Measurement of global spin alignment of ϕ and K^{*0} vector mesons at RHIC

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The 7th International Conference on Chirality, Vorticity and Magnetic Field in Heavy Ion Collisions

July 15-19, 2023, Beijing











Outline

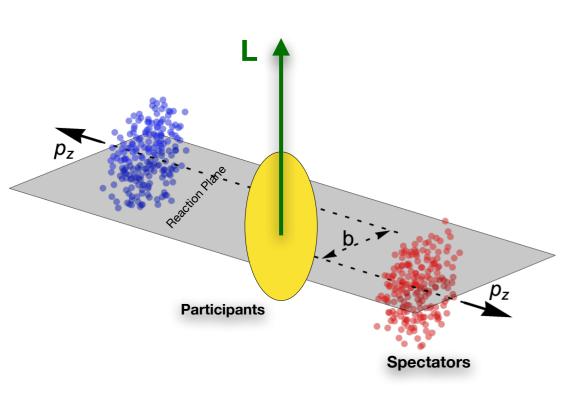
- Introduction
- STAR detector and analysis details
- Results from RHIC Beam Energy Scan
- Summary



Introduction



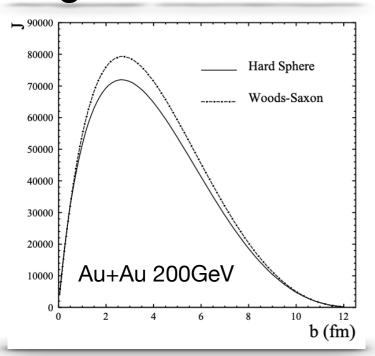
Initial condition in Heavy ion collision



In non-central heavy-ion collisions

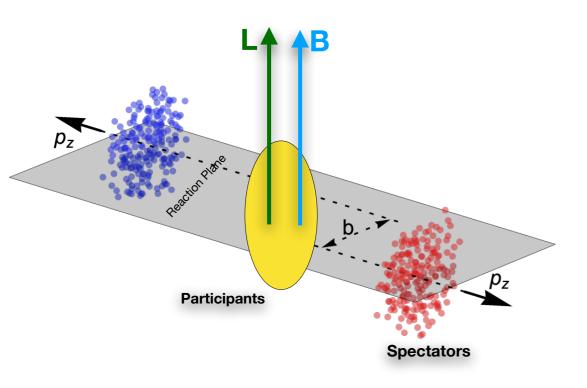
- A large orbital angular momentum (OAM) imparted into the system $L = r \times p \sim bA\sqrt{s_{\rm NN}} \sim 10^4~\hbar$
- Part of OAM transferred to QGP can polarize quarks and anti-quarks due to "spin-orbit" interaction.

Angular momentum





Initial condition in Heavy ion collision



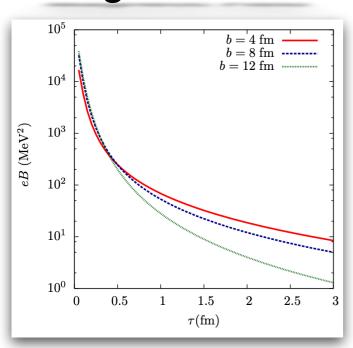
In non-central heavy-ion collisions

 Initial strong magnetic field (B) is expected

$$eB \sim m_{\pi}^2 \sim 10^{18}$$
 Gauss

 Such strong B field can also polarize quarks. Can induce different spin polarization for quarks and anti-quarks with different magnetic moments

Magnetic Field



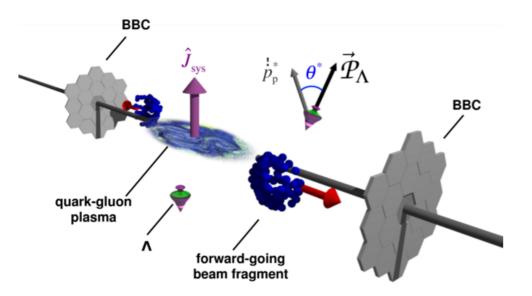
Through polarization measurements

- 1) Study roles of external rotation and magnetic field
- 2) Response of QCD matter under external strong fields



Global spin polarization of Λ

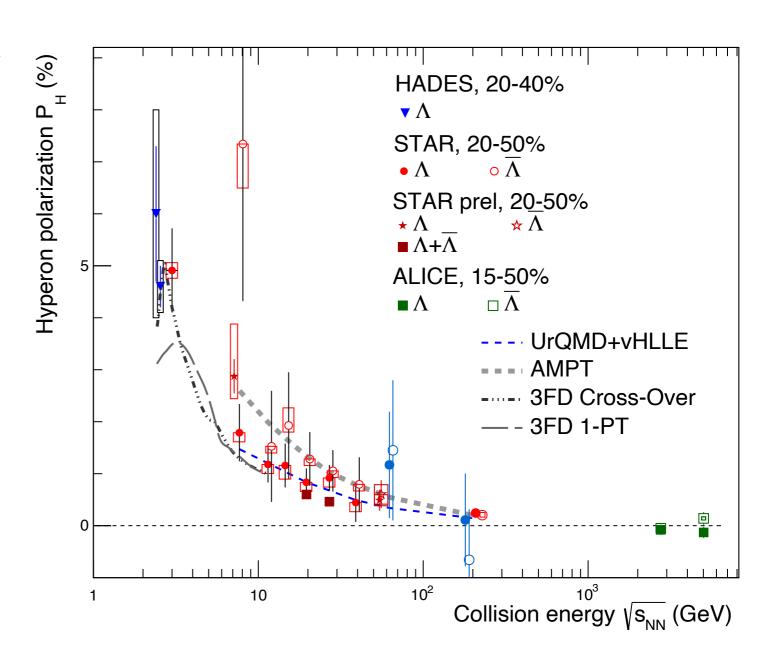
• Λ : parity violating weak decay and hence "self analyzing"



$$P_{\Lambda} = \frac{8}{\pi \alpha_{\Lambda}} \frac{\langle \sin(\Psi_1 - \phi_p^*) \rangle}{\text{Res}(\Psi_1)}$$

Thermal vorticity: $\omega = k_B T (P_{\Lambda} + P_{\bar{\Lambda}})/\hbar$

$$\omega \sim (9 \pm 1) \times 10^{21} \, s^{-1}$$



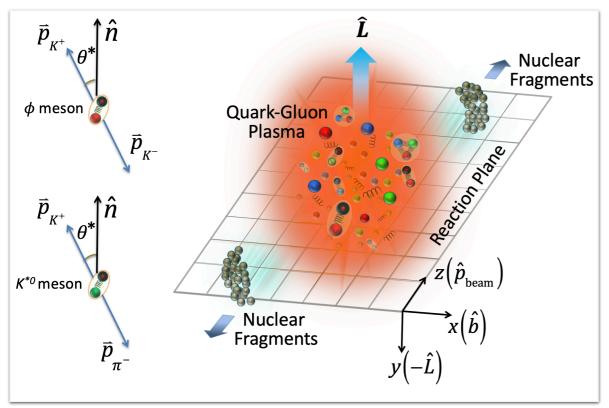
Presence of extreme vorticity induced by angular momentum

STAR: Phys. Rev. C 76, 024915 (2007); Nature 548, 62 (2017); Phys. Rev. C 101, 044611 (2020); Phys. Rev. C 104, 061901 (2021) ALICE: Phys. Rev. C 101, 044611 (2020); HADES: Phys. Lett. B 137506 (2022)



Spin alignment of vector mesons

 $\hat{n} \rightarrow \text{quantization axis}$



$$\rho^{V} = \begin{pmatrix} \rho_{11} & \rho_{10} & \rho_{1-1} \\ \rho_{01} & \rho_{00} & \rho_{0-1} \\ \rho_{-11} & \rho_{-10} & \rho_{-1-1} \end{pmatrix}$$

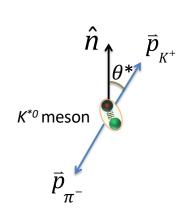
- Vector mesons (K* and ϕ): Spin 1-, Strong decay
- One can measure 00th component of spin density matrix from angular distribution of daughter

$$\frac{dN}{d\cos\theta^*} = N_0 \times \left((1 - \rho_{00}) + (3\rho_{00} - 1)\cos^2\theta^* \right)$$

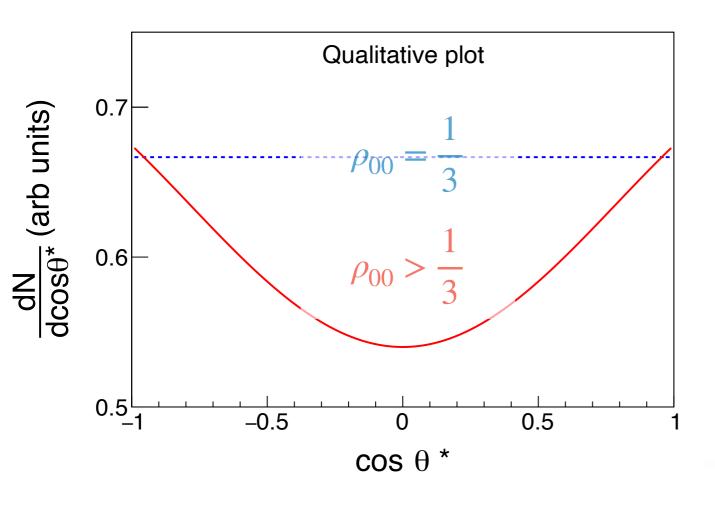
 $ho_{00}
ightarrow ext{Probability of finding a vector meson in spin state "0" out of (1, 0, -1)}$



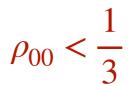
Angular distribution of daughter



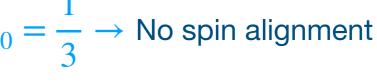
Qualitative expectation of angular distribution



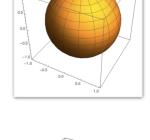
$$\rho_{00} = \frac{1}{3} \rightarrow \text{No spin alignment}$$

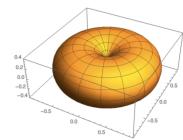


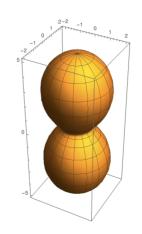
$$\rho_{00} > \frac{1}{3}$$



Spin alignment









Vector mesons vs hyperons

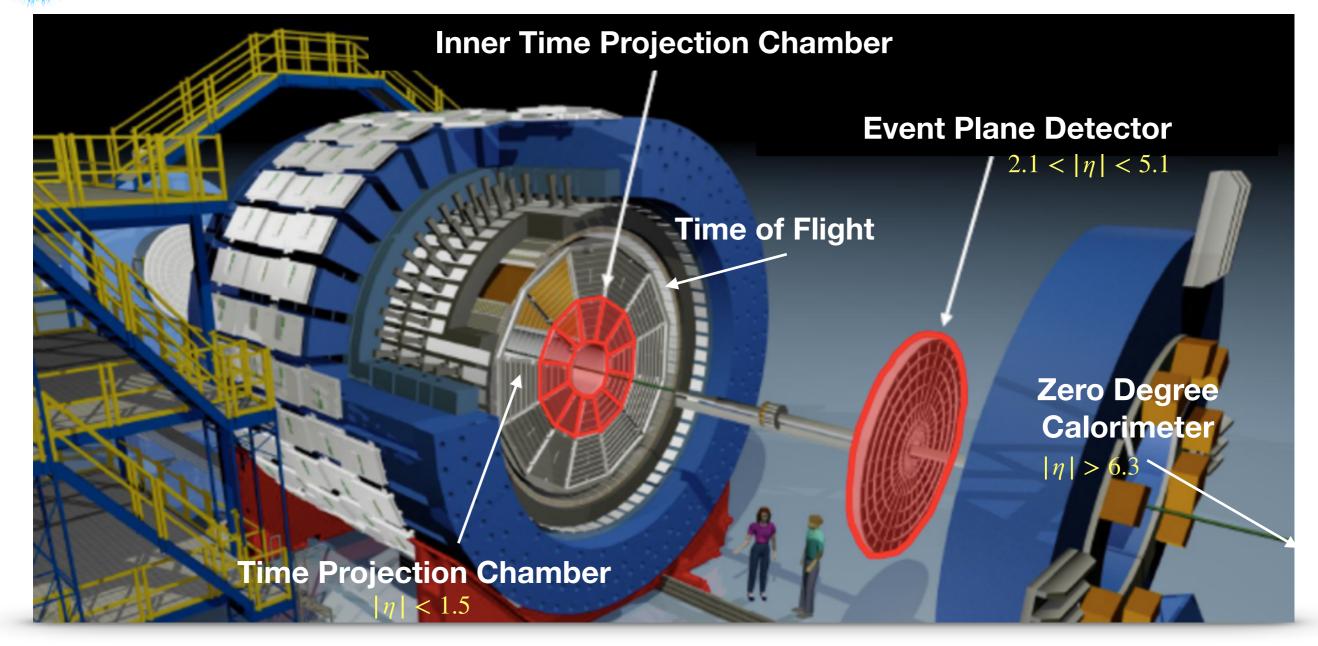
Species	Quark content	J Р	Decay	Feeddown	Direction of Ang. Mom.
K *0	ds	1-	Kπ (strong decay)	negligible (primordial)	not reqd. (Ψ ₂ , Ψ ₁)
φ	ss	1-	KK (strong decay)	negligible (primordial)	not reqd. (Ψ ₂ , Ψ ₁)
^	uds	1/2+	pπ (weak decay)	substantial	required (Ψ ₁)



STAR detector and analysis details



STAR detector



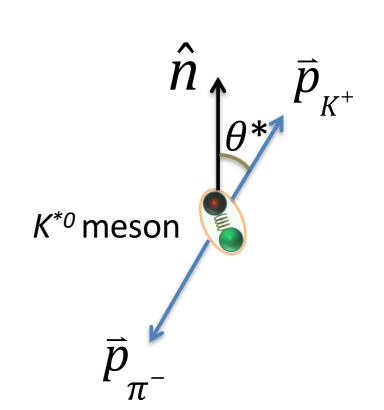
- Uniform acceptance, full azimuthal coverage, excellent PID capability
- <u>TPC</u>: tracking, centrality
- TPC, ZDC, EPD: event plane
- <u>TPC+TOF</u>: particle identification

Analysis Procedure

ρ₀₀ calculation from angular distribution of decay daughters:

- Yield of vector mesons is extracted in $\cos \theta^*$ bins
- Observed ρ_{00} obs is calculated from fitting the yield with function:

$$\frac{dN}{d(\cos\theta^*)} = N_0 \times [(1 - \rho_{00}) + (3\rho_{00} - 1)\cos^2\theta^*]$$



- Observed ρ₀₀ obs is corrected for detector acceptance and efficiency using detector simulation framework
- Observed ρ_{00} obs is corrected for event plane resolution

$$\rho_{00} - \frac{1}{3} = \frac{4}{1 + 3R} (\rho_{00}^{\text{obs}} - \frac{1}{3})$$
 Tang et.al. Phys. Rev. C 98, 044907 (2018)



Results and discussions

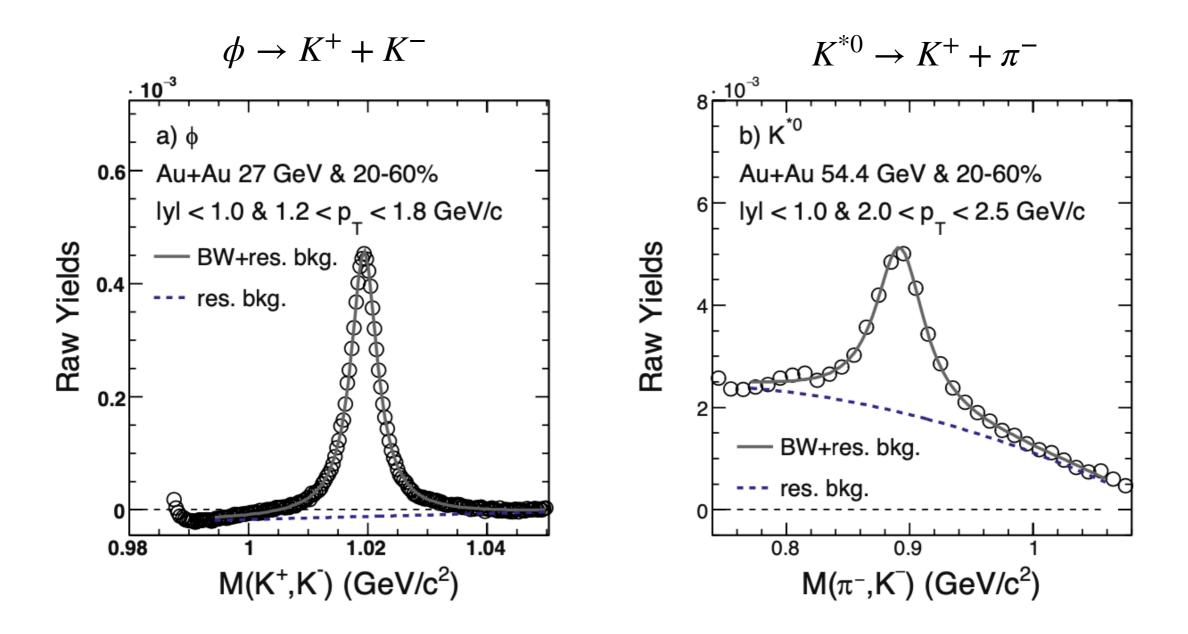


Dataset and analysis details

Collision system	Au+Au, Isobar (Ru+Ru and Zr+Zr)	
Beam energy	11.5, 14.5, 19.6, 27, 39, 54.4, 62.4 and 200 GeV	
Trigger	MInimum bias trigger	
Pair rapidity	y < 1.0	
Background	Mixed event and track rotation method	
Polarization axis	Perpendicular to TPC 2 nd order EP and ZDC 1 st order EP	
Consistency check	3D-random plane	



Invariant mass reconstruction

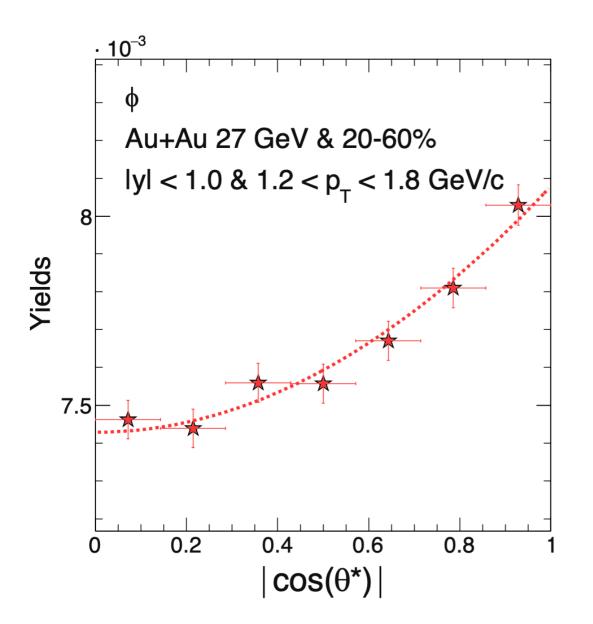


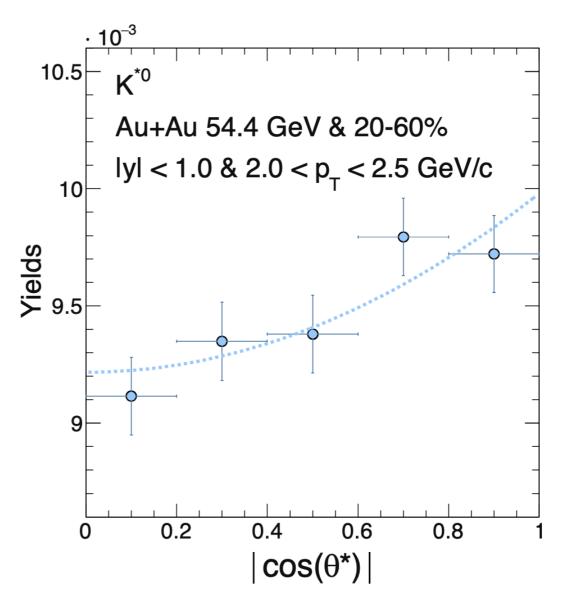
Invariant mass is reconstructed using mixed event and track-rotation methods Invariant mass signal obtained in $\cos\,\theta^*$ bins



Corrected yield vs. $\cos \theta^*$



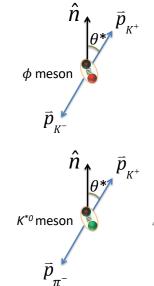




Yields in $\cos \theta^*$ bins are fitted with

$$\frac{dN}{d\cos\theta^*} = N_0 \times \left((1 - \rho_{00,\text{obs}}) + (3\rho_{00,\text{obs}} - 1)\cos^2\theta^* \right)$$

Apply event plane resolution correction: $\rho_{00} - \frac{1}{3} = \left(\rho_{00,\text{obs}} - \frac{1}{3}\right) \frac{4}{1 + 3R}$ 16





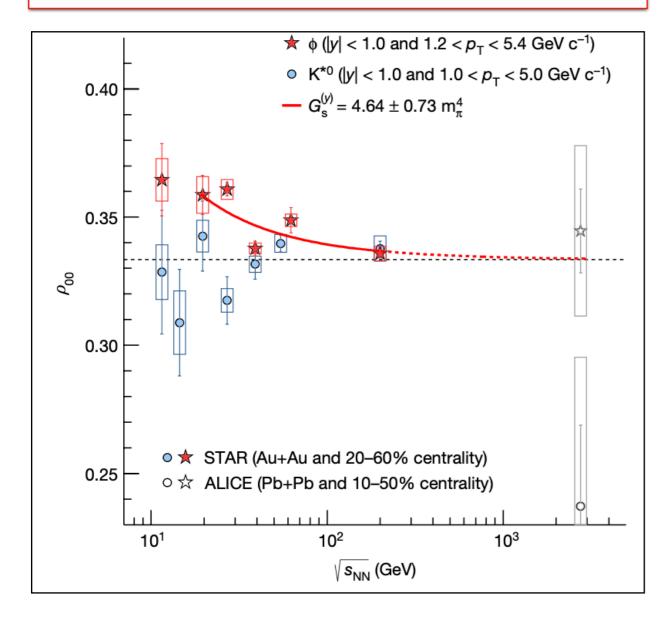
Spin alignment from RHIC BES

Article Published: 18 January 2023

Pattern of global spin alignment of ϕ and K^{*0} mesons in heavy-ion collisions

STAR Collaboration

Nature 614, 244–248 (2023) Cite this article



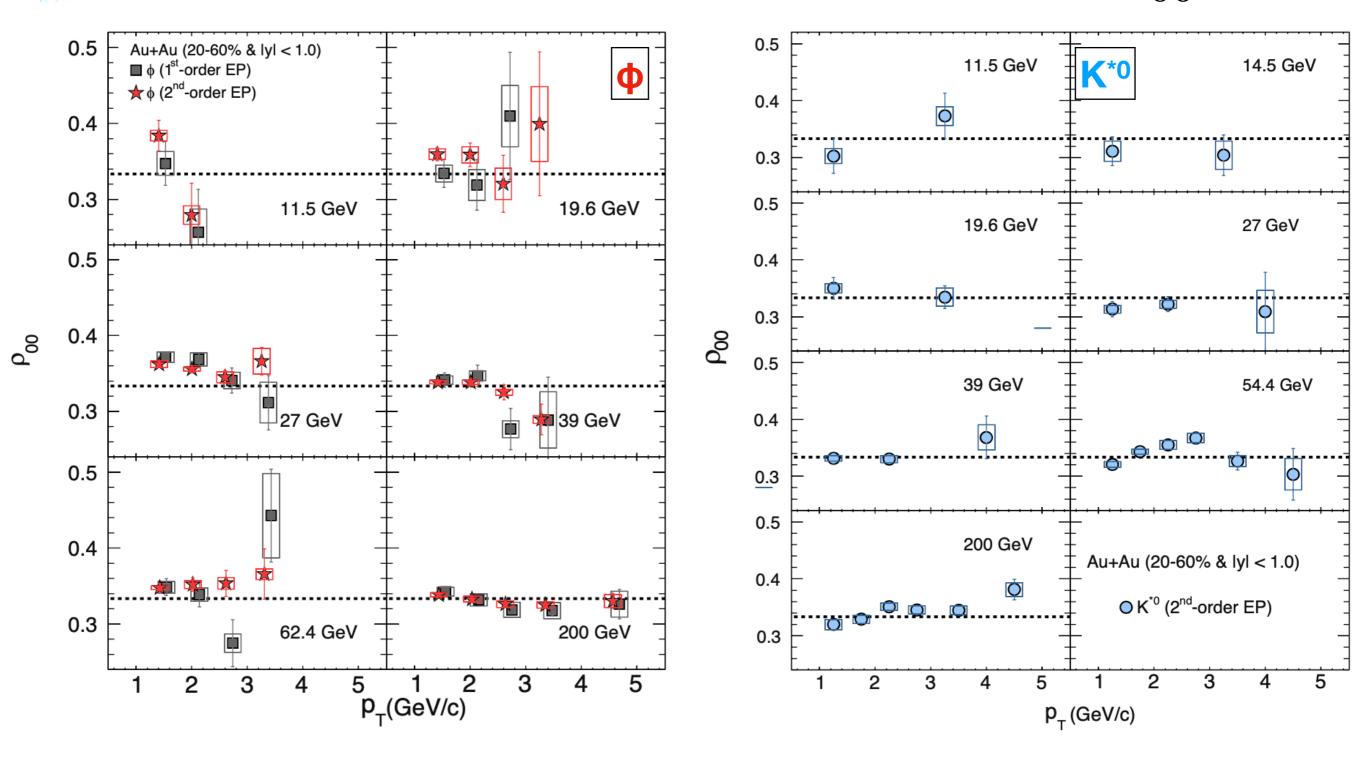
For 20-60%:

- For $\sqrt{s_{\rm NN}}$ <= 62.4 GeV:
- Φ ρ_{00} = 0.3512 \pm 0.0017 (stat.) \pm 0.0017 (sys.) ρ_{00} > 1/3 with 7.4σ
- For $\sqrt{s_{\rm NN}}$ <= 54.4 GeV:
- $^{\circ}$ K^{*0} ρ_{00} = 0.3356 ± 0.0034 (stat.) ± 0.0043 (sys.)

 $\rho_{00} \sim 1/3$



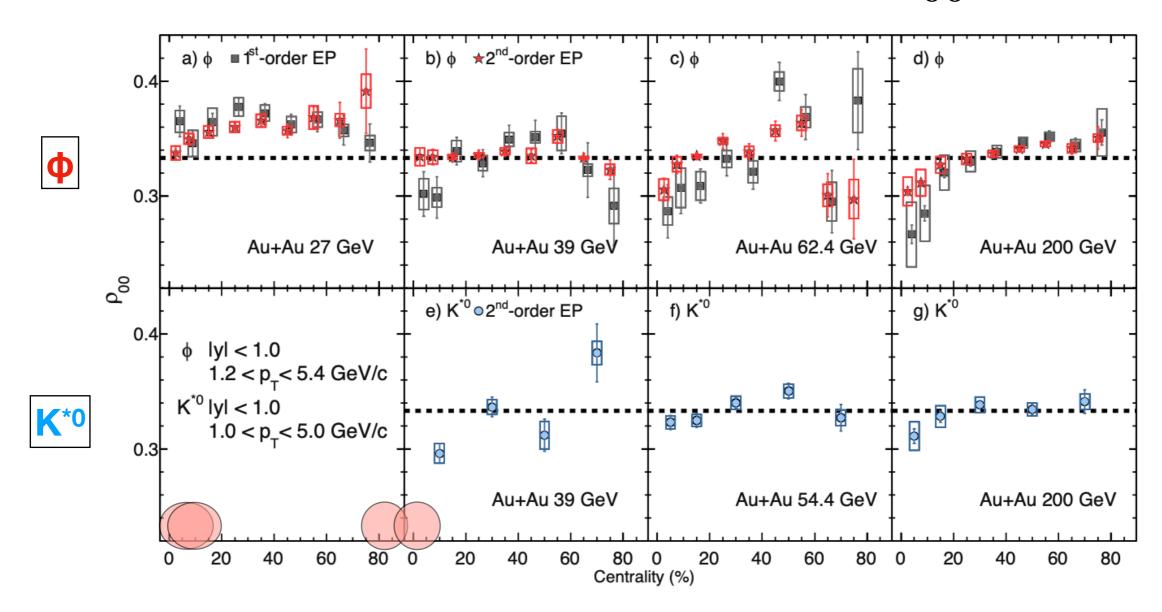
Transverse momentum dependence of ho_{00}



For 20-60%: non-trivial p⊤ dependence



Centrality dependence of ho_{00}



- For mid-central and peripheral:
- Φ , K^{*0} $\rho_{00} > \sim 1/3$

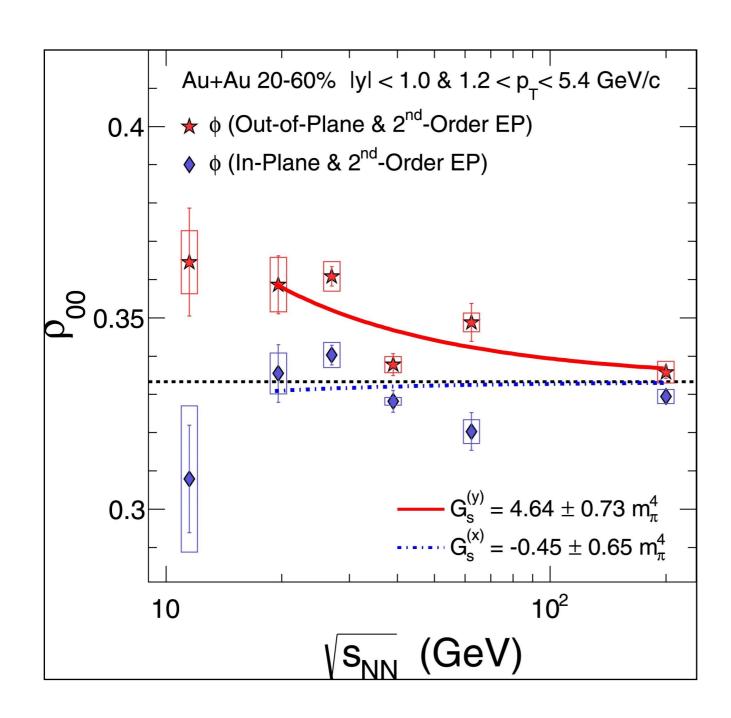
- For central at 200 GeV:
- φ, K*0 ρ₀₀ < 1/3
 Local spin alignment^[1]
 or, helicity contribution^[2]

^{[1].} Xia et al, Phys Lett B 817, 136325 (2021)

^{[2].} Gao, Phys Rev D 104, 076016 (2021)



ϕ meson ρ_{00} : in-plane vs out-plane



- For mid-central:
- ϕ ρ_{00} : in-plane < out-of-plane
- Effect from elliptic flow



Physics Mechanisms	(ρ ₀₀)
c _∧ : Quark coalescence vorticity & magnetic field ^[1]	< 1/3 (Negative ~ 10 ⁻⁵)
c _ε : Vorticity tensor ^[1]	< 1/3 (Negative ~ 10 ⁻⁴)
c _E : Electric field ^[2]	> 1/3 (Positive ~ 10 ⁻⁵)
c _F :Fragmentation ^[3]	> or, < 1/3 (~ 10 ⁻⁵)
c _L :Local spin alignment and helicity ^[4]	< 1/3
с т:Turbulent color field ^[5]	< 1/3
c _φ : Vector meson strong force field ^[6]	> 1/3 (Can accomodate large positive signal)

$$\rho_{00}(\phi) \approx \frac{1}{3} + c_{\Lambda} + c_{\epsilon} + c_{E} + c_{F} + c_{L} + c_{T} + c_{\phi};$$

- [1]. Liang et., al., Phys. Lett. B 629, (2005);
 Yang et., al., Phys. Rev. C 97, 034917 (2018);
 Xia et., al., Phys. Lett. B 817, 136325 (2021);
 Beccattini et., al., Phys. Rev. C 88, 034905 (2013)
 [2]. Sheng et., al., Phys. Rev. D 101, 096005 (2020);
- [2]. Sheng et., al., Phys. Rev. D 101, 096005 (2020); Yang et., al., Phys. Rev. C 97, 034917 (2018)
- [3]. Liang et., al., Phys. Lett. B 629, (2005)
- [4]. Xia et., al., Phys. Lett. B 817, 136325 (2021); Guo, Phys. Rev. D 104, 076016 (2021)
- [5]. Muller et., al., Phys. Rev. D 105, L011901 (2022)
- [6]. Sheng et., al., Phys. Rev. D 101, 096005 (2020); Sheng et., al., Phys. Rev. D 102, 056013 (2020)



_		
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c _L :Local spin alignment and helicity ^[4]	< 1/3	
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c _Φ : Vector meson strong force field ^[6]	> 1/3 (Can accomodate large positive signal)	

$$\rho_{00}(\phi) \approx \frac{1}{3} + c_{\Lambda} + c_{\epsilon} + c_{E} + c_{F} + c_{L} + c_{T} + c_{\phi};$$

$$c_{\Lambda} = \frac{4}{9} \langle P_{\Lambda} P_{\bar{\Lambda}} \rangle = -\frac{1}{9} \langle \omega_y^2 \rangle + \frac{Q_s^2}{9 m_s^2 T_{\text{eff}^2}} \langle B_y^2 \rangle$$

(Contribution from vorticity and magnetic field are constrained by Λ polarization)

Conventional mechanisms that can capture global spin polarization of Λ and its energy dependence can not accommodate large deviations observed in ρ_{00}



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c _∧ : Quark coalescence vorticity & magnetic field ^[1]	< 1/3 (Negative ~ 10 ⁻⁵)
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c _φ : Vector meson strong force field ^[6]	> 1/3 (Can accomodate large positive signal)

$$\rho_{00}(\phi) \approx \frac{1}{3} + c_{\Lambda} + c_{\epsilon} + c_{E} + c_{F} + c_{L} + c_{T} + c_{\phi};$$

$$\rho_{00}^{\phi} = \frac{1}{3} - \frac{4}{9} \langle P_{\Lambda} P_{\bar{\Lambda}} \rangle$$

$$+ \frac{1}{27m_s^2} \langle p^2 \rangle_{\phi} \langle \epsilon_x^2 + \epsilon_z^2 \rangle$$

Conventional mechanisms that can capture global spin polarization of Λ and its energy dependence can not accommodate large deviations observed in ρ_{00}



Physics Mechanisms	(ροο)
c _∧ : Quark coalescence vorticity & magnetic field ^[1]	< 1/3 (Negative ~ 10 ⁻⁵)
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$$\rho_{00}(\phi) \approx \frac{1}{3} + c_{\Lambda} + c_{\epsilon} + c_{E} + c_{E} + c_{F} + c_{L} + c_{T} + c_{\phi};$$

$$\rho^{\phi} = \frac{1}{2} - \frac{4}{2} \langle P, P_{\tau} \rangle$$

$$\rho_{00}^{\phi} = \frac{1}{3} - \frac{4}{9} \langle P_{\Lambda} P_{\bar{\Lambda}} \rangle$$

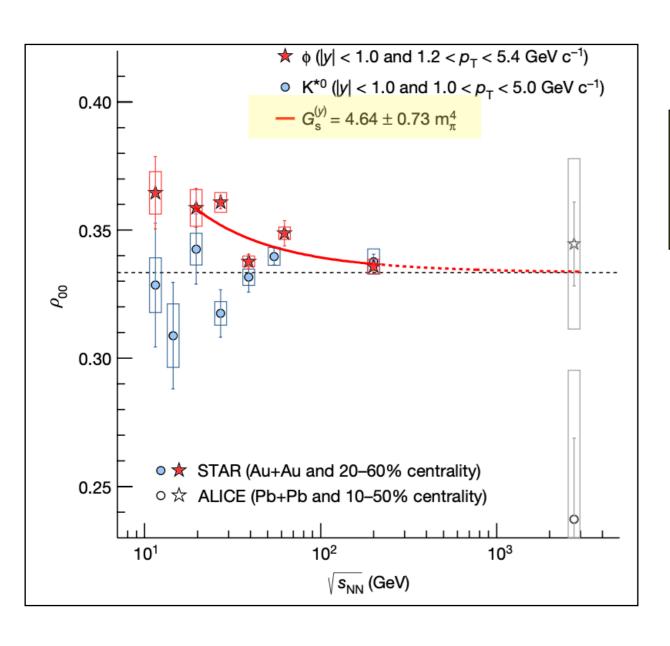
$$+ \frac{1}{27m_s^2} \langle p^2 \rangle_{\phi} \langle \epsilon_x^2 + \epsilon_z^2 \rangle$$

$$- \frac{g_{\phi}^4}{27m_s^4 m_{\phi}^4 T_{\text{eff}}} \langle E_{\phi,z}^2 + E_{\phi,x}^2 \rangle$$

- Like electric charges in motion can generate an EM field, s and \bar{s} quarks in motion can generate an effective ϕ -meson field
- The ϕ -meson field can polarize s and \bar{s} quarks with a large magnitude due to strong interaction (large coupling constant g_{ϕ})



Energy dependence of spin alignment



A model with vector meson strong force field can accommodate the large positive signal of ϕ meson ρ_{00}

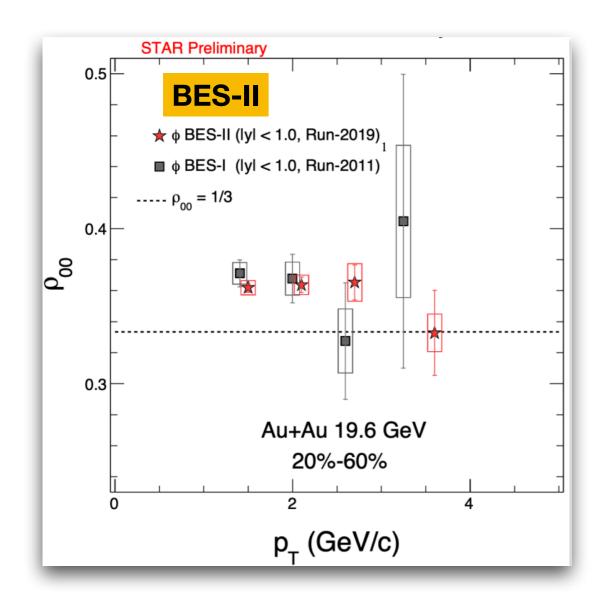
$$\rho_{00}(\phi) \approx \frac{1}{3} + c_{\Lambda} + c_{\epsilon} + c_{E} + c_{\phi};$$

$$c_{\phi} = \frac{G_{s}^{y}}{27m_{s}^{2}T_{\text{eff}}^{2}}$$

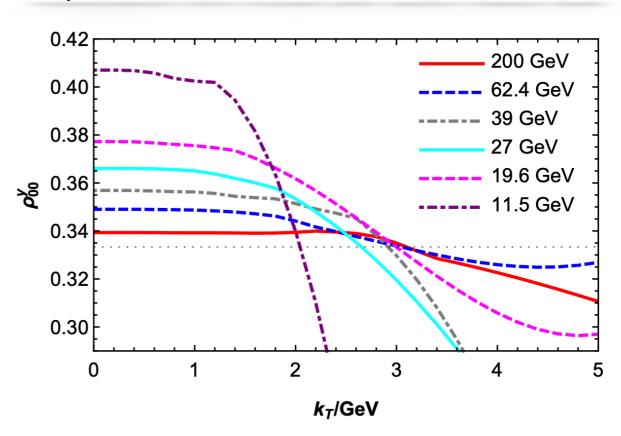
$$G_{s}^{y} = g_{\phi}^{2} \left[3\langle B_{\phi,y}^{2} \rangle - \frac{\langle P^{2} \rangle_{\phi}}{m_{s}^{2}} \langle E_{\phi,z}^{2} + E_{\phi,x}^{2} \rangle \right]$$



Transverse momentum dependence of $\phi \, \rho_{00}$



Expectation from model with vector meson field



Sheng el. al., arXiv: 2205.15689

Sheng el. al., Phys Rev D 101, 096005 (2020)

Sheng el. al., Phys Rev D 102, 056013 (2020)

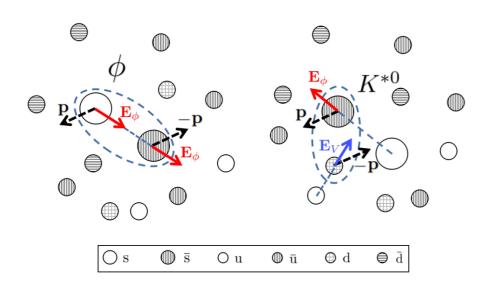
Precision from BES-II:

$$\rho_{00} > \frac{1}{3}$$
 (about $\sim 5\sigma$ significance)



About K* ρ_{00}

Expectation from model with vector meson field



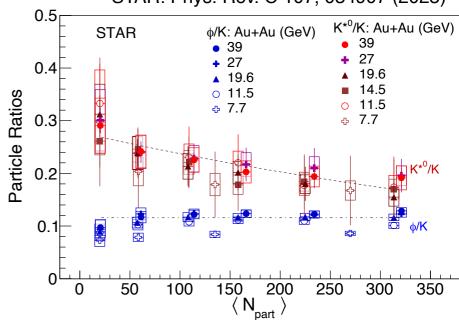
We note that the above arguments are only valid for primary K^{*0} . The life time of K^{*0} is much shorter and the interaction of K^{*0} with the surrounding matter is much stronger than the ϕ meson. This may bring other contributions to $\rho_{00}^{K^*}$ from the interaction of K^{*0} with medium. A caveat is that the above arguments are based on the approximation that different fields do not have large correlation in space as compared with the correlation between the same fields. This seems to work for ρ_{00}^{ϕ} since there are squares of the same vector meson field. However it is not the case for $\rho_{00}^{K^*}$ that all terms of vector meson fields are mixture of differenct fields which are thought to be equally small. In this case, in order to justify the approximation, we may need to evaluate these terms and compare their magnitudes with the negative comtribution from vorticity tensor fields. This is beyond the scope of this paper and will be studied in the future.

Sheng el. al., arXiv: 2205.15689

Sheng el. al., Phys Rev D 101, 096005 (2020) Sheng el. al., Phys Rev D 102, 056013 (2020)

Species	Quarks	J₽	Decay	Lifetime (fm/c)
K*0	ds	1-	Κπ	4
φ	SS	1-	KK	45





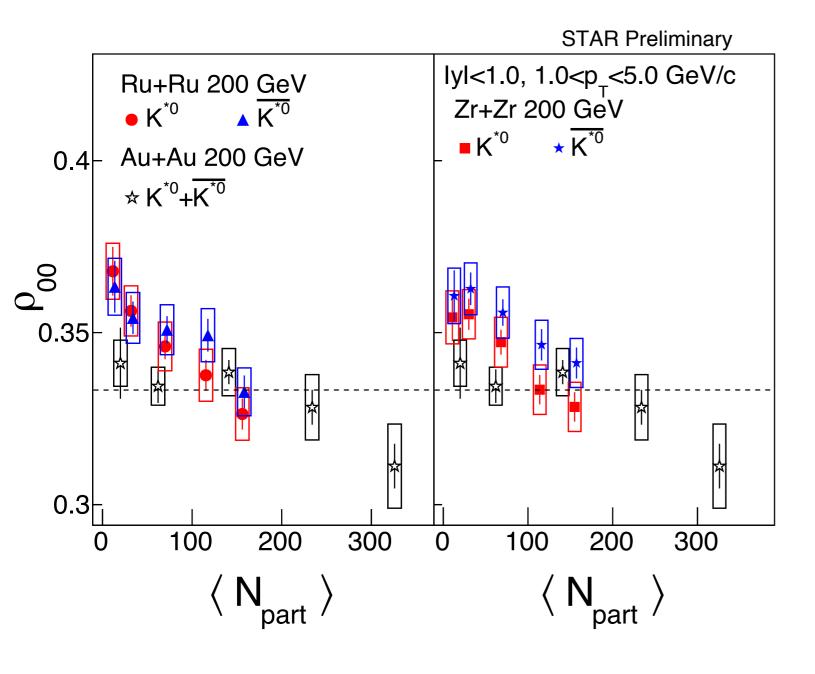
- (K*/K)_{central} < (K*/K)_{peripheral}
- (ϕ/K) ~ centrality independent

Compared to ϕ :

- K* have un-equal constituent quarks, have different response to vector meson field
- K* shorter lifetime, more susceptible to in-medium interactions



$\mathsf{K}^{\star}\,\rho_{00}$ from isobar collisions



- Species dependence:
- $K^{*0} \rho_{00} \sim \text{anti-} K^{*0} \rho_{00}$
- System size dependence:
- ρ_{00} : Au+Au ~ Zr+Zr ~ Ru+Ru
- ρ_{00} ~ scales with $\langle N_{\rm part} \rangle$



Summary

- At RHIC BES:
- For the first time significantly large ϕ meson $\rho_{00} > 1/3$ ($\sim 7.4\sigma$)
- K* meson ρ_{00} ~ 1/3
- Surprising magnitude of $\phi \rho_{00}$: can not be explained by conventional mechanisms of polarization (e.g. driven by vorticity, EM fields ...)
- Many players and new developments:
- Influence of vector meson strong force field (Sheng et. al.)
- Effects from gluon and glasma field (Muller et. al.)
- Contribution from local spin alignment, helicity (Xia et. al., Gao)
- Shear induced tensor polarization (Liu et. al.)

—...

Outlook

Experiment side: High precision measurements from RHIC BES-II and other datasets; ρ_{00} of different vector meson species $(\rho, \omega, J/\Psi)$ at RHIC are ongoing Stay tuned



Article Published: 18 January 2023

Pattern of global spin alignment of ϕ and K^{*0} mesons in heavy-ion collisions

STAR Collaboration

Nature 614, 244–248 (2023) | Cite this article

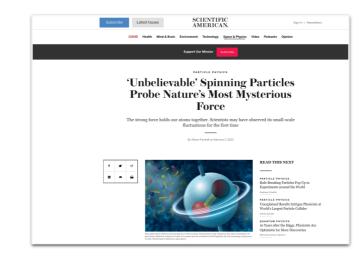
2988 Accesses | 8 Citations | 165 Altmetric | Metrics

Pattern of global spin alignment of φ and K*0 mesons in heavy-ion collisions

STAR Collaboration • M.S. Abdallah (American U., Cairo) et al. (Apr 5, 2022)

Published in: Nature 614 (2023) 7947, 244–248, Nature (2023) • e-Print: 2204.02302 [hep-ph]

Dip pdf links DOI cite datasets claim reference search 227 citations





Thank you for your attention

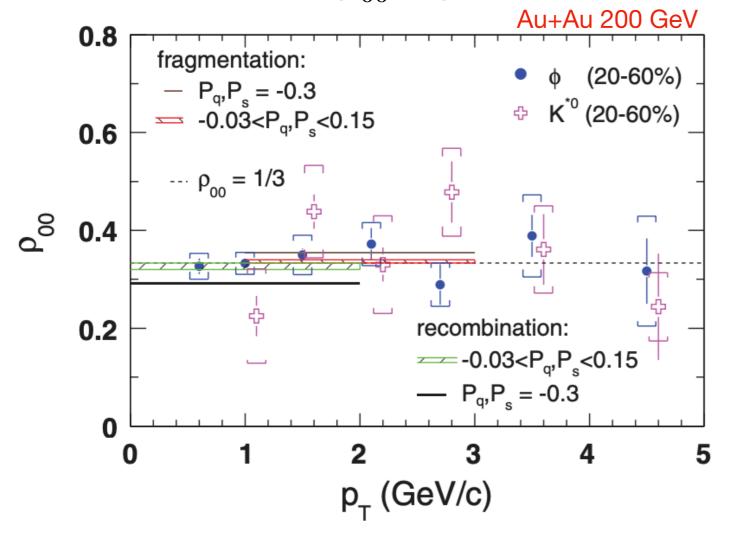


Back up slides



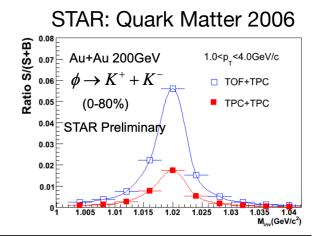
First spin alignment from RHIC

• STAR reported first ho_{00} of ϕ and K* wrt reaction plane in 2008 using ~ 23 million events:



- For both ϕ and K*:
- results were in-conclusive, dominated by large uncertainty

STAR: Phys. Rev. C 77, 061902 (2008)

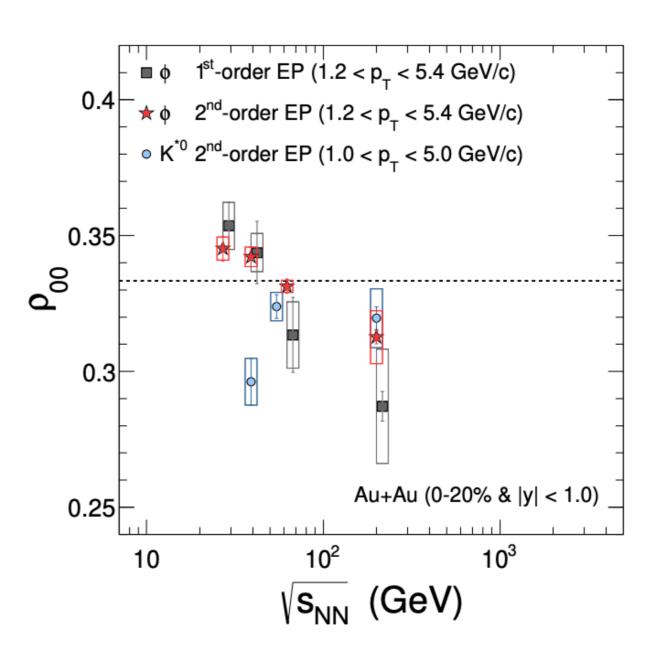


Advantages in measurements from BES

- STAR upgrade: Benefitted from Time of Flight (ToF), better particle identification, improved signal-to-noise ratio
- Analysis method: Improved understanding of detector corrections (efficiency, acceptance and event plane resolutions)

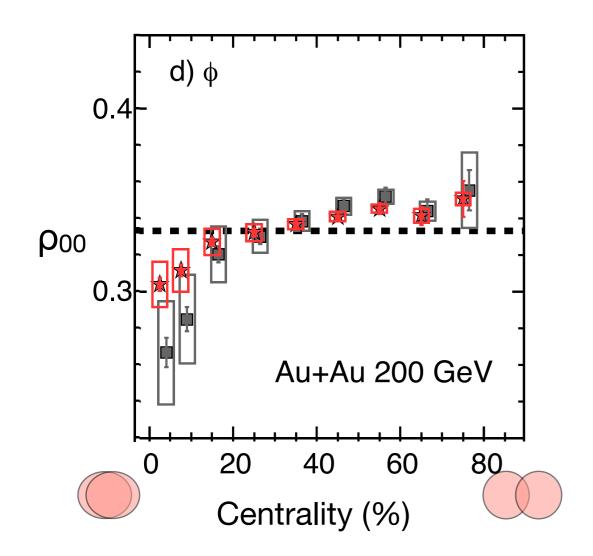


Energy dependence of ϕ and $K^* \, \rho_{00}$ in central collisions





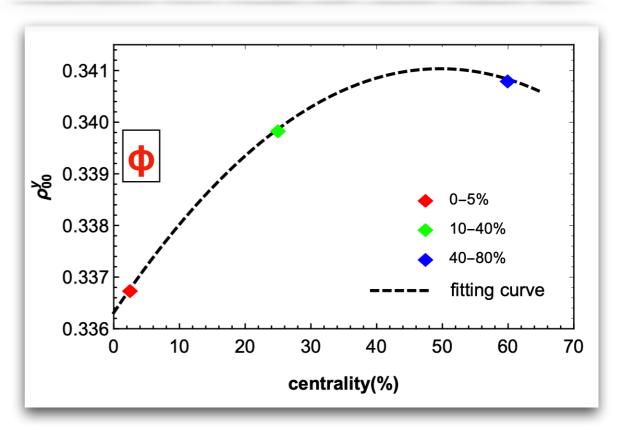
Centrality dependence of $\phi \, \rho_{00}$



Is the contribution from local spin alignment dominant in central collisions and at higher energies?

Xia et., al., Phys Lett B 817, 136325 (2021)

Expectation from model with vector meson field



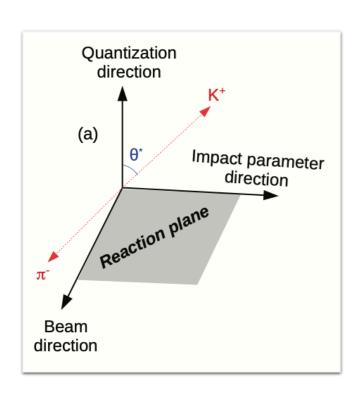
Can accommodate positive deviation in mid-central and peripheral collisions

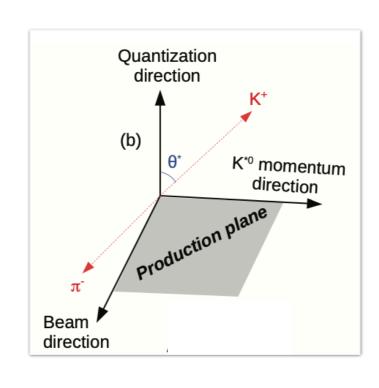
Sheng el. al., arXiv: 2205.15689 Sheng el. al., Phys Rev D 101, 096005 (2020) Sheng el. al., Phys Rev D 102, 056013 (2020)

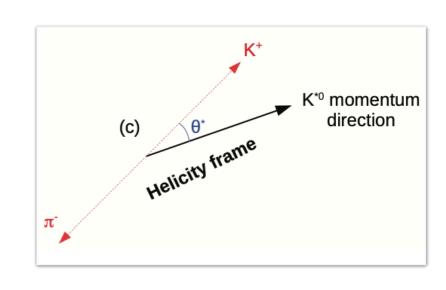


Quantization axis

Choice of Quantization direction (or Polarization axis)







Reaction Plane

Production Plane

Helicity

Experimentally, we can't measure impact parameter direction. We use "event plane" as a proxy for reaction plane



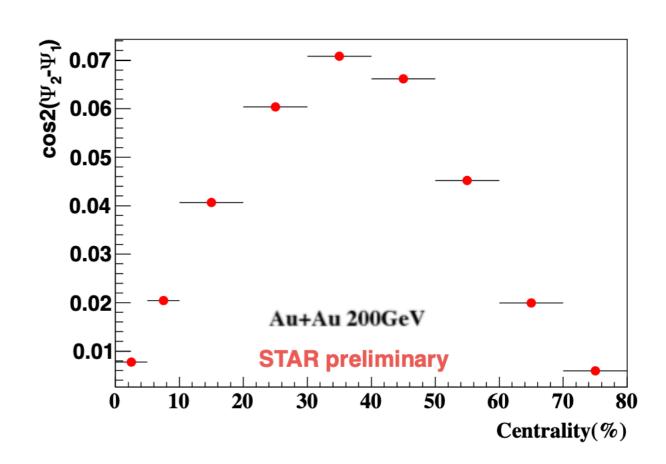
Event plane de-correlation effect

 ho_{00} can be measured using Ψ_2 (2nd order event plane, TPC) and Ψ_1 (1st order event plane, ZDC)

$$\rho_{00,\text{obs}}^{\Psi_2} - \frac{1}{3} = \left(\rho_{00}^{\Psi_2} - \frac{1}{3}\right) \frac{1 + 3R_2}{4}$$

$$\rho_{00,\text{obs}}^{\Psi_2} - \frac{1}{3} = \left(\rho_{00}^{\Psi_1} - \frac{1}{3}\right) \frac{1 + 3D_{12} R_1}{4}$$

$$\rho_{00}^{\Psi_2} - \frac{1}{3} = \left(\rho_{00}^{\Psi_1} - \frac{1}{3}\right) \frac{1 + 3D_{12} R_1}{1 + 3R_2}$$



Event plane de-correlation effect needed to be considered carefully

Tang et.al. Phys. Rev. C 98, 044907 (2018)

STAR

Polarization wrt event plane

Reconstruct event plane (Ψ_n) :

$$Q_{n,y} = \sum_{i} w_{i} \sin(n\phi_{i}); \quad Q_{n,x} = \sum_{i} w_{i} \cos(n\phi_{i})$$

$$\Psi_n = \operatorname{arctan2}(Q_{n,y}, Q_{n,x})/n$$

Unit vectors in reaction plane and event plane are related by:

$$\hat{x} = \cos(\Psi) \ \hat{x}' - \sin(\Psi) \ \hat{y}'; \ \hat{y} = \sin(\Psi) \ \hat{x}' + \cos(\Psi) \ \hat{y}'; \ \hat{z} = \hat{z}'$$



$$\frac{dN}{d\cos(\theta^{*'})} \propto (1 - \rho_{00}^{EP}) + (3\rho_{00}^{EP} - 1)\cos(\theta^{*'})$$



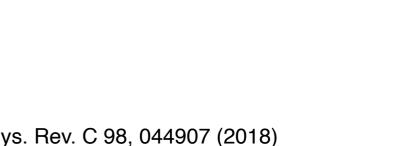
$$\rho_{00}^{EP} = \rho_{00}^{RP} - \frac{1}{2} (3 \ \rho_{00}^{RP} - 1) \sin^2(\Psi)$$



$$\rho_{00}^{RP} - \frac{1}{3} = \left(\rho_{00}^{EP} - \frac{1}{3}\right) \frac{4}{1 + 3R}$$

Tang et.al. Phys. Rev. C 98, 044907 (2018) Mohanty et., al, Mod. Phys. Lett. A, 36, 2130026 (2021)

Ζ



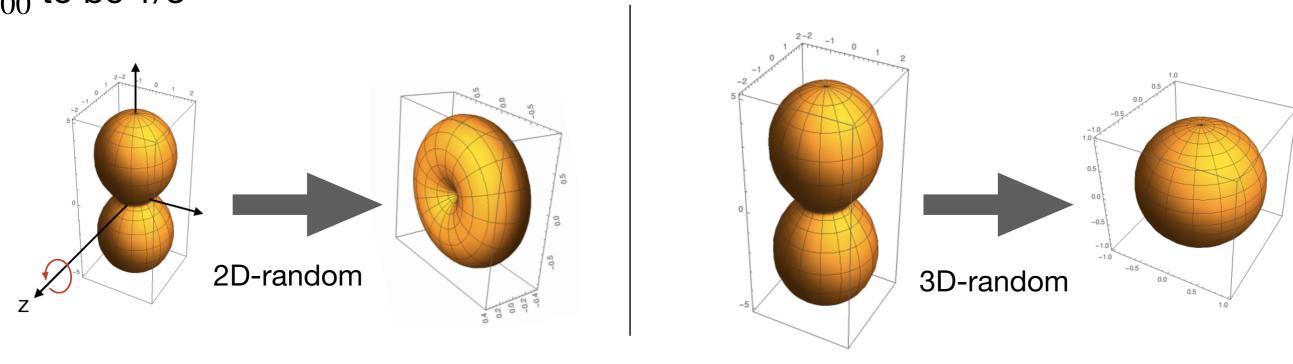
$$\frac{4}{1+3R}$$
 $R \to \text{Event plane resolution}$



Polarization wrt random plane

 ho_{00} can be measured using a random plane; Naive expectation $ho_{00}^{\rm random} \sim \frac{1}{3}$ (e.g. using random number but depends on how it is randomized)

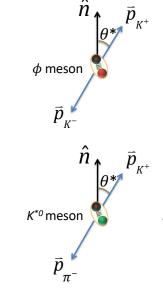
If the quantization axis is randomized in 2D-plane (x, y), it does not guarantee the ρ_{00} to be 1/3

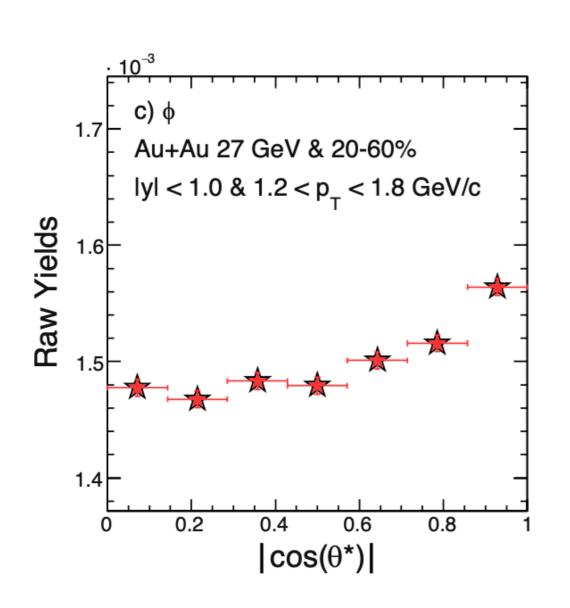


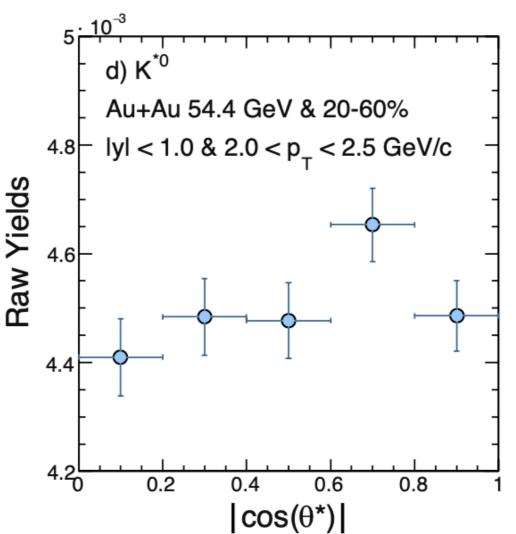
So the quantization axis has to be randomized in 3D-plane (x, y, z), to get the ρ_{00} to be ~ 1/3



Raw yield vs. $\cos \theta^*$



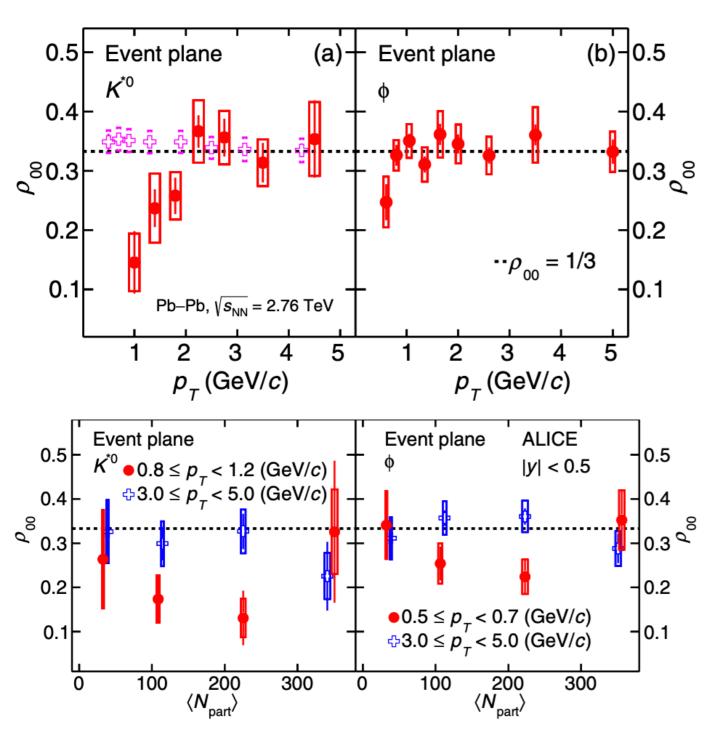




Raw yield is calculated in $\cos\,\theta^*$ bins

Spin alignment from Pb+Pb collisions at LHC

In Pb+Pb collisions, we measured ρ_{00} wrt event plane



ALICE: Phys. Rev Lett 125, 012301 (2020)

• For both K* and ϕ :

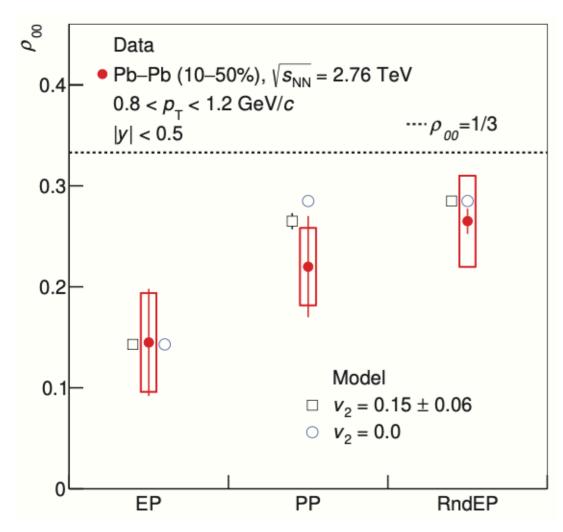
•
$$\rho_{00}^{EP} \le \frac{1}{3}$$

Significance of deviation for $K^* \leq 2.6\sigma$ for $\phi \leq 1.9\sigma$



Spin alignment between PP and EP

 ho_{00} wrt production (PP) and event plane (EP)



Event Production 2D-random plane plane plane

ALICE: Phys. Rev Lett 125, 012301 (2020)

• $ho_{00}^{\it EP}$ and $ho_{00}^{\it PP}$ are related by:

$$\bullet \left(\rho_{00}^{PP} - \frac{1}{3}\right) = \left(\rho_{00}^{EP} - \frac{1}{3}\right) \frac{1 + 3v_2}{4}$$



First measurement of charged K* ρ_{00}

Naive expectation

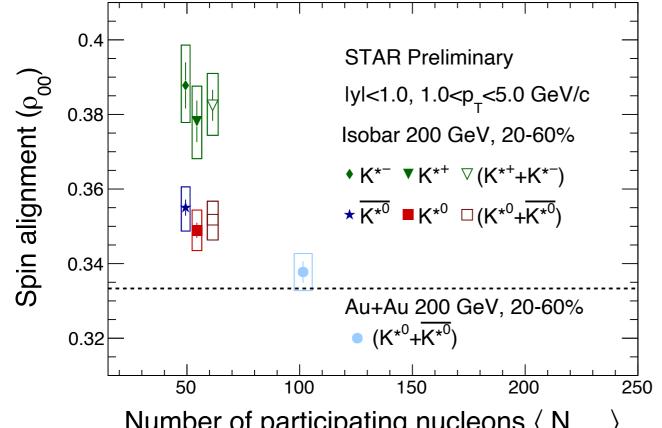
Particle Species	Magnetic moment (μ _N)
$K^{*0}(d\bar{s})$	$\mu_d \approx -0.97, \mu_{\bar{s}} \approx 0.61 \mu_N$
$K^{*+}(u\bar{s})$	$\mu_d \approx 1.85, \mu_{\bar{s}} \approx 0.61 \mu_N$

$$\rho_{00}(K^{*0}) > \rho_{00}(K^{*\pm})$$

$$\rho_{00}(B) \approx \frac{1}{3} - \frac{4}{9} \beta^2 \mu_{q_1} \mu_{q_2} B^2$$

Yang, et. al., Phys. Rev. C 97, 034917 (2018)

Observation from STAR



Number of participating nucleons $\langle N_{part} \rangle$

STAR: Quark Matter 2022

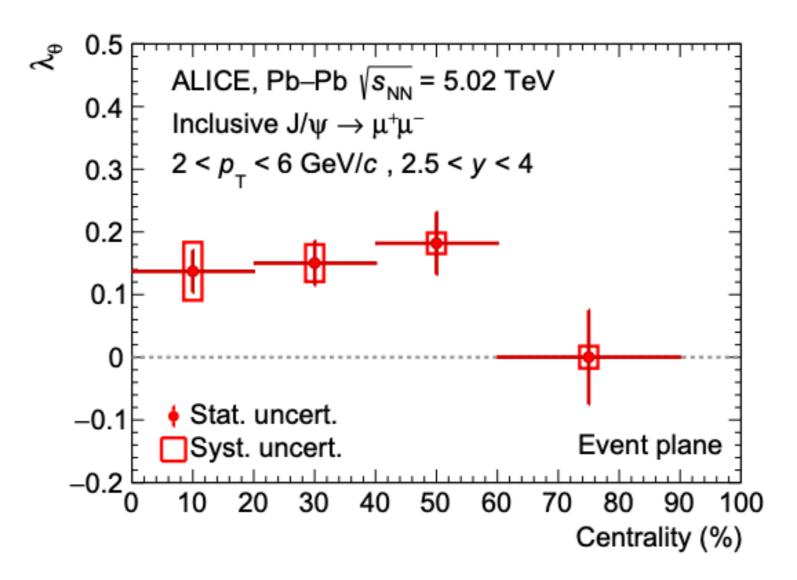
- First observation of $K^{*\pm} \rho_{00}$ in HIC
- Surprising ordering in ρ_{00}

$$K^{*+} + K^{*-} >> K^{*0} + \overline{K}^{*0}$$

(Opposite to naive expectation from B-field)



J/Ψ spin alignment in Pb+Pb collisions at LHC



First observation of global spin alignment of J/Ψ at LHC

$$J/\psi: \lambda_{\theta} \sim 0.2, \, \rho_{00} \sim 0.37 \, (>\frac{1}{3})$$

$$\lambda_{\theta} \propto \frac{3\rho_{00} - 1}{1 - \rho_{00}}$$

•
$$K^{*0} \rho_{00} \sim -0.2$$
, $\phi \rho_{00} \sim -0.1$