Studying the QCD Phase Diagram in RHIC-BES at STAR

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✓ Introduction

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- Beam Energy Scan
- The STAR detector

Experimental results

- Freeze-out conditions
- The 1st-order phase transition
- **Critical point** •
 - Higher-order fluctuations
 - Light nuclei production
- Crossover •
- ✓ BES-II

✓ Summary







\checkmark Need to investigate the QCD phase structure in wide (μ_B ,T) region.



A. Bzdak et al, 1906.00936









✓ Large data sets in various collision energies. ✓ Large and homogeneous acceptance, especially important for fluctuation analysis.

√S _{NN} (GeV)	Events (10 ⁶)	Year
200	350	2010
62.4	67	2010
54.4	1200	2017
39	39	2010
27	70	2011
19.6	36	2011
14.5	20	2014
11.5	12	2010
7.7	4	2010







The STAR detector

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Large & uniform acceptance

Excellent particle identification







 \checkmark The combined PID with m² from TOF is used at high p_T region.



 \checkmark dE/dx measured with TPC is used for proton identification at low p_T region.

















✓ Chemical freeze-out

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- Weak temperature dependence
- Centrality dependence of μ_B



✓ Kinetic freeze-out

- Central collisions \rightarrow lower value of T_{kin} and larger collectivity $<\beta>$
- Stronger collectivity at higher energy, even for peripheral collisions.







Freeze-out conditions



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✓ Both the larger separation of the freeze-out temperature (T_{ch}-T_{kin}) and stronger collectivity imply a longer hadronic interactions at higher collision energies.

PRC96, 44904(2017) : STAR Collaboration





V1 Versus collision energy



PRL120, 62301(2018) : STAR Collaboration

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v₁ slope versus collision energy

300



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STAR: 1708.07132; PRL120, 62301(2018)

- ✓ Minimum at $\sqrt{s_{NN}} = 14.5$ GeV for net-proton and net- Λ , but net-kaon v₁ slope continue decreasing as energy decreases.
- \checkmark At low energy, or in the region where the net-baryon density is large, repulsive force is expected, v_1 slope is large and positive.
- Softest point only for baryons? \checkmark
- Need model to explain \checkmark

M. Isse, A. Ohnishi et al, PRC72, 064908(05) Y. Nara, A. Ohnishi, H. Stoecker, PRC94, 034906(16)







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Higher-order fluctuations

Moments and cumulants are mathematical measures of "shape" of a distribution which probe the fluctuation of observables.

- \checkmark
- S and k are non-gaussian fluctuations. \checkmark



Cumulant *⇒* **Moment** \checkmark

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 $<\delta N>=N-<N>$ $C_1 = M = \langle N \rangle$ $C_2 = \sigma^2 = \langle (\delta N)^2 \rangle$ $C_3 = S\sigma^3 = \langle \delta N \rangle^3 >$ $C_4 = \kappa \sigma^4 = \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2$

Moments: mean (M), standard deviation (σ), skewness (S) and kurtosis (κ).



Cumulant : additivity \checkmark

 $C_n(X+Y) = C_n(X) + C_n(Y)$

proportional to volume









$$C_2 = \langle \delta N \rangle^2 >_c \approx \xi^2 \qquad C_5 = \langle \delta N \rangle^5 >_c \approx \xi^6$$

$$C_3 = \langle \delta N \rangle^3 >_c \approx \xi^{4.5} \quad C_6 = \langle \delta N \rangle^6 >_c \approx \xi^1$$

$$C_4 = \langle \delta N \rangle^4 >_c \approx \xi^7$$

$$S\sigma = \frac{C_3}{C_2} = \frac{\chi_3}{\chi_2} \quad \kappa \sigma^2 = \frac{C_4}{C_2} = \frac{\chi_4}{\chi_2}$$
$$\chi_n^q = \frac{1}{VT^3} \times C_n^q = \frac{\partial^n p / T^4}{\partial \mu_q^n}, \quad q = B, Q, S$$





Data analysis method



- X.Luo, J. Xu, B. Mohanty and N. Xu. J. Phys. G40,105104(2013)
- M. Kitazawa : PRC.86.024904, M. Kitazawa and M. Asakawa : PRC.86.024904
- Bzdak and V. Koch : PRC.86.044904, PRC.91.027901, X. Luo : PRC.91.034907
- T. Nonaka, M. Kitazawa, S. Esumi : PRC.95.064912
- X. Luo, T. Nonaka : PRC.99.044917

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\checkmark At $\mu_B < 210$ MeV, the 4th-order fluctuation is found to be flat as a function of collision energy.



PRL105, 022302(2010): STAR Collaboration







Net proton from BES-I



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✓ Deviation below Poisson baseline (unity).

✓ Both 3rd- and 4th-order fluctuations have their minima at $\sqrt{s_{NN}} = 19.6$ GeV.

PRL112, 032302(2014): STAR Collaboration









 $error(\kappa\sigma^2) \propto \frac{\sigma^2}{\varepsilon^2} \frac{1}{\sqrt{N_{evts}}}$

Large statistical uncertainties, need more data.







\checkmark p_T region can be extended up to 2.0 GeV by using TOF as well as TPC. \checkmark (Anti)proton statistics is doubled with respect to the published results.









Cumulants in BES-I energies



STAR, PoS CPOD2014 (2015)019

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STAR, PoS CPOD2014 (2015)019

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- $\sqrt{\kappa\sigma^2 (C_4/C_2)}$ shows a non-monotonic behaviour. The trend is consistent with the theoretical calculation.
- ✓ Measurement at lower energies are important.









New data set : 54.4 GeV

Collision system and energy	Au+Au and 54.4 GeV		
Baryon Chemical Potential	~ 90 MeV		
No. of events	~ 553 M		
Collision centrality	0-5%, 5-10%, 10-20%, 20-30%, 30-40%, 40 50%, 50-60%, 60-70%, 70-80%		
Centrality	η < 1; charged particles other than protons and antiprotons		
Z Vertex	+/- 30 cm		
Vertex radial position	2 cm		
Detectors	Time Projection Chamber and Time-of-Flig		
Particle Type	Proton and antiprotons		
Rapidity	+/- 0.5		
Transverse Momentum Range	0.4 to 2.0 GeV/c		
Secondary proton backgrounds	DCA < 1cm		

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Colliding energy √s_{NN} (GeV)

✓ Results at 54.4 GeV follow the trend very well.







 \checkmark At $\sqrt{s_{NN}} < 10$ GeV, data shows $\kappa \sigma^2 > 1$, while model shows $\kappa \sigma^2 < 1$ No model simulations can explain the enhancement at low beam energies.



Z. Feckova, et al., PRC92, 064908(2015). J. Xu, et. al., PRC94, 024901(2016). X. Luo et al., NPA931, 808(14), P.K. Netrakanti et al. 1405.4617, NPA947, 248(2016), P. Garg et al. PLB 726, 691(2013). S. He, et. al., PLB762, 296 (2016). S. He, X. Luo, PLB 774, 623 (2017).







 \checkmark There isn't yet any experimental evidence for the smooth crossover at $\mu_{\rm B} \sim 0$ MeV.

Sixth-order cumulants of net-charge and net-baryon distributions are predicted to be negative if the chemical freeze-out is close enough to the phase transition, which is the characteristic signal for $\sqrt{s_{NN}}$ >60 GeV.



Positive sign is predicted in $\sqrt{s_{NN}}$ <60 GeV

C.Schmidt,Prog.Theor.Phys.Suppl.186,563–566(2010) Cheng et al, Phys. Rev. D 79, 074505 (2009) Friman et al, Eur. Phys. J. C (2011) 71:1694

 χ_6^Q/χ_2^Q $\chi_4^{\rm B}/\chi_2^{\rm B}$ $\chi_6^{\rm B}/\chi_2^{\rm B}$ $\chi_4^{\rm Q}/\chi_2^{\rm Q}$ Freeze-out conditions HRG ~ 2 ~ 10 QCD: $T^{\rm freeze}/T_{pc} \lesssim 0.9$ $\gtrsim 1$ $\gtrsim 1$ ~ 10 ~ 2 QCD: $T^{\text{freeze}}/T_{pc} \simeq 1 \qquad \sim 0.5$ ~ 1 <0 <0 Predicted scenario for this measurement 1.2





Net-proton at $\sqrt{s_{NN}} = 200 \text{ GeV}$

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✓ Much better precision as compared to net-charge. Negative values are observed systematically from midcentral to central collisions, which seems consistent with the theoretical prediction.



$$error(C_r) \propto \frac{\sigma^r}{\sqrt{N_{\rm eve}}}$$

Used statistics

	0-10%	10-80%
Run10	~160M	~200M
Run11	~50M	~450M





Lower beam energy?

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✓ Positive C₆ is predicted in $\sqrt{s_{NN}}$ <60 GeV (µ_B/T<0.5). ✓ STAR collected ~550M events at $\sqrt{s_{NN}} = 54.4$ GeV in 2017.

Friman et al, Eur. Phys. J. C (2011) 71:1694

If freeze-out occurs close to the chiral crossover temperature the sixth order cumulant of the net baryon number fluctuations will be negative at LHC energies as well as for RHIC beam energies $\sqrt{s_{NN}} \gtrsim 60$ GeV, corresponding to $\mu_B/T \leq 0.5$. This is in contrast to hadron resonance gas model calculations which yield a positive sixth order cumulant.



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Net-proton C_6/C_2 at $\sqrt{s_{NN}} = 54.4$ GeV

 \checkmark Positive values are observed systematically from peripheral to central collisions.









\checkmark Clear separation and opposite signs between two energies in 0-40%.













✓ Clear separation and opposite signs between two energies in 0-40%. ✓ UrQMD result shows positive signs for all centralities at $\sqrt{s_{NN}}$ = 200 GeV.





Acceptance dependence

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Monotonic decrease with enlarging the acceptance.

✓ p_T dependence seems to be saturated at 0.4< p_T <1.7 GeV/c.







\checkmark Results from the central bin of 200 GeV Au+Au collisions are consistent with the LQCD prediction \rightarrow remittance of chiral phase transition?









or 1st-order phase transition



K. J. Sun, L. W. Chen, C. M. Ko, Z. Xu, Phys. Lett. B774, 103 (2017). K. J. Sun, L.W. Chen, C. M. Ko, J. Pu, Z. Xu, Phys. Lett. B781, 499 (2018). Edward Shuryak and Juan M. Torres-Rincon, arXiv:1805.04444

Server A Baryon density fluctuation becomes large near the critical point

Coalescence + nucleon density fluctuation $N_{
m d} \;=\; rac{3}{2^{1/2}} \left(rac{2\pi}{m_0 T_{
m eff}}
ight)^{3/2} \; N_p \langle n
angle (1 + lpha \Delta n),$ $N_{^{3}\mathrm{H}} = \frac{3^{3/2}}{4} \left(\frac{2\pi}{m_{0}T_{\mathrm{eff}}}\right)^{3} N_{p} \langle n \rangle^{2} [1 + (1 + 2\alpha)\Delta n],$ $N_t \cdot N_p / N_d^2 \approx g(1 + \Delta n)$

> Neutron density fluctuation $\Delta n = \langle (\delta n)^2 \rangle / \langle n \rangle^2$





Light nuclei production



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✓ Deuteron and triton yields have been measured in **BES-I energies.**











Non-monotonic energy dependence is observed in Au+Au collisions with a peak around 20-30 GeV.







√s _{NN} (GeV)	Events (10 ⁶)	BES-II / BES-I	Weeks	μ _Β (MeV
200	350	2010		25
62.4	67	2010		73
54.4	1200	2017		
39	39	2010		112
27	70	2011		156
19.6	400 / 36	2019-21 / 2011	3	206
14.5	300 / 20	2019-21 / 2014	2.5	264
11.5	230 / 12	2019-21 / 2010	5	315
9.2	160 / 0.3	2019-21 / 2008	9.5	355
7.7	100 / 4	2019-21 / 2010	14	420

✓ BES-II has started this year.

T _{CH} (MeV)
166
165
164
162
160
156
152
140
140

- Luminosity has been improved with the electron cooling system.
- ✓Inner TPC has been fully integrated, which extends the pseudorapidity coverage from 1.0 to 1.5
- ✓ New centrality definition by EPD.
- ✓ eTOF for fixed-target program
- ✓ Higher-order fluctuation measurement with small errors and large acceptance.









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✓ Higher-order fluctuation measurement with small errors and large acceptance.

✓ Statistical uncertainties will be dramatically reduced.

✓ Can we measure a possible "peak" structure?

- freeze-out in various collision energies.
- measuring v_1 slope in various collision energies.
- to the theoretical prediction.
- Non-monotonic energy dependence is observed for neutron density fluctuation in central Au+Au collisions.
- BES-II will shrink the statistical uncertainties at 7.7–19.6 GeV region.

• Freeze-out conditions have been determined in terms of chemical/kinetic

Search for the softest point as a signal from the 1st-order phase transition, by

• Non-monotonic behavior has been observed for net-proton $\kappa\sigma^2$ in the central Au+Au collisions as a function of collision energy, which is qualitatively similar

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Thank you for your attention

Back up

\checkmark Single-particle tracking efficiencies for $\pi/K/p$ have been estimated by embedding simulation.

✓ TOF matching efficiency is obtained from the real data.

Efficiency

- Fror estimation using a simple toy model with **Skellam distributions.**
- One order of cumulant increases, statistical errors becomes x10. ($C_4 \rightarrow C_6$: x100 err)
- ✓ C₆ measurement is more challenging than C₄.

STAR Non-binomial efficiencies

Experimental effects

✓ The detector efficiency may not be binomial, which would be due to the particle misidentification, track splitting/merging effects, and many other reasons.

2. Multiplicity dependent efficiency

✓ Residual dependence of efficiency inside one multiplicity bin (for centrality) needs to be taken into account.

A. Bzdak, R. Holzmann, V. Koch : PRC.94.064907

→ One example of nonbinomial distribution, Beta-binomial, is wider distribution than binomial

Unfolding

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- ✓ We performed MC simulations by embedding protons and antiprotons, e.g., N_p=60 and N_{pbar}=15 (which would be an extreme number), and see whether those particles can be reconstructed or not.
- The response matrix is close to the beta-binomial distribution, which is wider than binomial.
 - "Urn model" for beta-binomial distribution, where the parameter α controls the deviation from binomial.

N_w : white balls, N_b: black balls, ϵ : efficiency $N_w = \alpha N_p \quad \epsilon = N_w/(N_w + N_b)$

