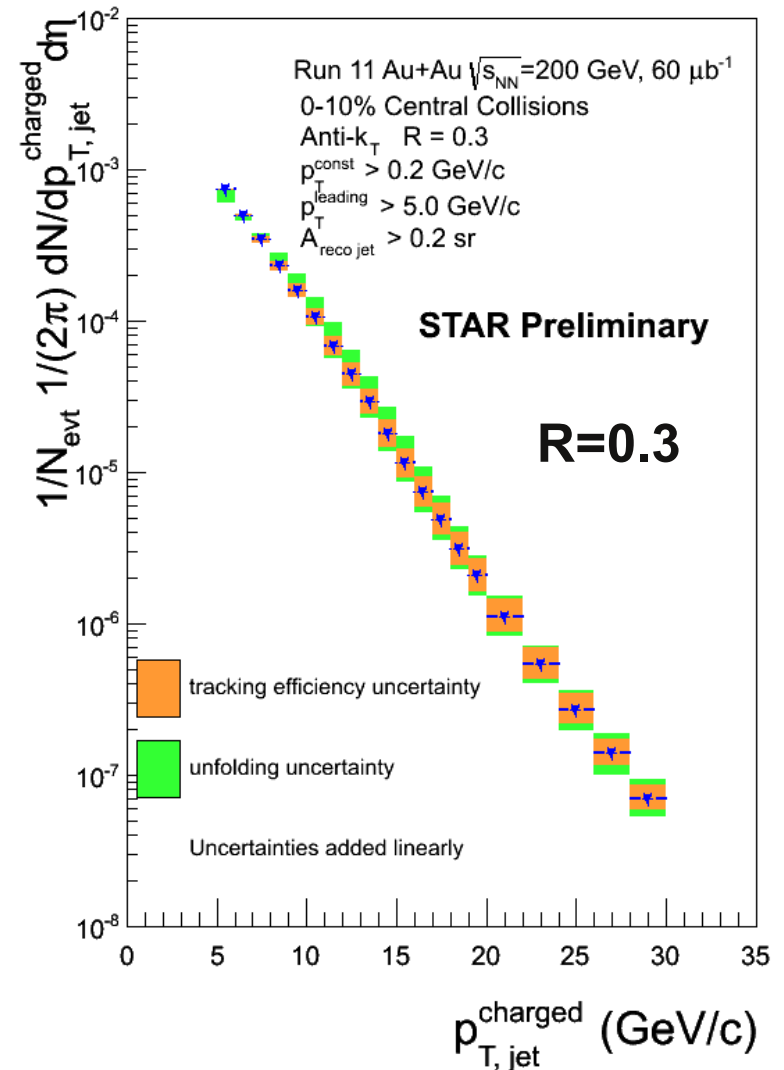
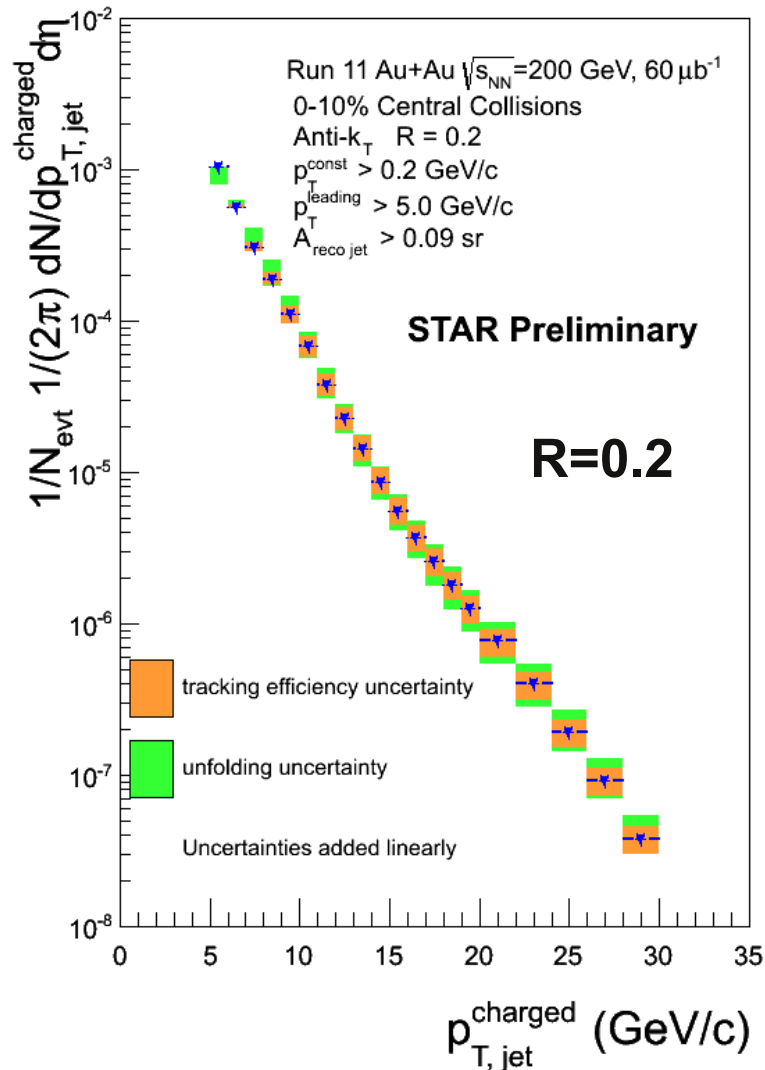
Overall response matrix:

Corrected Spectra

$$p_T^{\text{leading}} > 5 \text{ GeV}/c$$

Charged jets

Charged jets



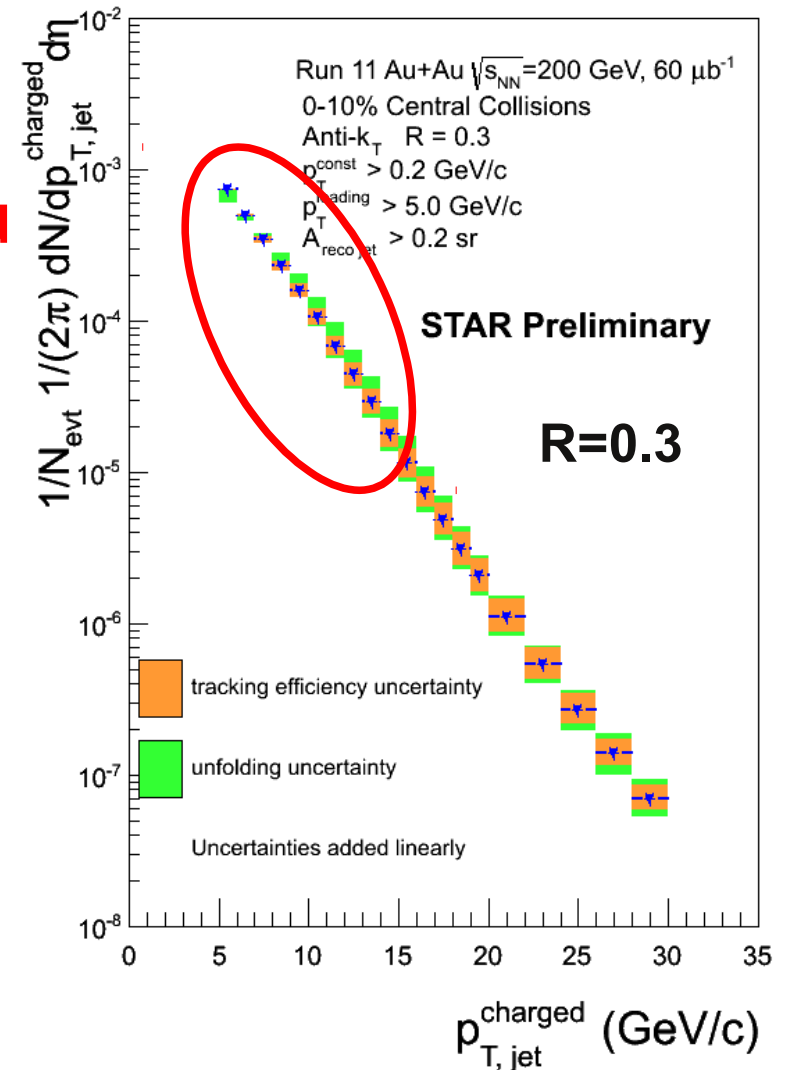
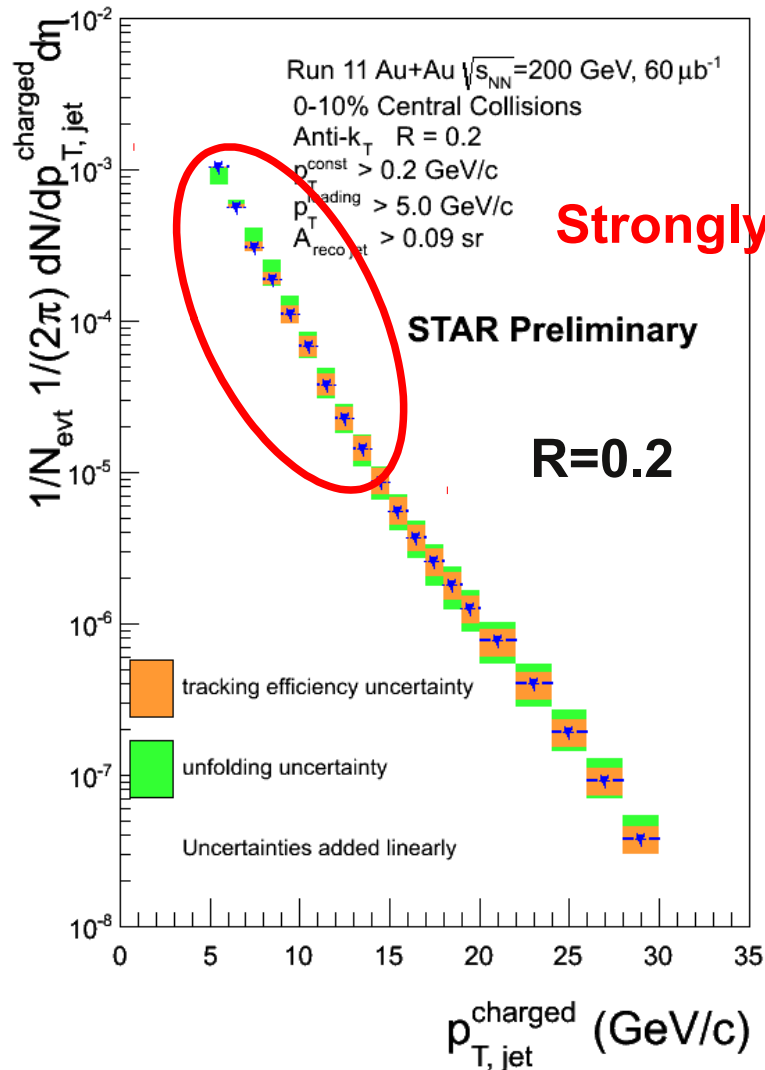
Jet energy scale resolution: $\sim 5\%$ (due to tracking efficiency uncertainty)

Corrected Spectra

$$p_T^{\text{leading}} > 5 \text{ GeV}/c$$

Charged jets

Charged jets

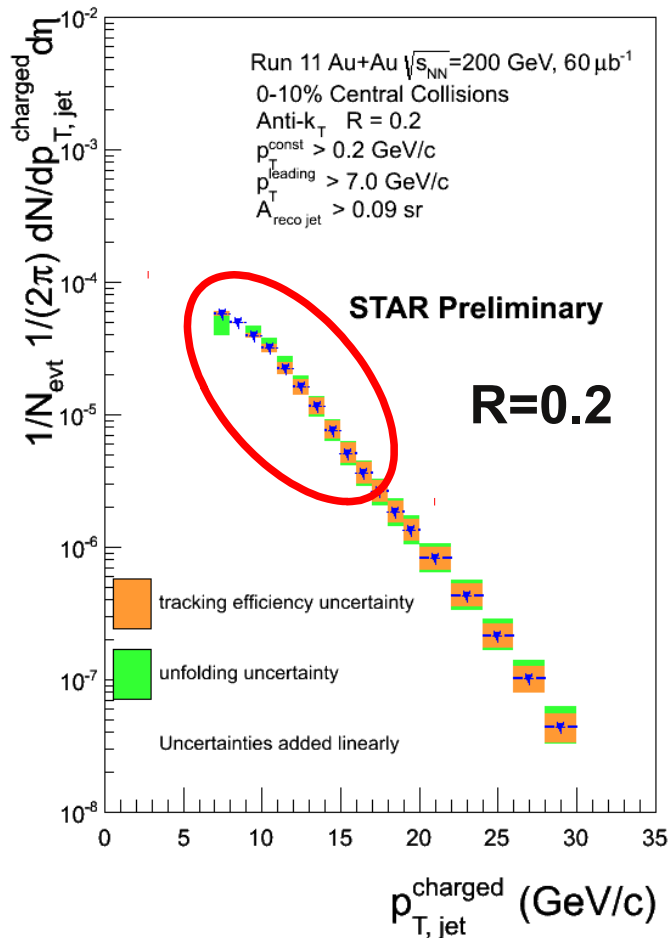


Jet energy scale resolution: $\sim 5\%$ (due to tracking efficiency uncertainty)

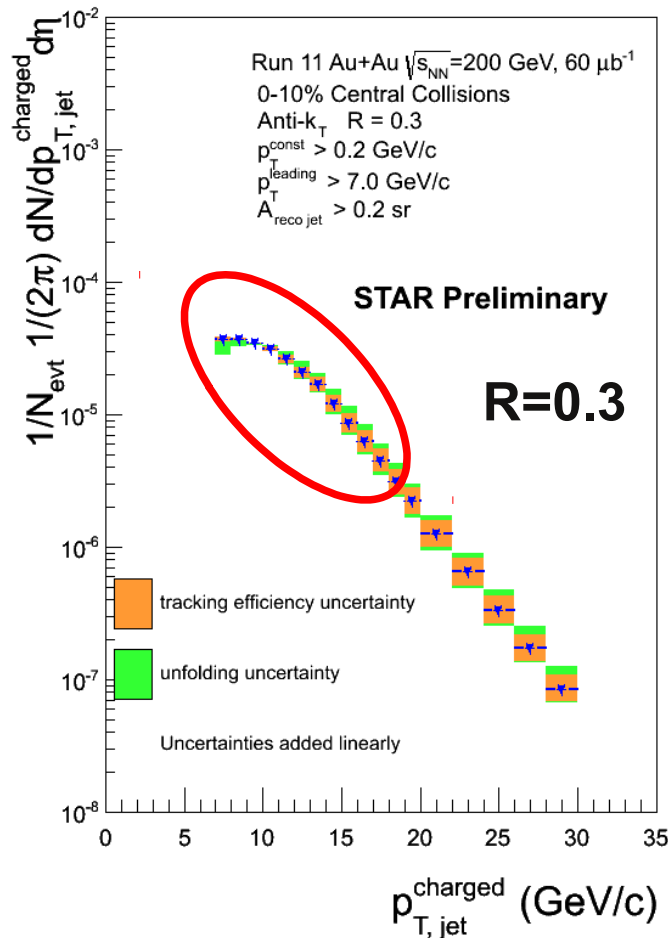
Corrected Spectra

$$p_T^{\text{leading}} > 7 \text{ GeV}/c$$

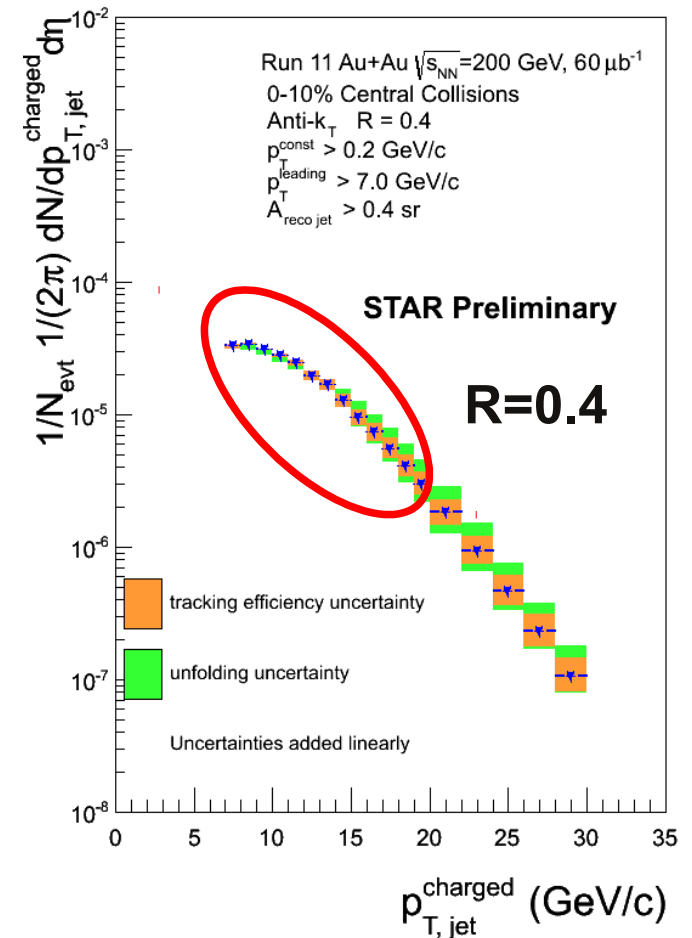
Charged jets



Charged jets

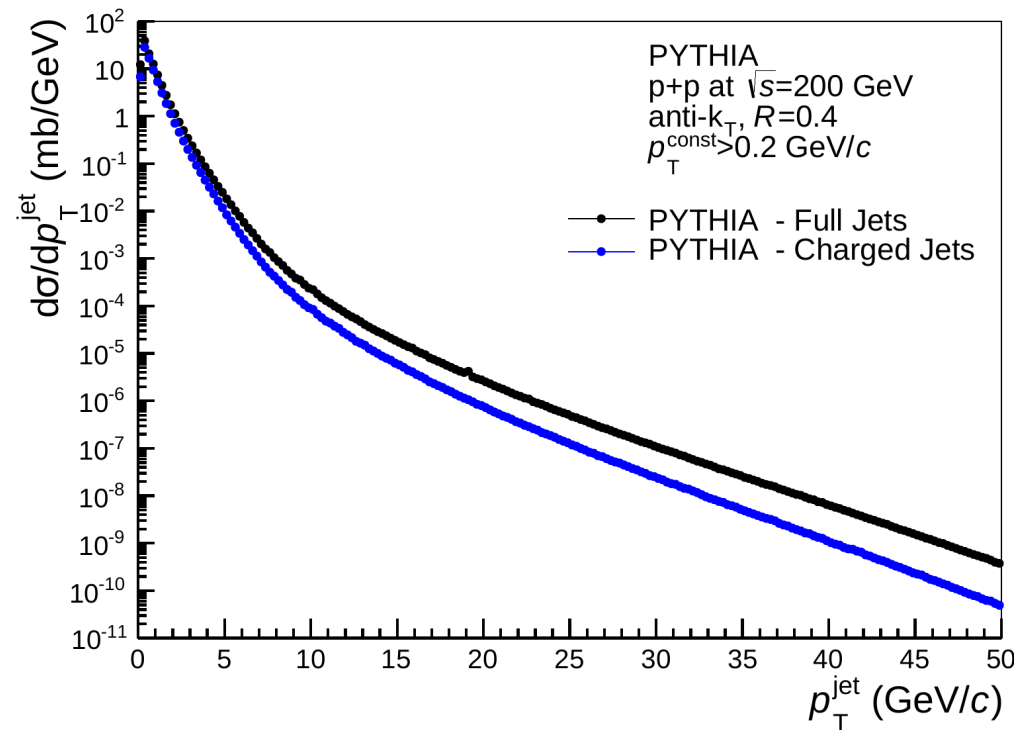
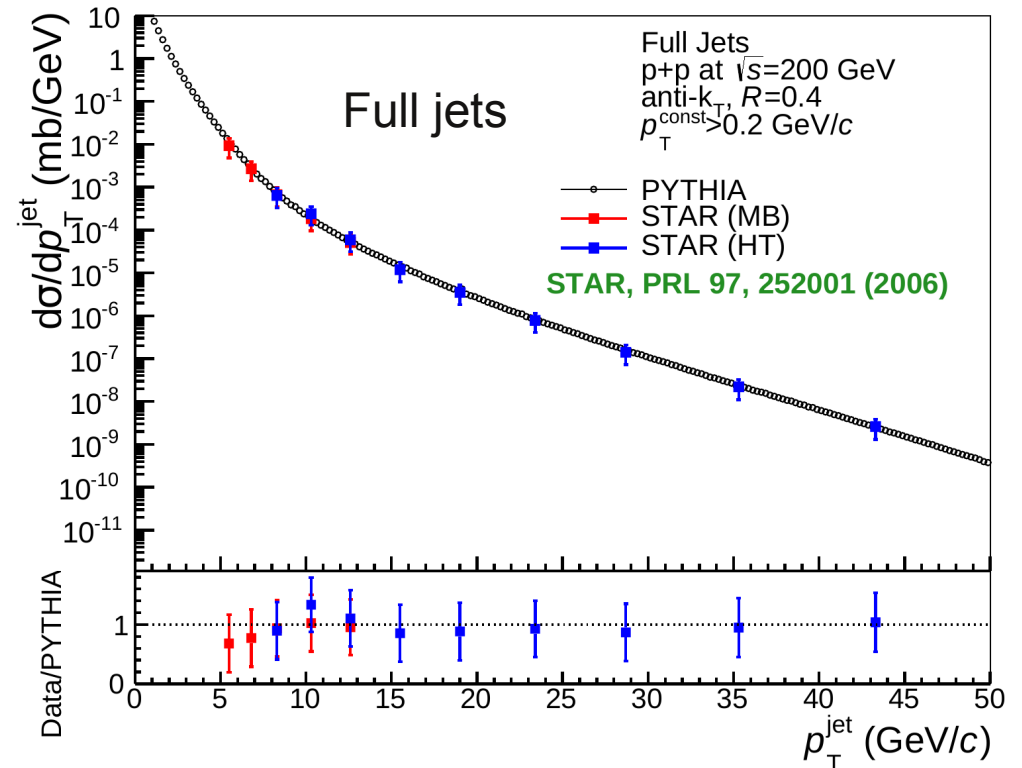


Charged jets

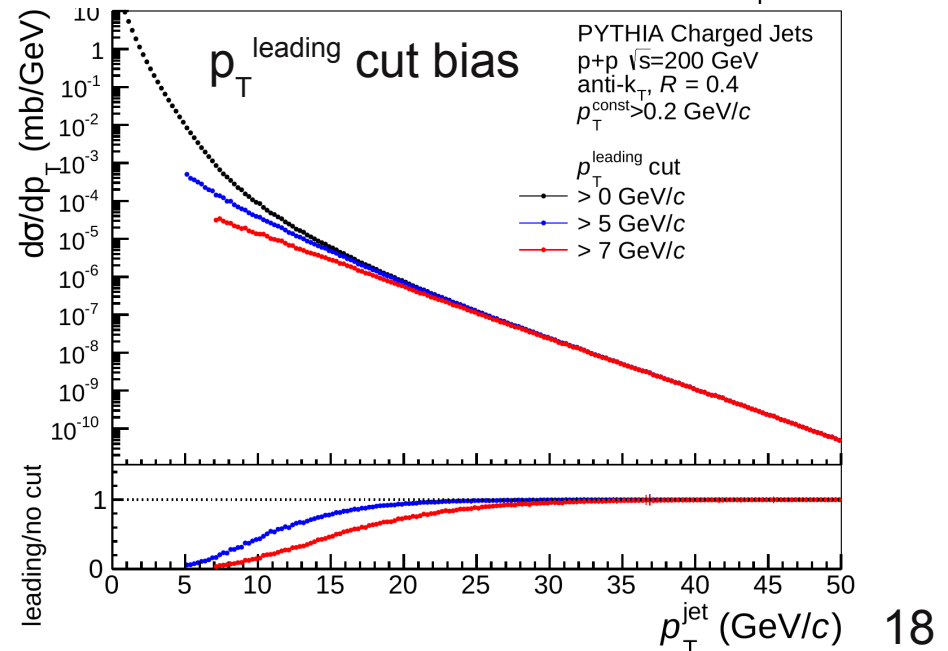


p+p reference

- Pythia 6.4 - Tune A used as p+p reference agrees well with STAR p+p data (for full jets)

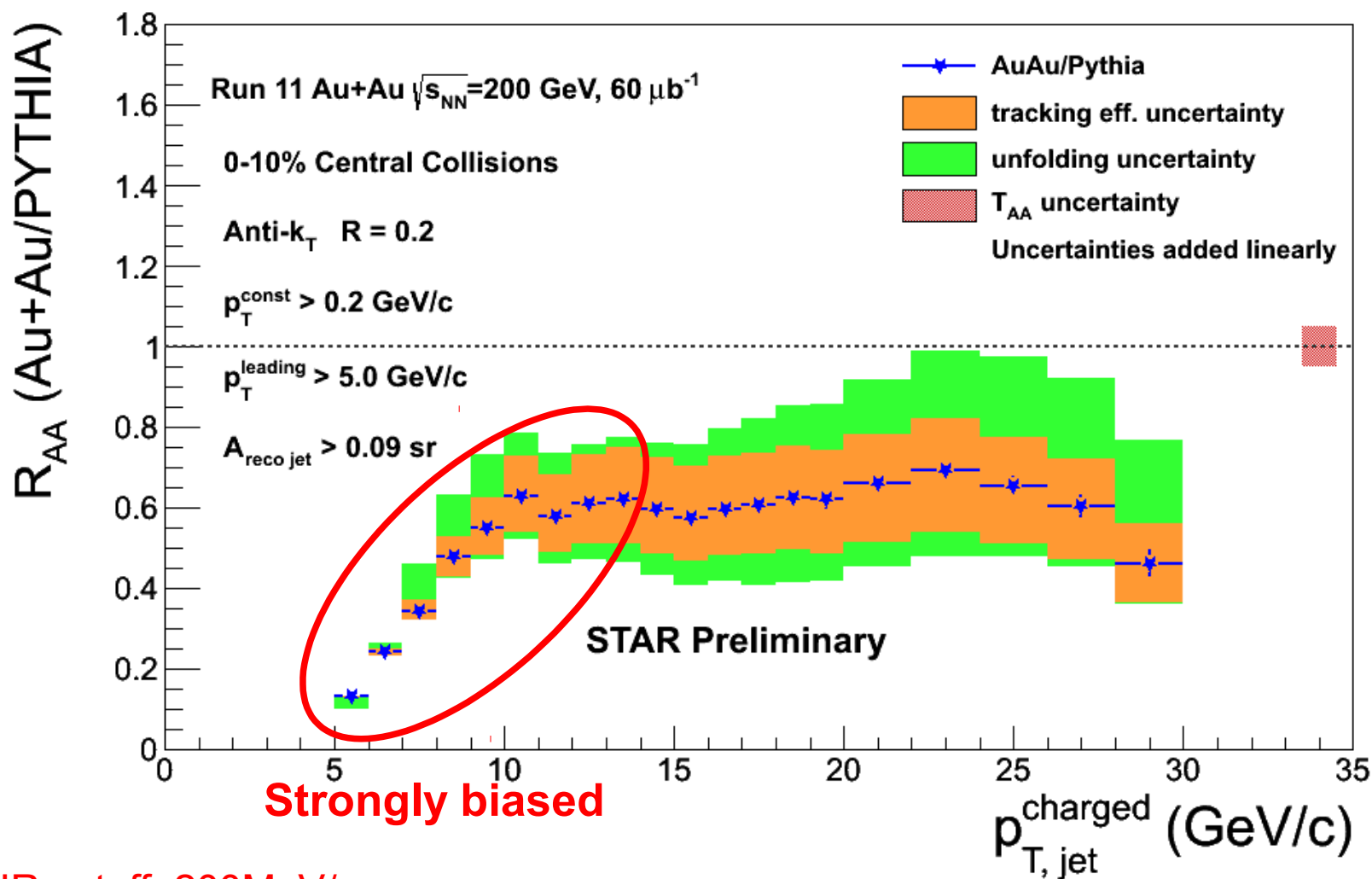


- Pythia: full jets \rightarrow charged jets (accept only final particles which are charged)
- p_T^{leading} bias – up to 3-4x p_T^{leading} ($R=0.2 - 0.4$)



$R=0.2, p_T^{\text{leading}} > 5 \text{ GeV}/c$

Charged jets

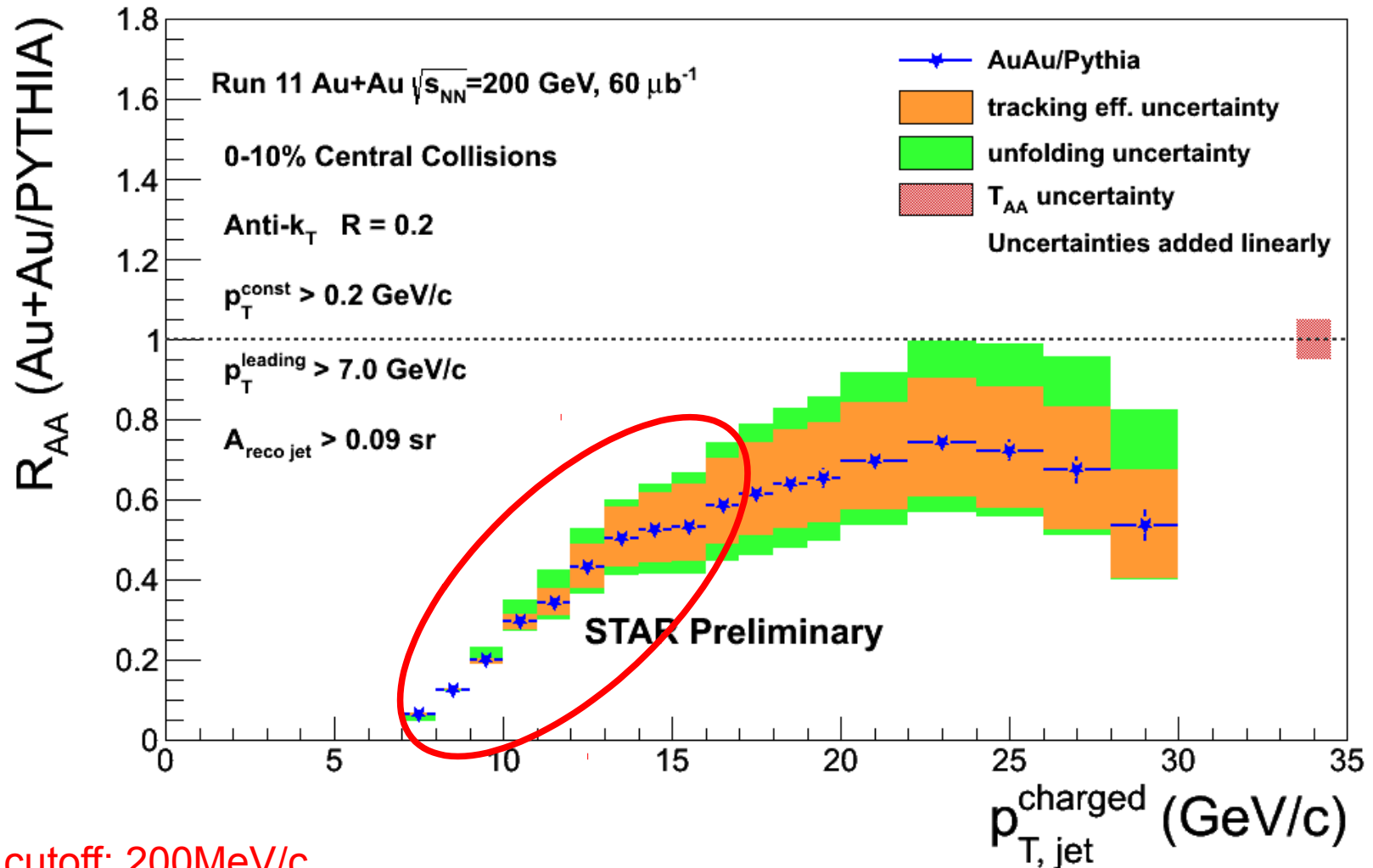


IR cutoff: 200MeV/c

Yield-weighted jet R_{AA} is consistent with inclusive hadron R_{AA} at $p_T \sim 5 \text{ GeV}$

$R=0.2, p_T^{\text{leading}} > 7 \text{ GeV}/c$

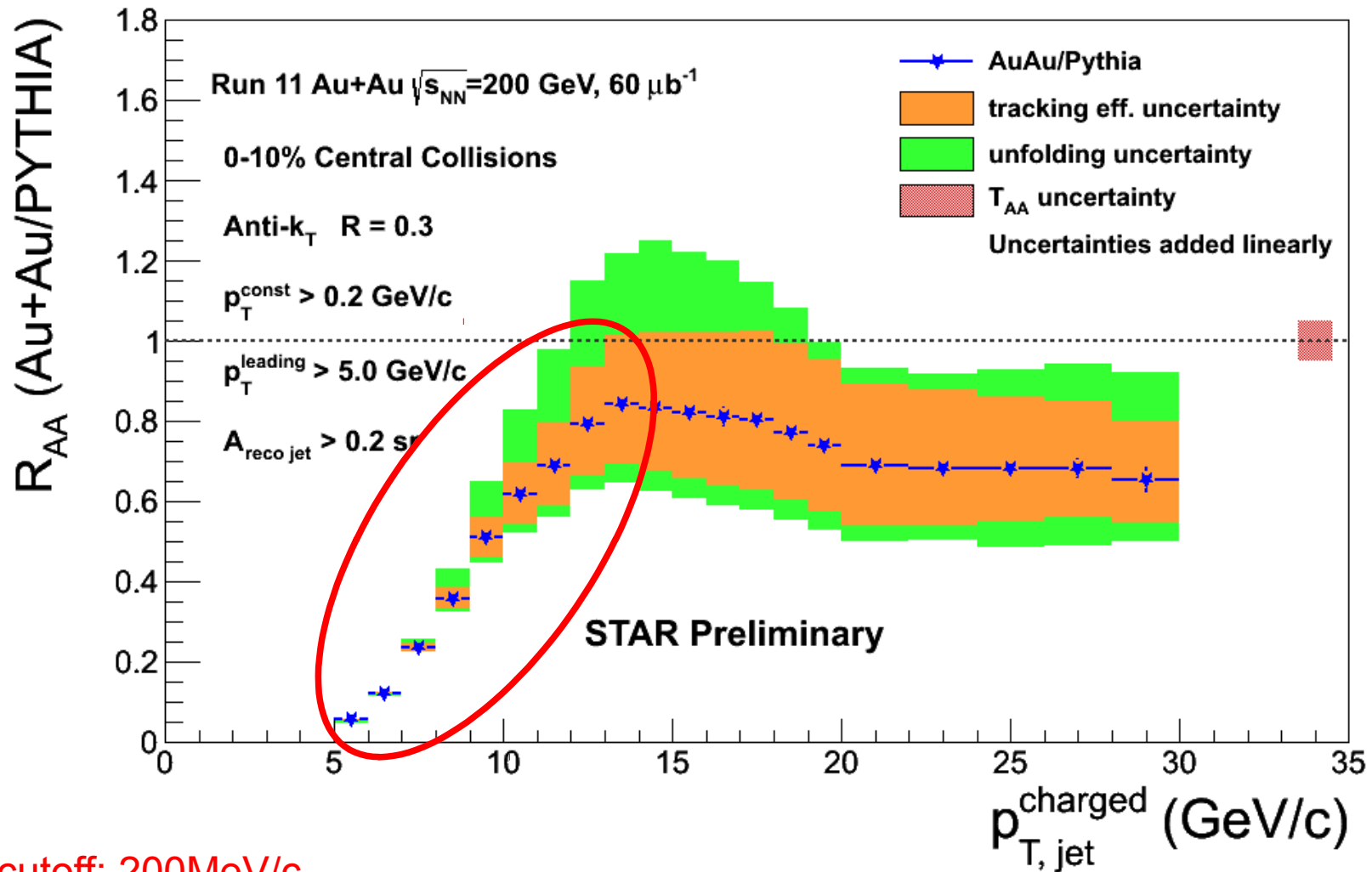
Charged jets



IR cutoff: 200MeV/c

$R=0.3, p_T^{\text{leading}} > 5 \text{ GeV}/c$

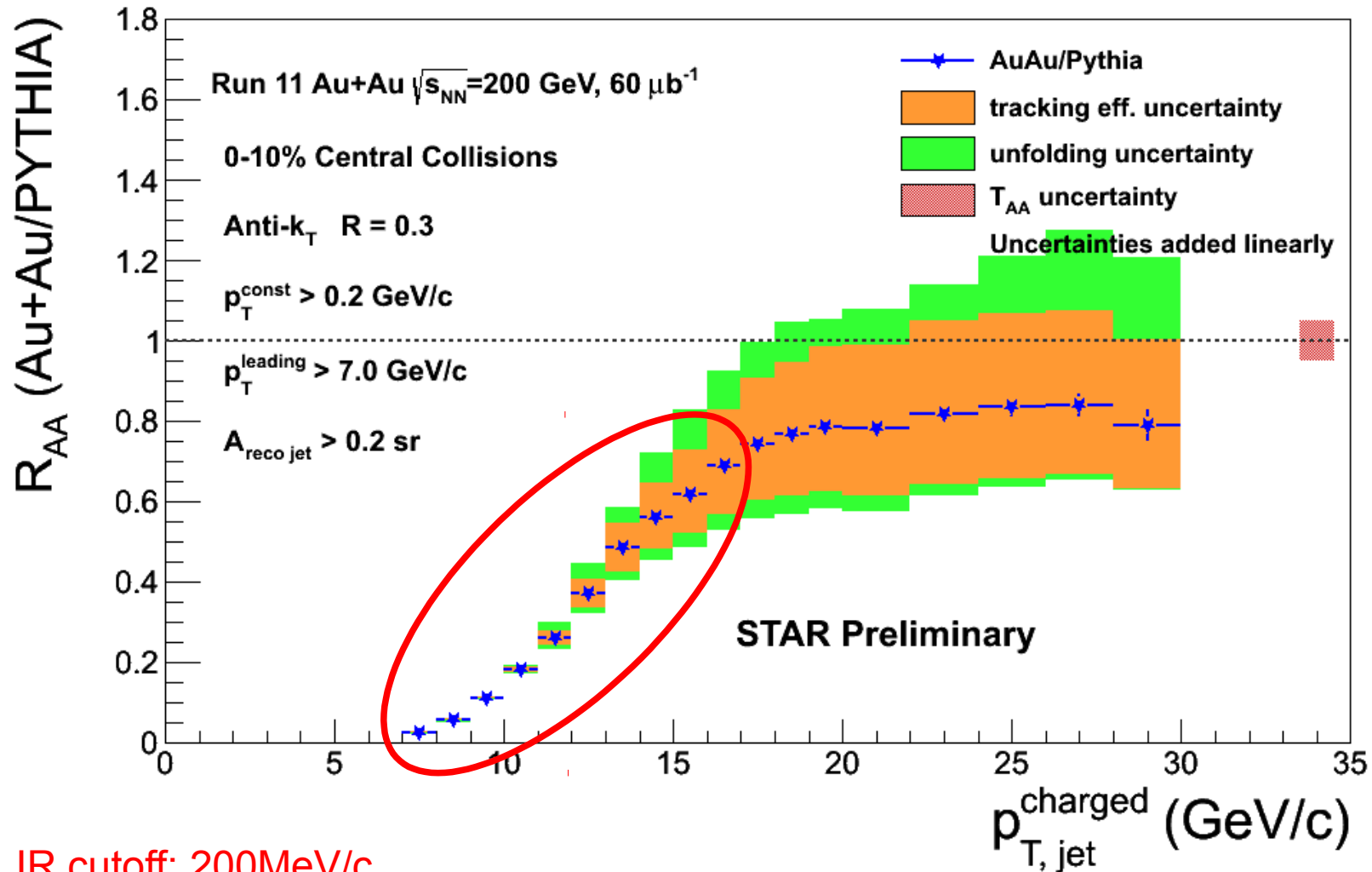
Charged jets



IR cutoff: 200MeV/c

$R=0.3, p_T^{\text{leading}} > 7 \text{ GeV}/c$

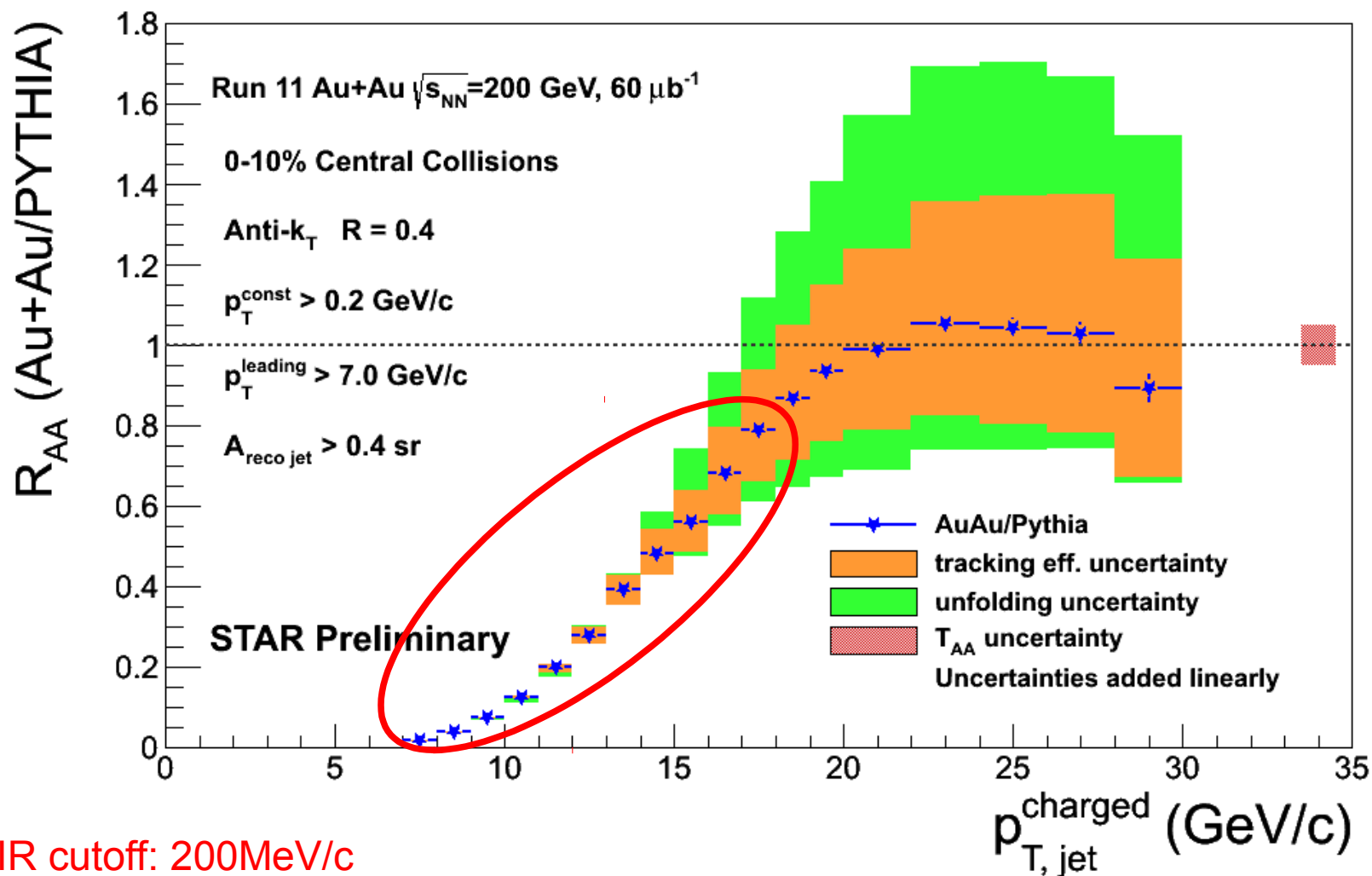
Charged jets



IR cutoff: 200MeV/c

$R=0.4, p_T^{\text{leading}} > 7 \text{ GeV}/c$

Charged jets

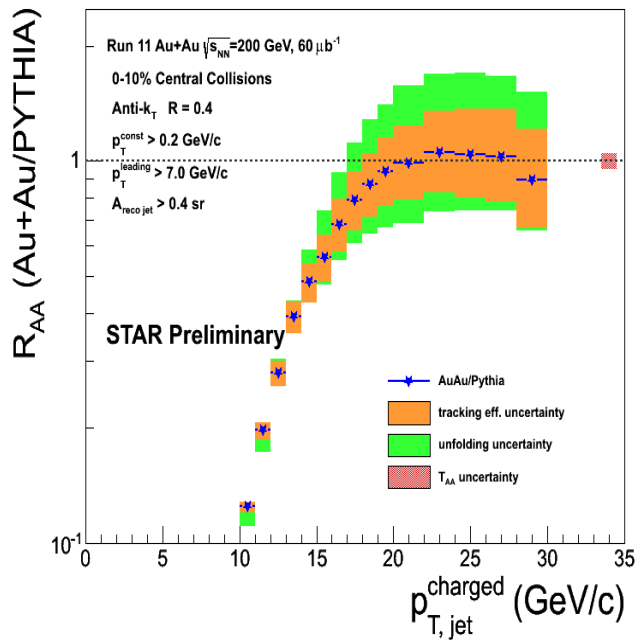


IR cutoff: 200MeV/c

Comparison of old and new STAR results

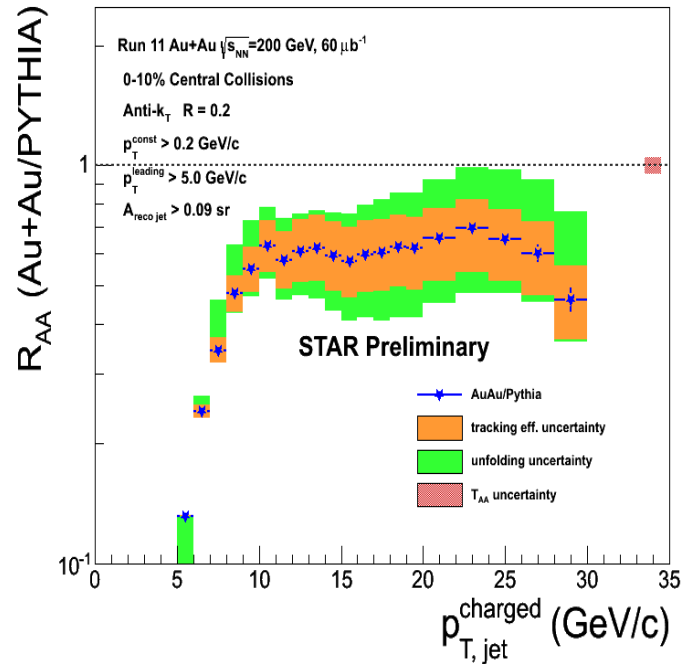
R=0.4

Charged jets



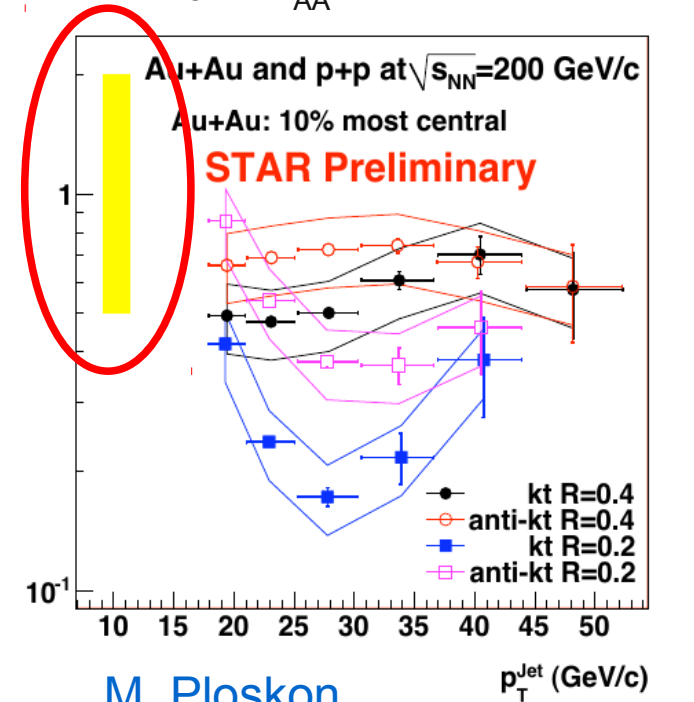
R=0.2

Charged jets



QM '09 results

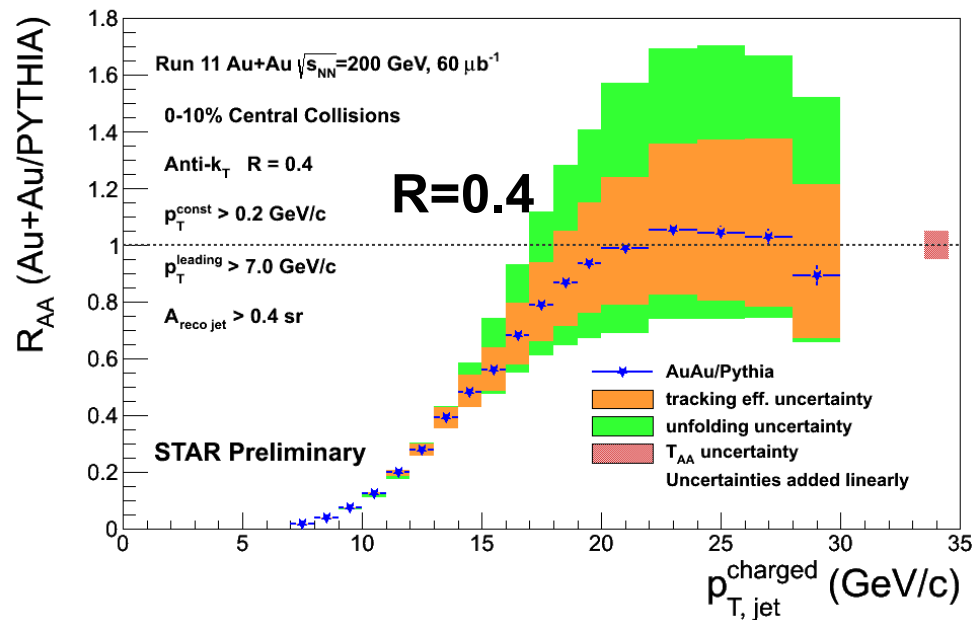
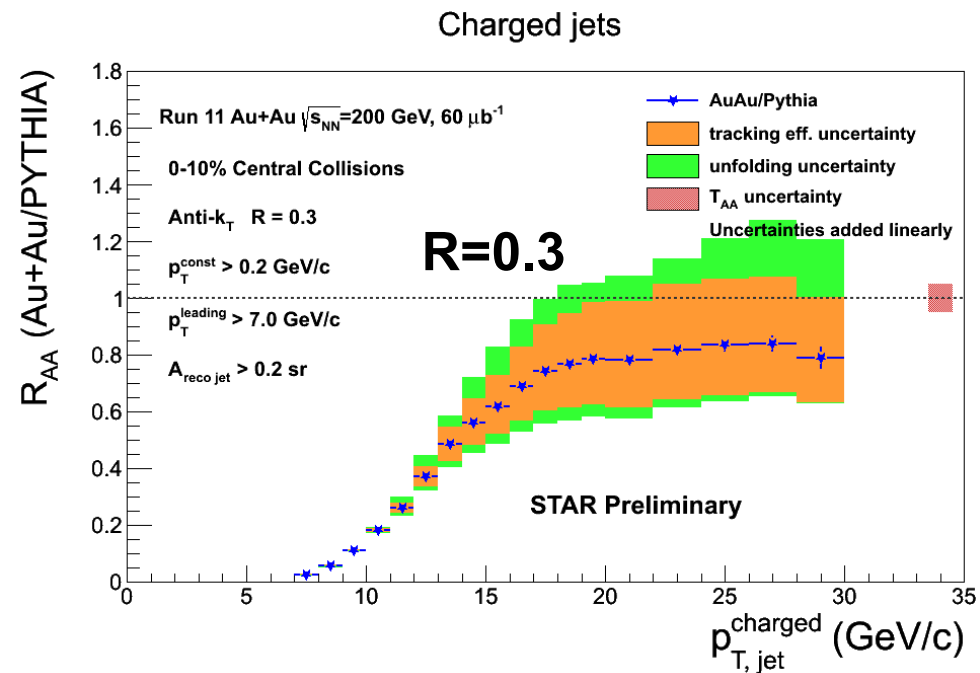
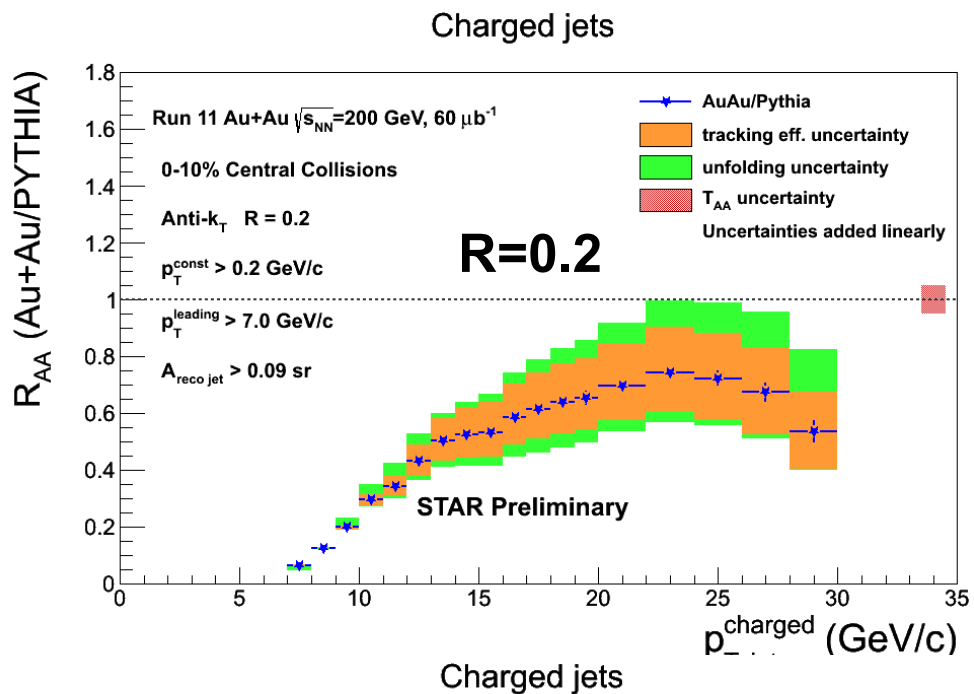
Full jet R_{AA}



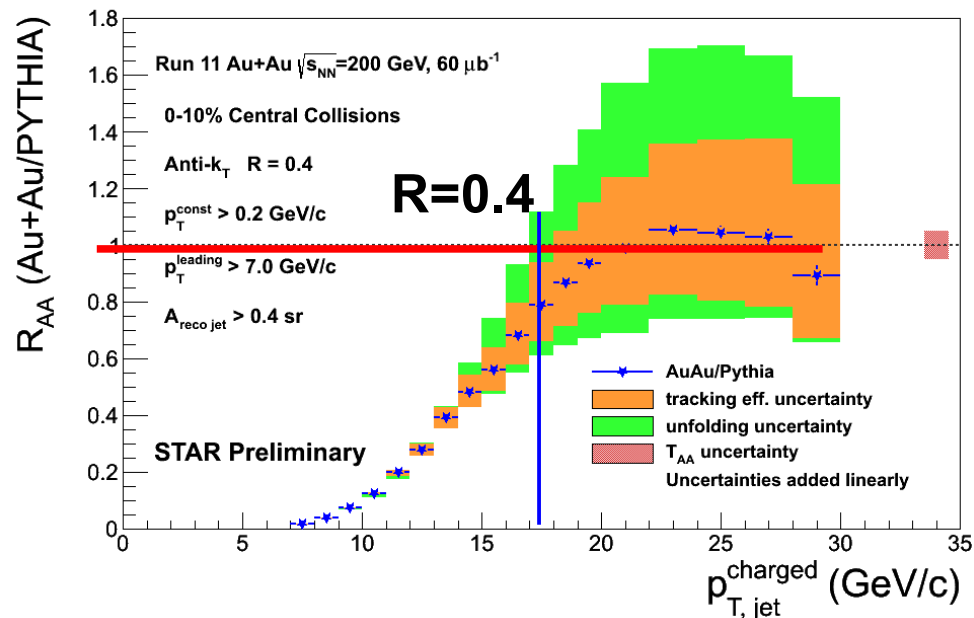
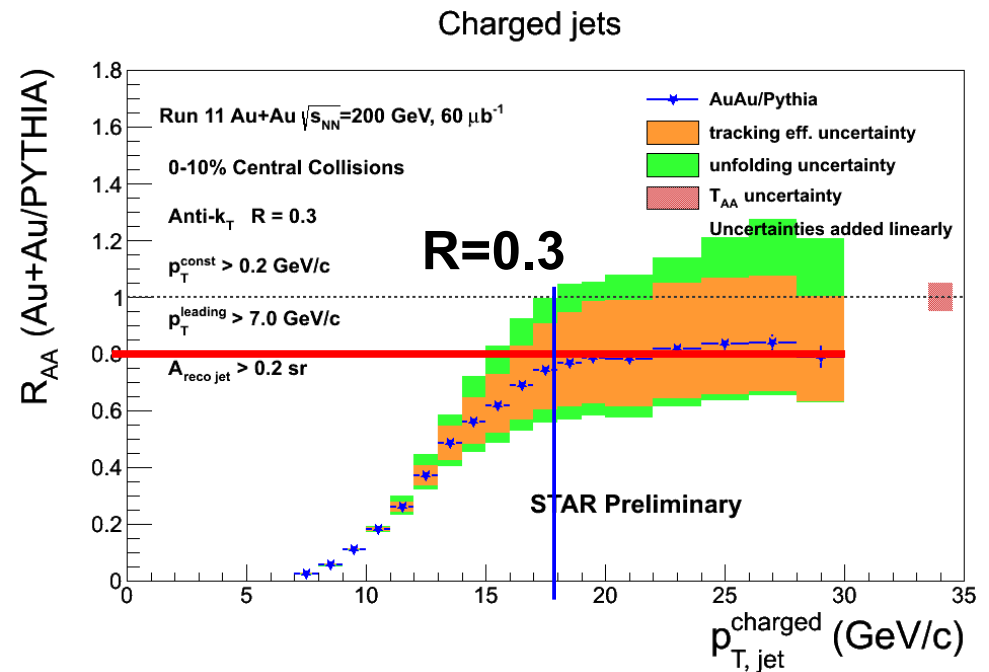
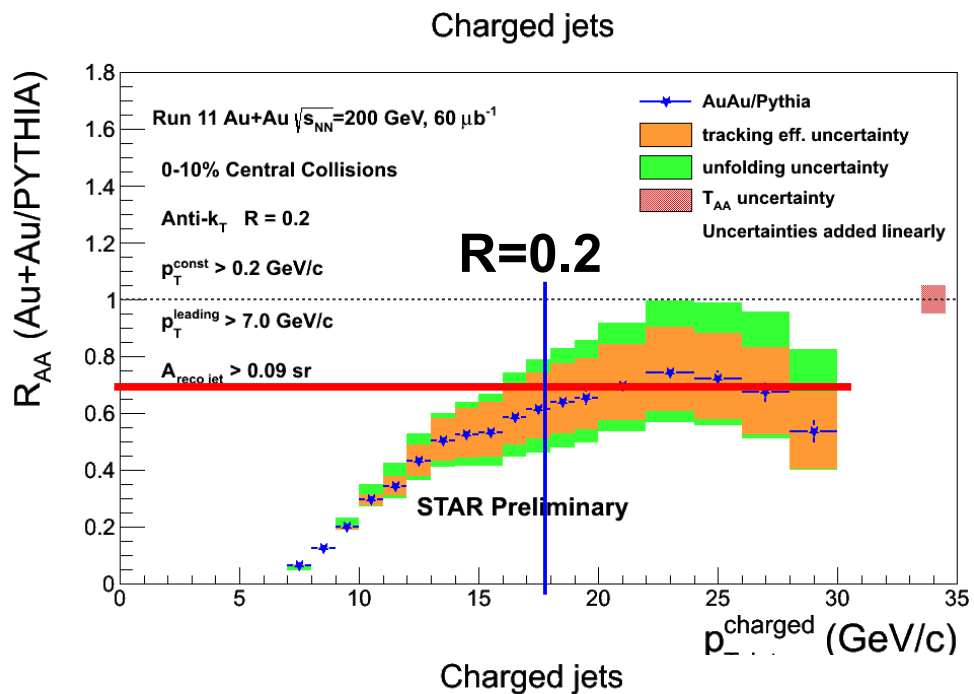
M. Ploskon,
Nucl.Phys.A830:255c-258c,
2009

- Results are consistent
- Old results – large systematic uncertainty on normalization
- New methodology – better understood systematics

Jet R_{AA} for different R

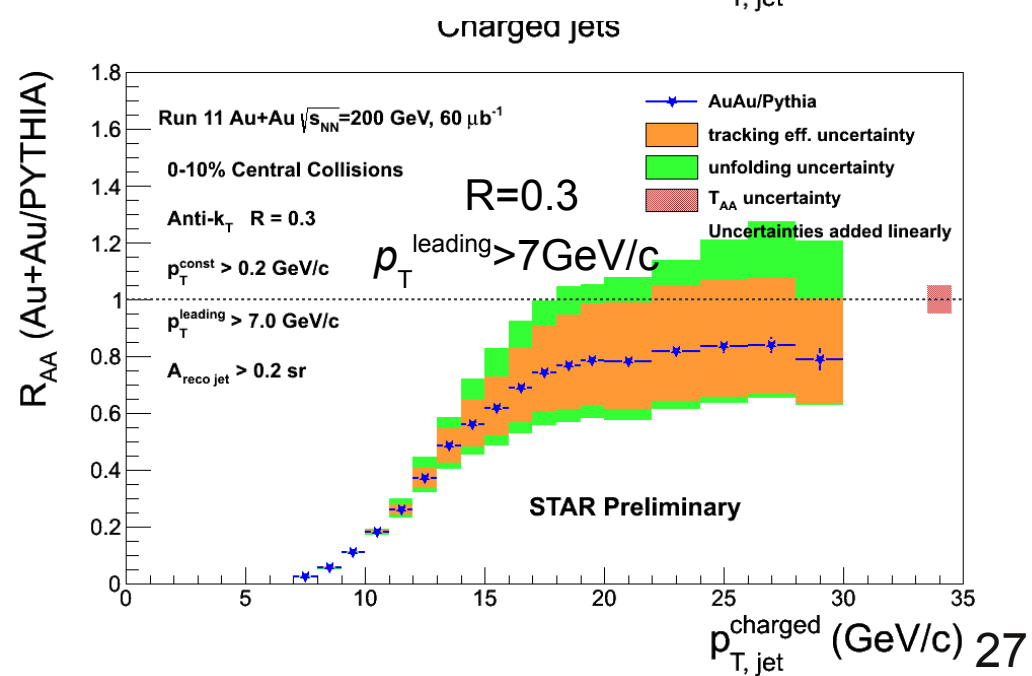
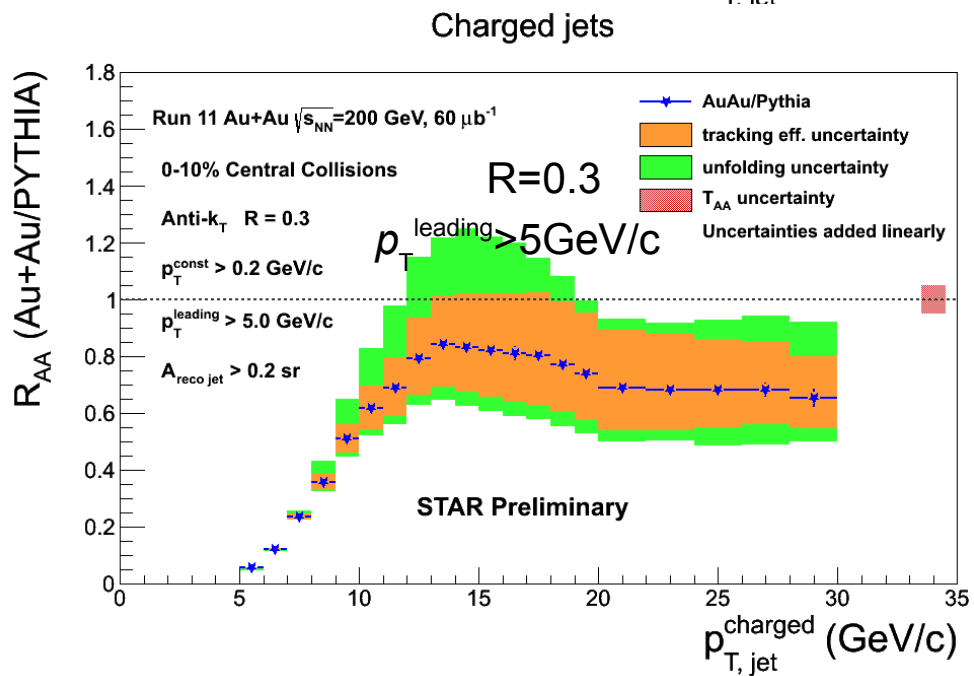
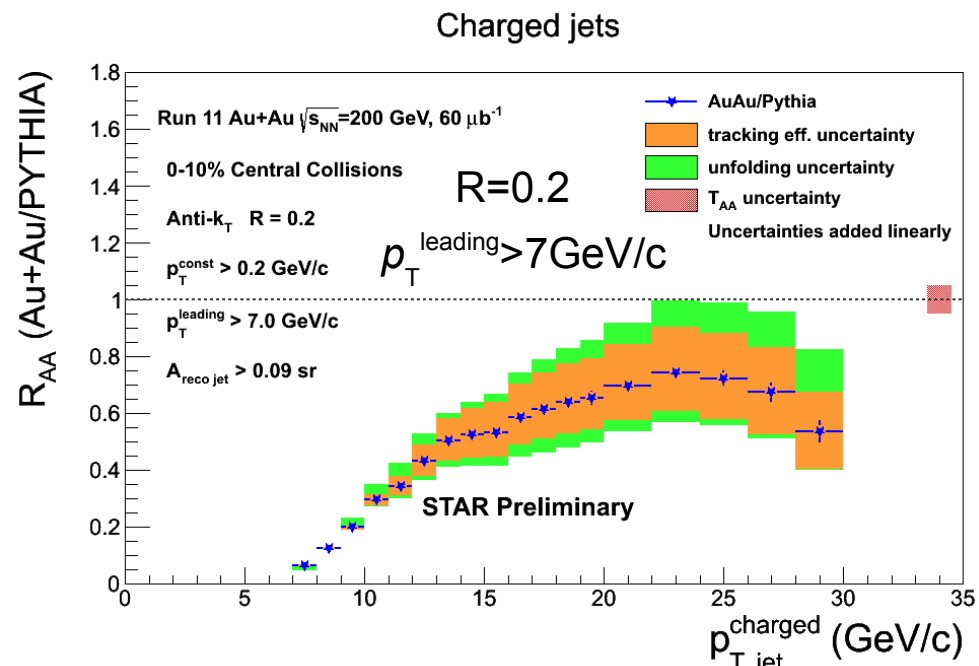
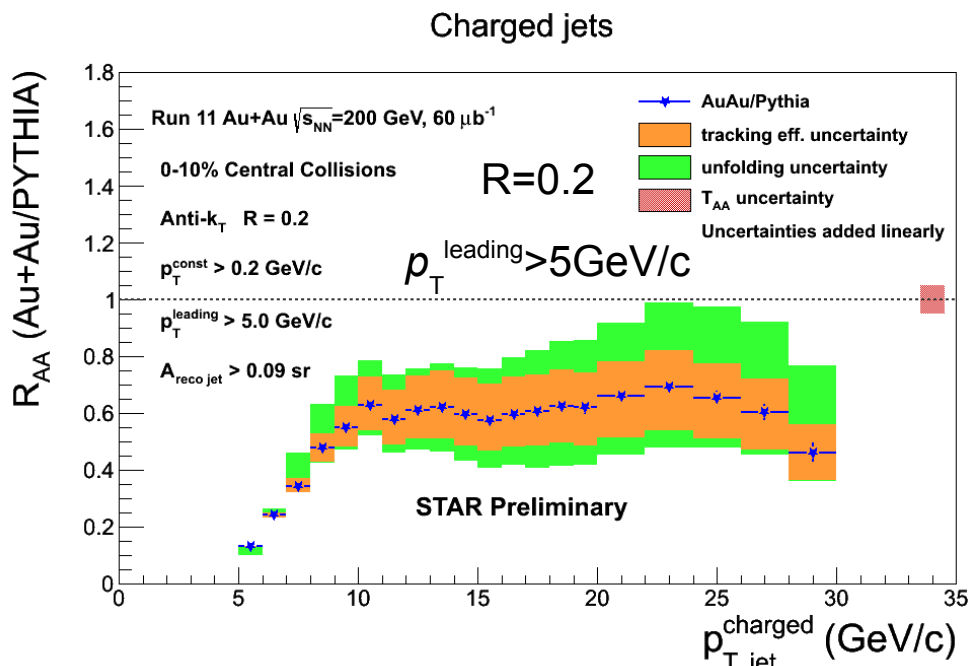


Jet R_{AA} for different R

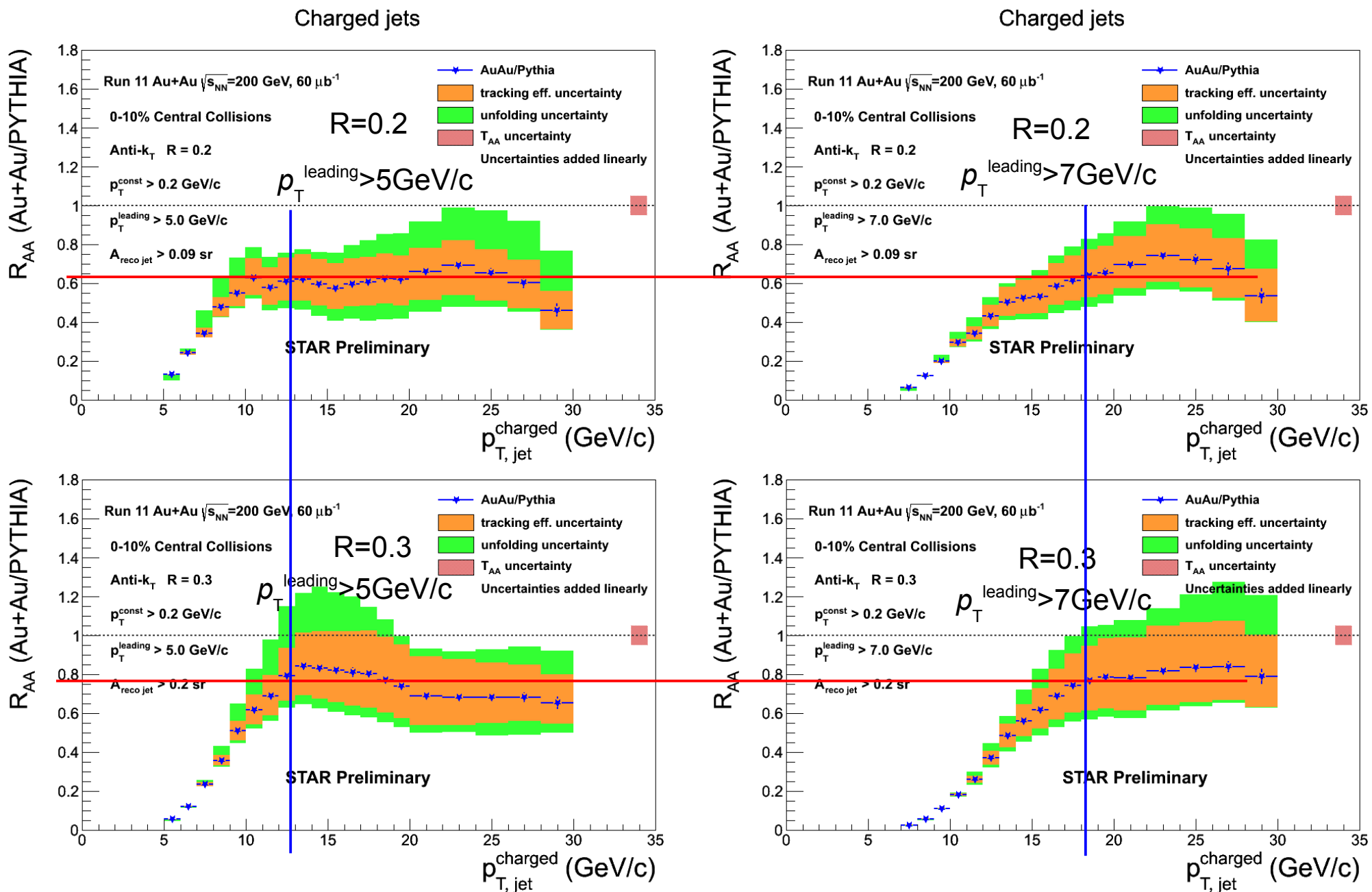


- Growth in R_{AA} central value with rising R ?
- Uncertainties too large for definitive statement
- To do: Ratios of distributions (correlated uncertainties cancel)

Jet R_{AA} for different p_T^{leading} cut



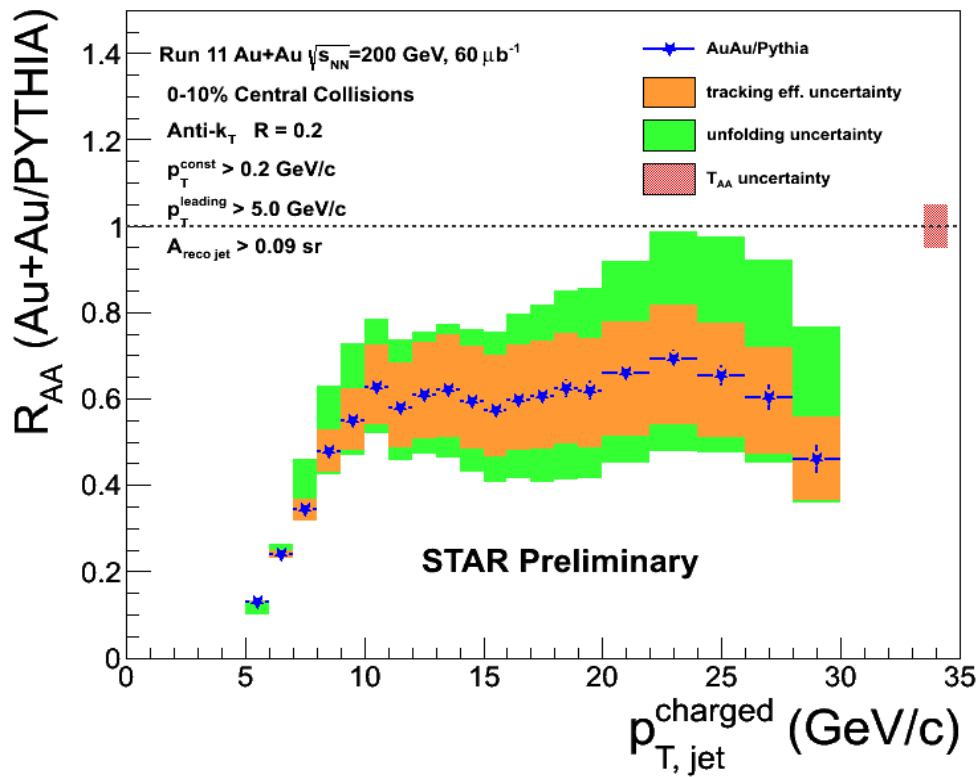
Jet R_{AA} for different p_T^{leading} cut



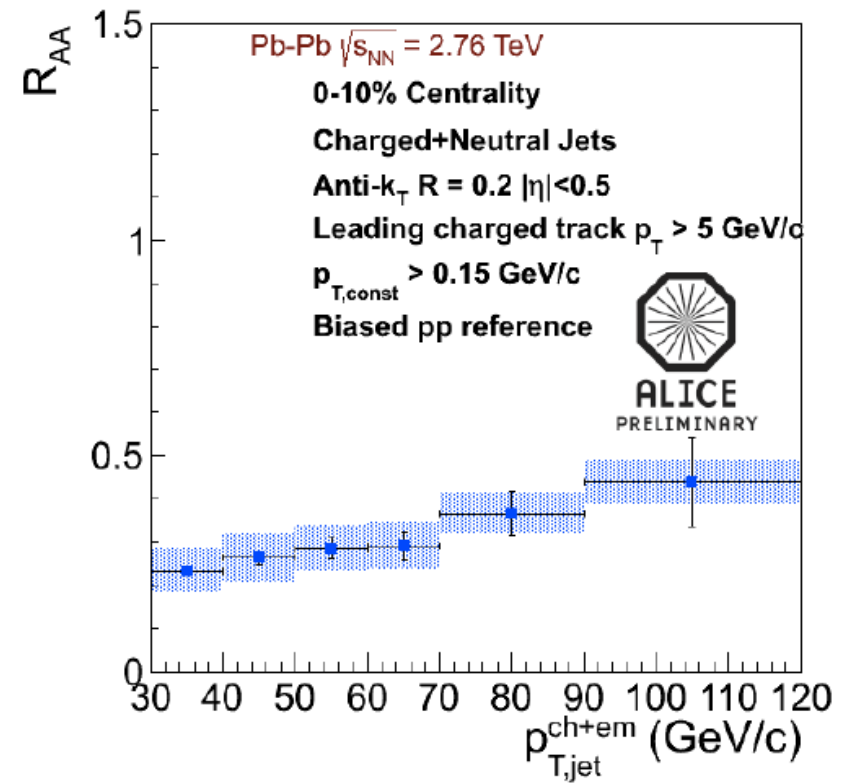
central value of R_{AA} approximately same for both values of p_T^{leading} cut

Comparison with ALICE results

Charged jets



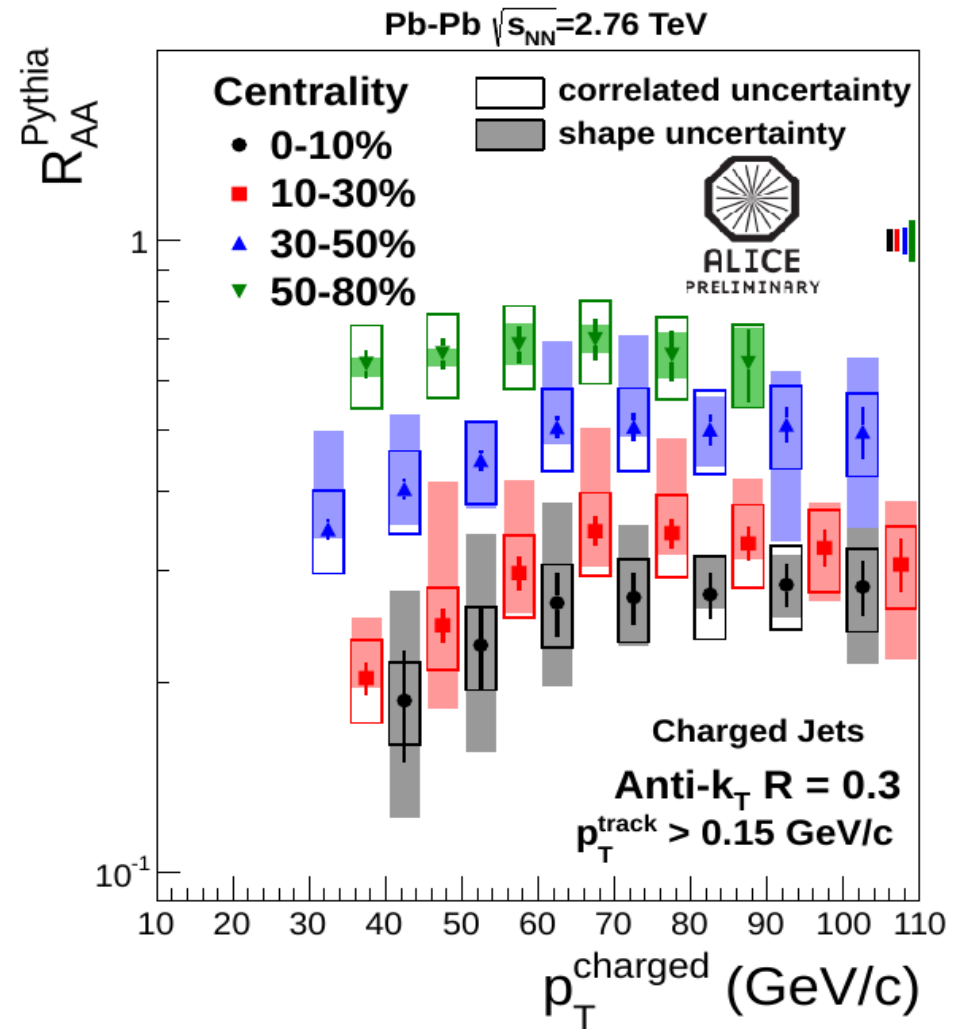
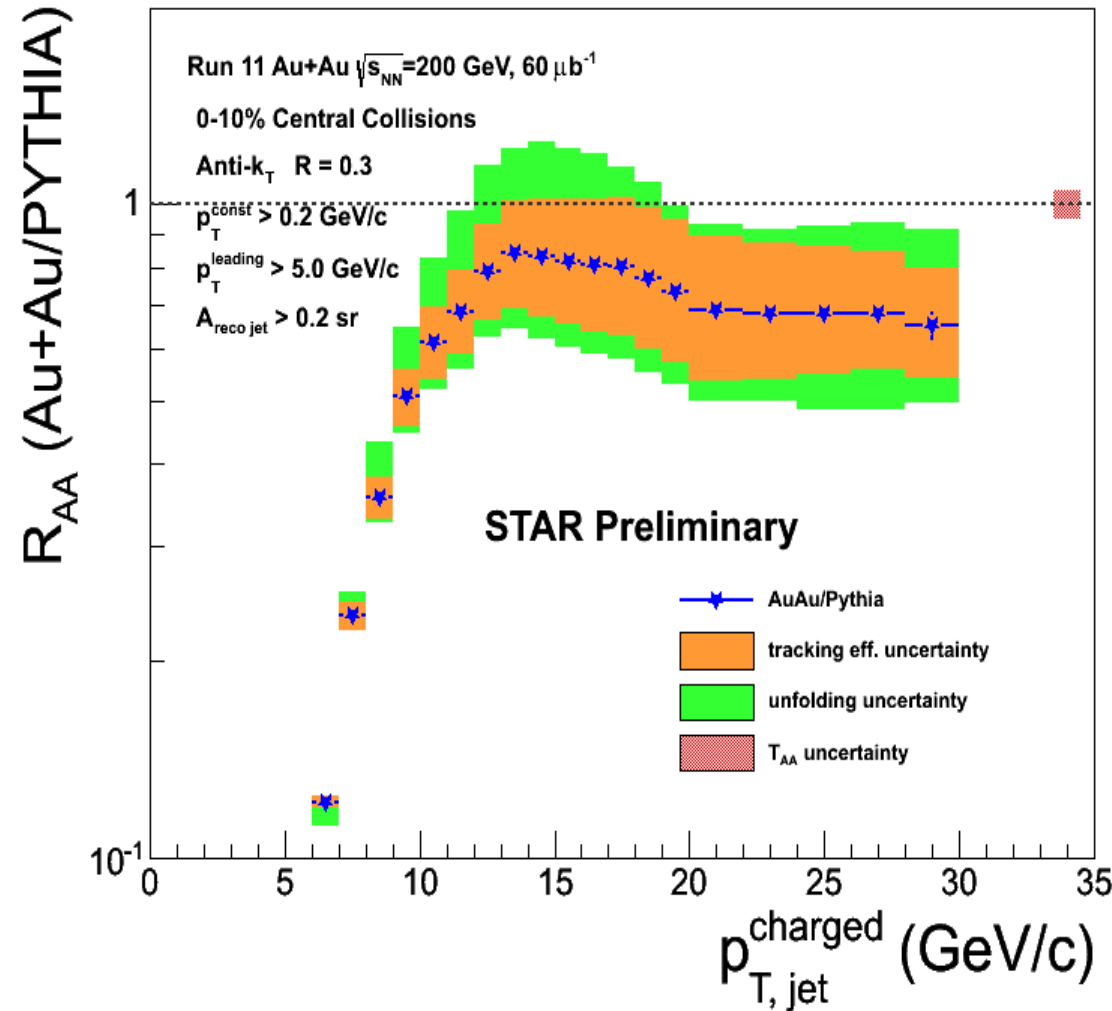
Full jets, biased p+p reference



[ALICE: arXiv:1304.5945]

Comparison with ALICE results

Charged jets



- R_{AA} at RHIC appears to be larger than at the LHC

M. Verweij (ALICE)
 NPA 904-905 (2013) 1015c

Conclusions

New analysis of inclusive charged jet R_{AA} in central Au+Au collisions at $\sqrt{s} = 200\text{GeV}$, utilizing the high statistics and high quality Run11 data (so far only 15% of total int. luminosity)

Dominant instrumental uncertainty is tracking efficiency (JES uncert. $\sim 5\%$)

Charged jet R_{AA} :

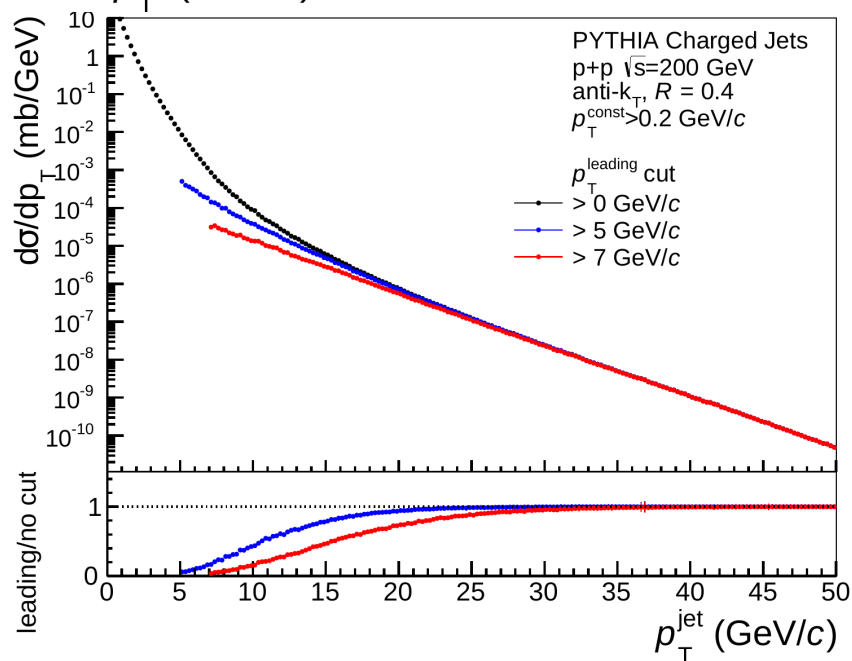
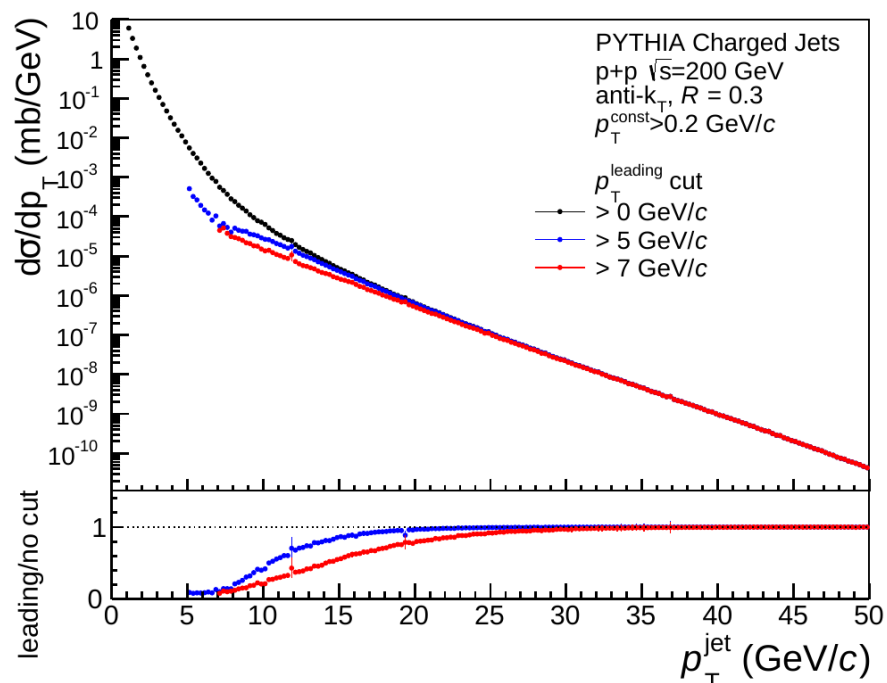
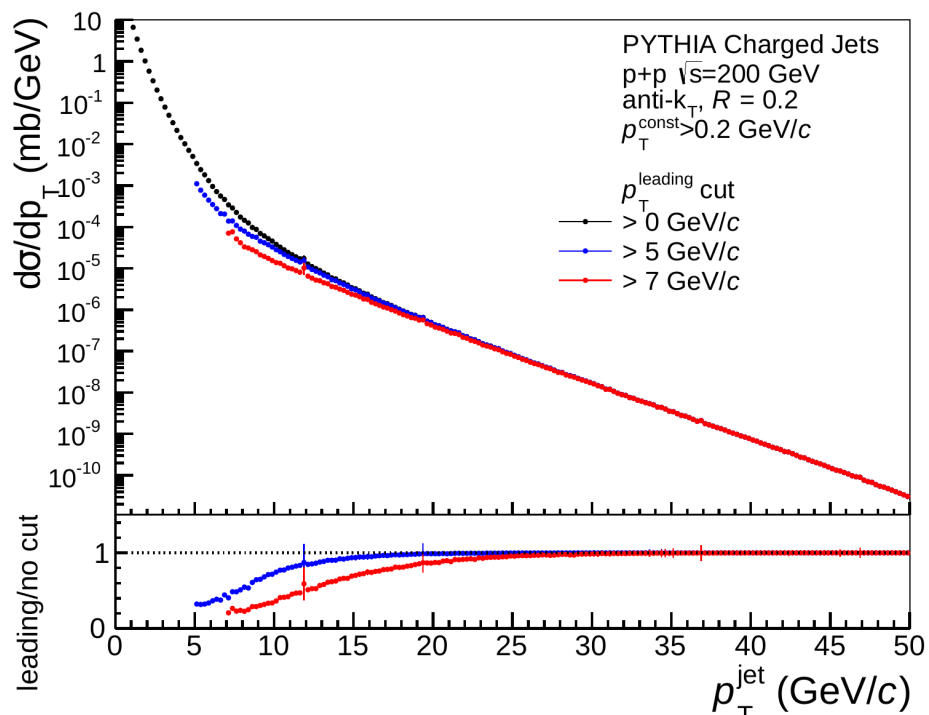
- $R_{AA} \gtrsim 0.5$
- R_{AA} appears to be larger at RHIC (25-30GeV/c) than at the LHC (30-40GeV/c)
- growth in central value with increasing R ?
- (need to assess correlated and uncorrelated uncert. for definitive statement)

Next steps:

- Cross check with other unfolding methods (Singular Value Decomposition)
- Full jets - incorporating BEMC, use triggered data

BACKUP SLIDES

p_T^{leading} cut bias



Tracking Efficiency Correction

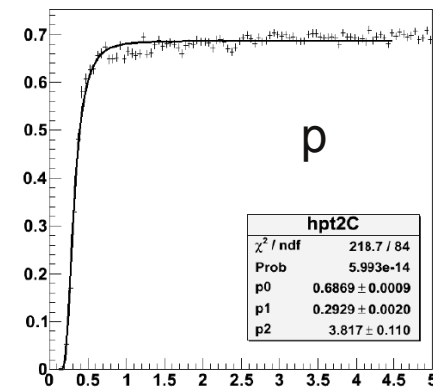
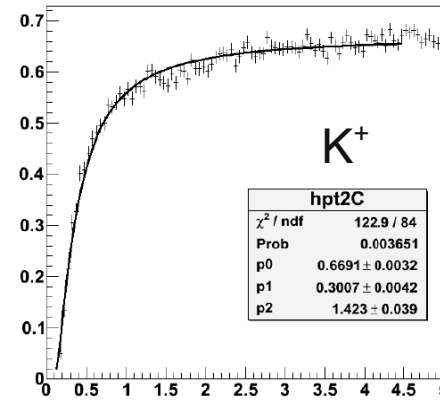
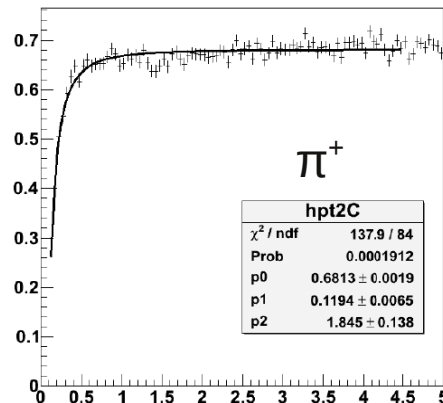
charged hadron reconstruction efficiency:

- from embedding of protons, pions and kaons (Run11)

$$eff_{hadron} = \frac{\sum_i sp_i eff_i}{\sum_j sp_j}$$

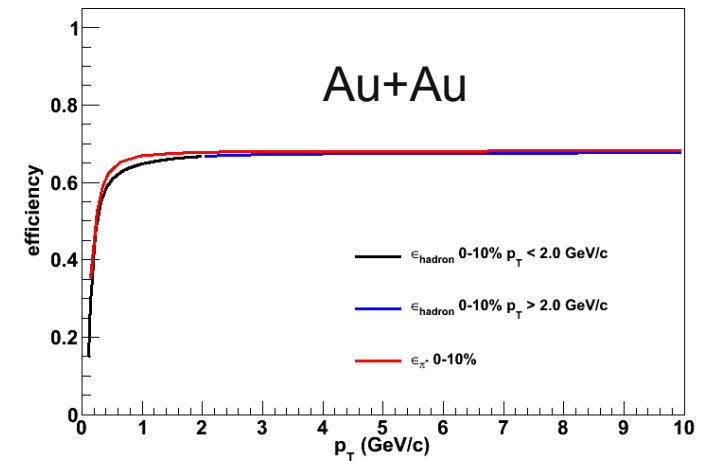
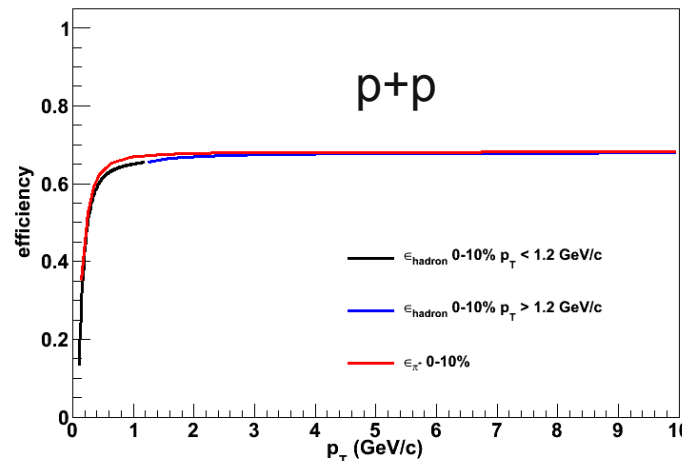
where sp_i is the efficiency corrected single particle spectra and eff_i is the single particle efficiency

$$eff_i = p0 * \exp(-\text{pow}(p1/pT), p2)$$



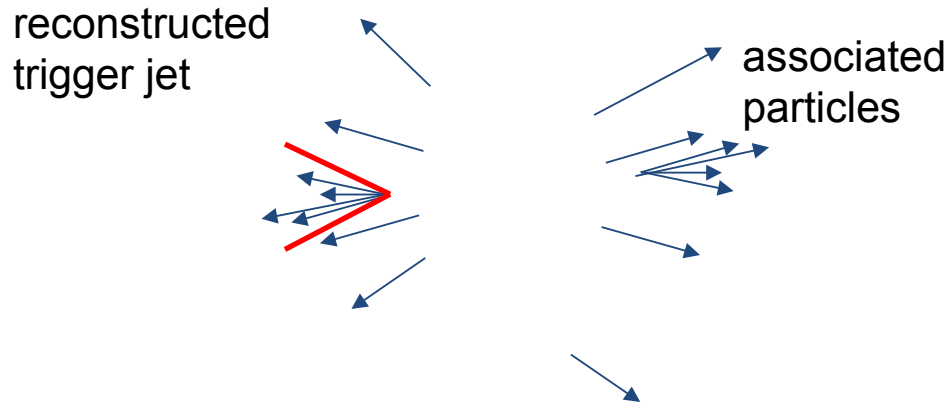
pion/proton/kaon ratios:

- p+p
- Au+Au



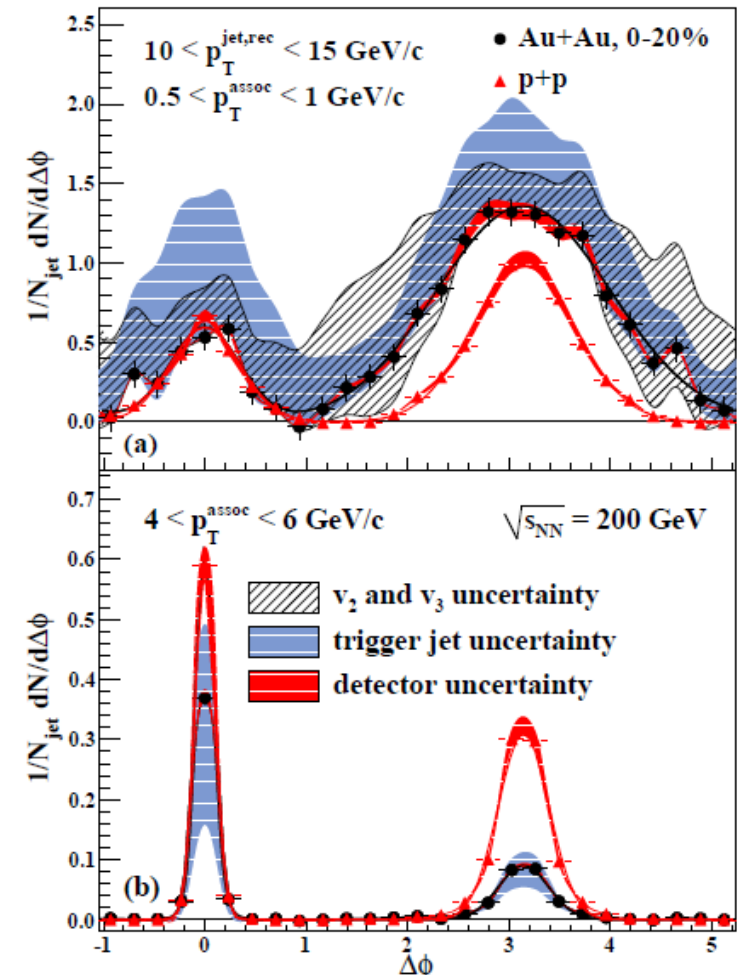
- do unfolding for both scenarios separately and compare results
- both eff. are relatively similar => main source of error will come from embedding uncertainties

Jet-hadron correlations



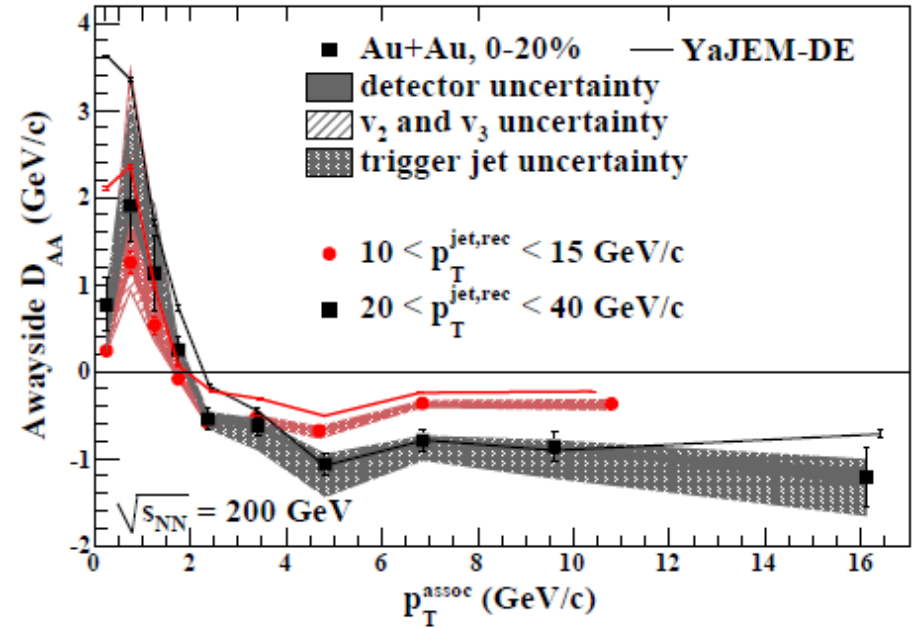
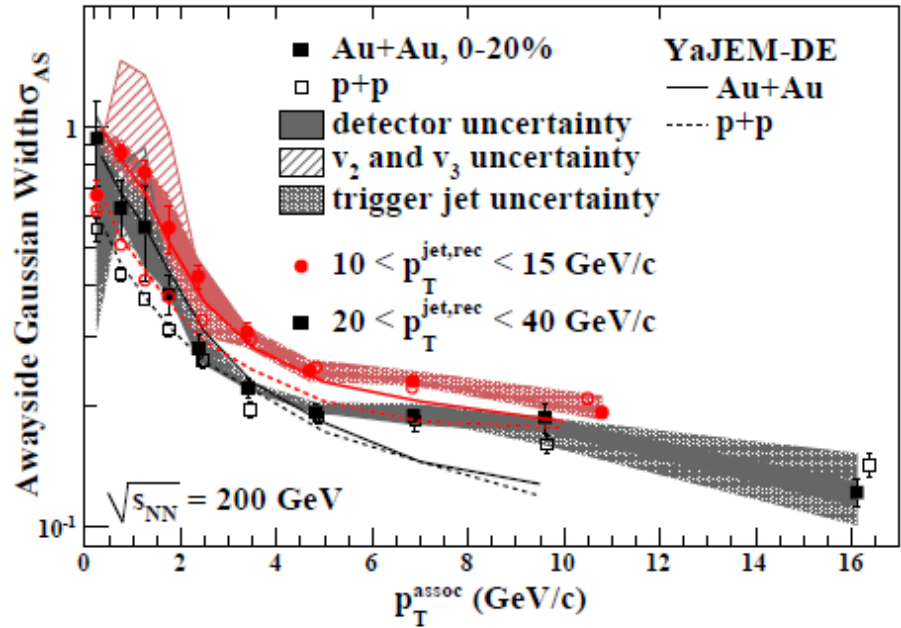
larger kinematic reach compared to di-hadron correlations

- jet reconstruction: anti- k_T algorithm ($R = 0.4$) with a trigger tower in BEMC with $E_T > 6$ GeV and constituent tracks/towers with $p_T(E_T) > 2$ GeV to reduce background fluctuations (NOTE: “a highly biased jet population”)
- recoil (away-side) jet fragmentation is unbiased



Quantifying properties of jet-hadron correlations

STAR, arXiv: 1302.6184 ; YaJEM-DE, T. Renk, PRC 87 (2013) 2, 024905



Away-side Gaussian width:
 suggestive of medium induced broadening
BUT!
 highly depends on $v_3 \rightarrow$ further information
 on v_{2jet} and v_{3jet} needed

Away-side energy balance:

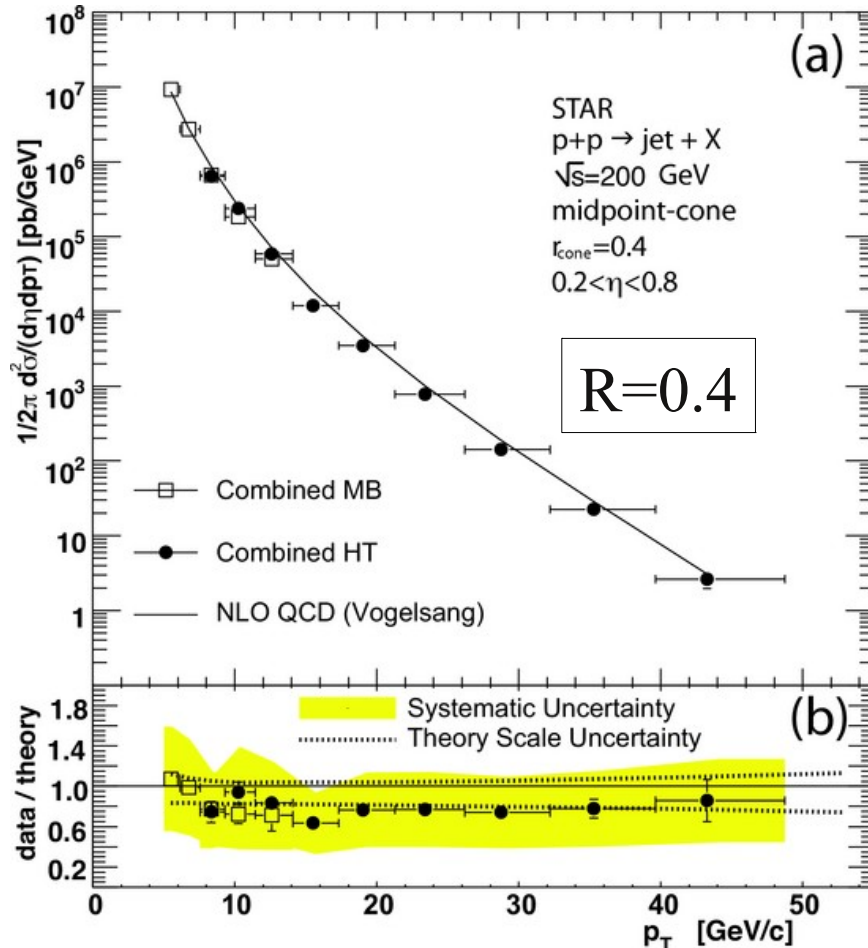
$$D_{AA}(p_T^{assoc}) = Y_{AuAu}(p_T^{assoc}) \cdot \langle p_T^{assoc} \rangle_{AuAu} - Y_{pp}(p_T^{assoc}) \cdot \langle p_T^{assoc} \rangle_{pp}$$

Suppression of high- p_T associated hadron yield \sim balanced by enhancement at low- p_T .

The redistribution of energy quantitatively consistent with YaJEM-DE model of in-medium shower modification.

Estimate of jet yields in Run11 Au+Au data

p+p jet spectrum (from Run 3 + 4)



[STAR, Phys. Rev. Lett. 97 (2006) 252001]

Run 11 Au+Au integrated luminosity:
 $\sim 2.8 \text{nb}^{-1}$

Estimate jet production yield
 (i.e. $R_{AA}=1$):

$$\sim T_{AA} \cdot \frac{d\sigma_{pp}^{\text{jet}}}{dp_T d\eta}$$

10% central Au+Au in Run11:
 We expect $\sim 2\text{K}$ jets with $p_T > 50$ GeV/c
 (full jets with BEMC).

Our analysis:

charged-only jets, no trigger, 40 million minimum bias events ($\sim 60 \mu\text{b}^{-1}$)

$$R_{AA} = \frac{1/N_{events} \frac{dN_{AuAu}}{dp_T d\eta}}{T_{AA} \frac{d\sigma_{pp}}{dp_T d\eta}}$$

$$T_{AA} = \frac{\langle N_{bin} \rangle}{\sigma_{inelastic}} \simeq 22.2 \pm 1 \text{ mb}^{-1} \quad \text{for 0-10\% Au+Au}$$