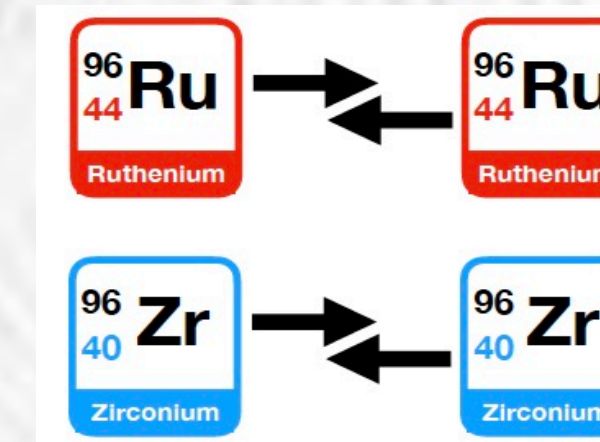


Search for the Chiral Magnetic Effect: Recent results from isobar collisions at RHIC



Sergei A. Voloshin



for the  STAR ☆ Collaboration

Search for the Chiral Magnetic Effect with Isobar Collisions at $\sqrt{s_{NN}} = 200$ GeV by the STAR Collaboration at RHIC

STAR Collaboration arXiv:2109.00131v1 [nucl-ex] 1 Sep 2021

+ one slide on results from

Search for the chiral magnetic effect via charge-dependent azimuthal correlations relative to spectator and participant planes in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV

STAR Collaboration arXiv:2106.09243v1 [nucl-ex] 17 Jun 2021

Outline:

CME: observables and background

First measurements

STAR isobar blind analysis results

Implications

Supported by



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ENERGY

Office of
Science

D. Kharzeev, Parity violation in hot QCD: Why it can happen and how to look for it, *Phys. Lett. B* **633**, 260 (2006).

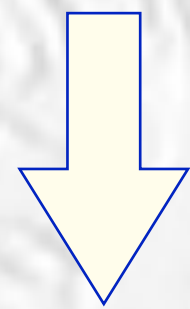
S. A. Voloshin, Parity violation in hot QCD: How to detect it, *Phys. Rev. C* **70**, 057901 (2004).

Chirality imbalance

$$\mu_5 \propto \mathbf{E} \cdot \mathbf{B}$$

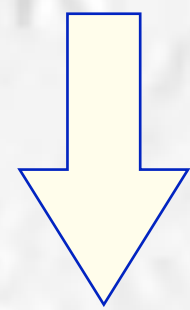
+

Magnetic field

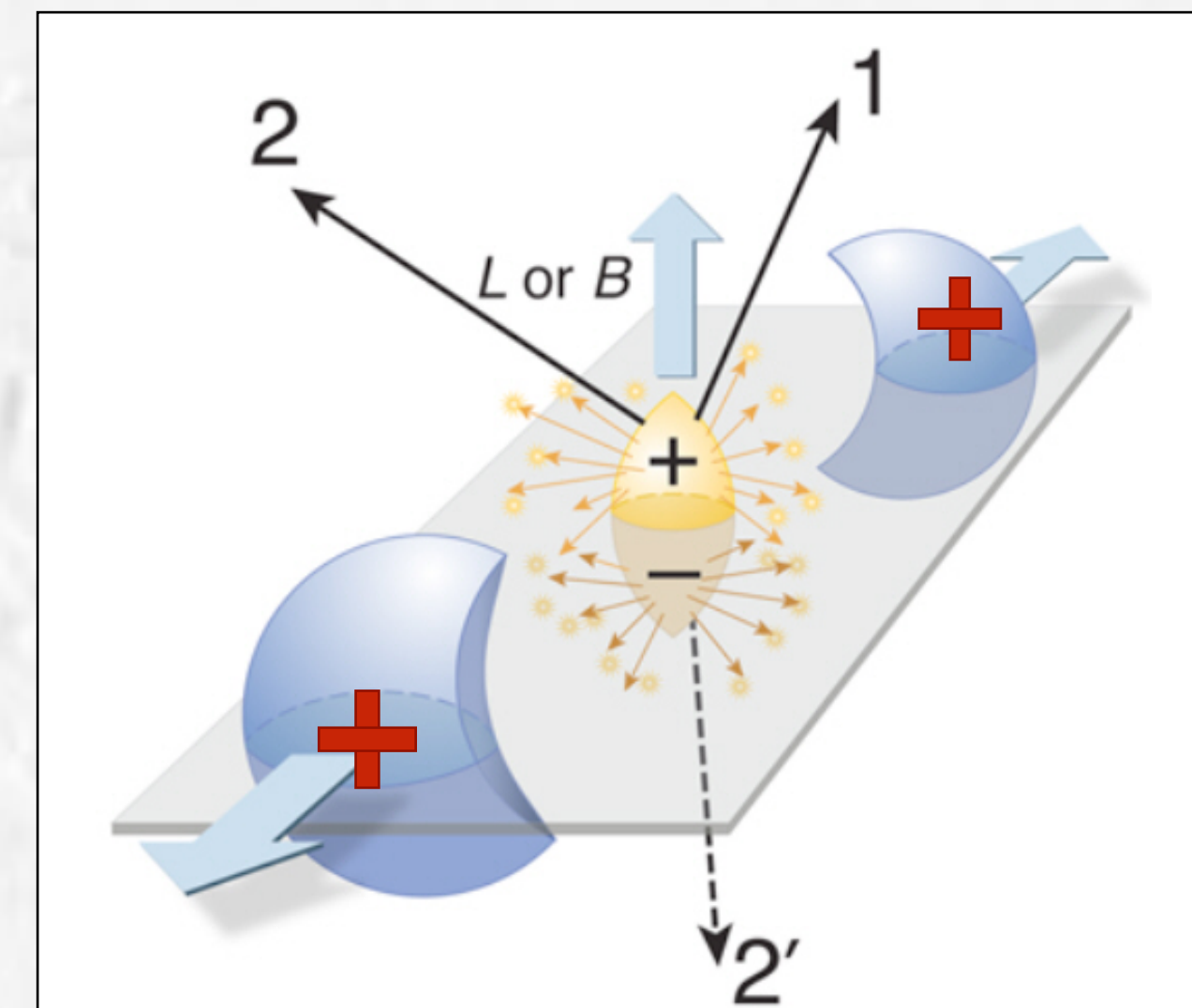
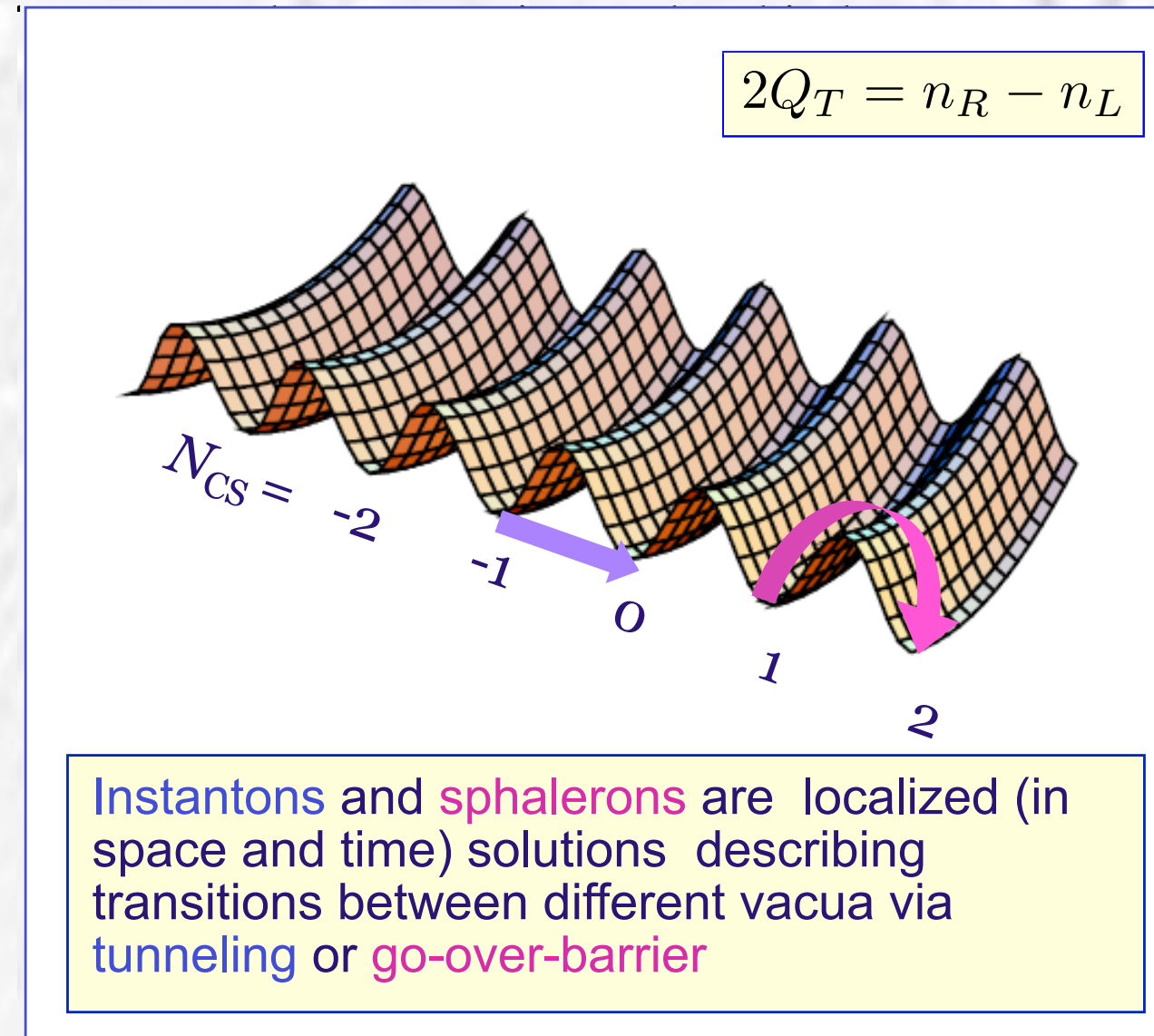


Anomalous transport,
Electric current along
magnetic field

$$\mathbf{J} = \frac{e^2}{2\pi^2} \mu_5 \mathbf{B}$$



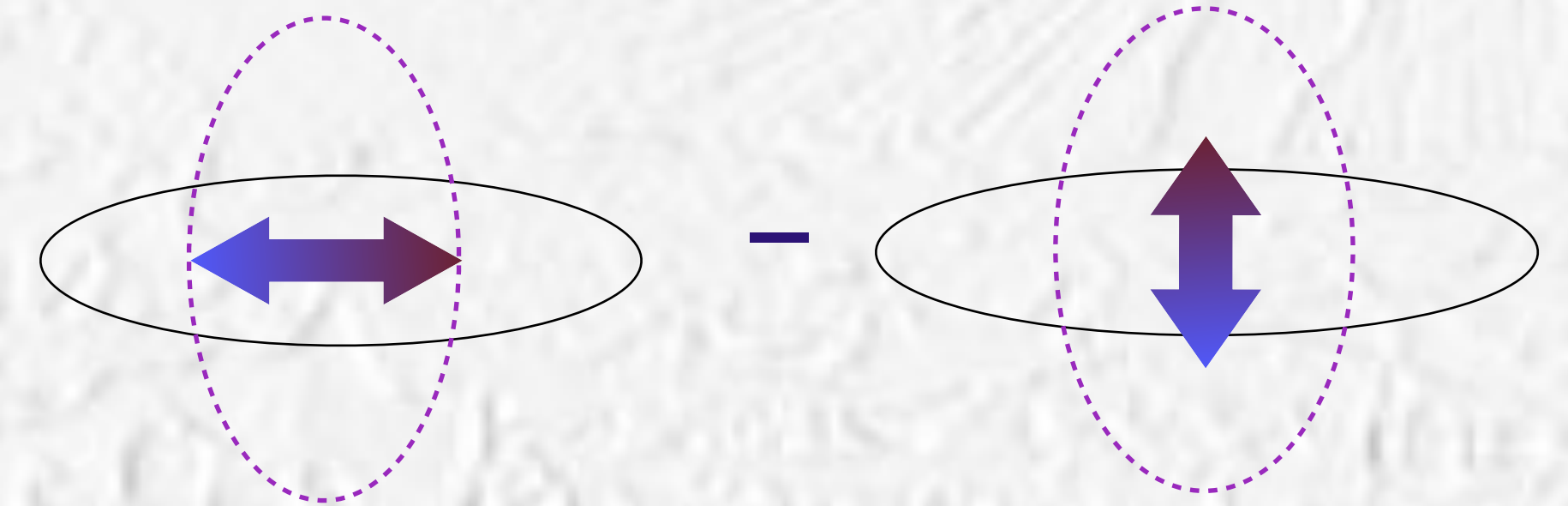
Charge separation
along \mathbf{B} direction



Effective particle distribution

$$\frac{dN_{\pm}}{d\phi} \propto 1 + 2v_1 \cos(\Delta\phi) + 2v_2 \cos(2\Delta\phi) + \dots + 2a_{1,\pm} \sin(\Delta\phi) + \dots ; \quad \Delta\phi = \phi - \Psi_{RP}$$

$$\begin{aligned} \gamma_{\alpha,\beta} &\equiv \langle \cos(\phi_{\alpha} + \phi_{\beta} - 2\Psi_{RP}) \rangle \\ &= \langle \cos \Delta\phi_{\alpha} \cos \Delta\phi_{\beta} \rangle - \langle \sin \Delta\phi_{\alpha} \sin \Delta\phi_{\beta} \rangle \\ &= [\langle v_{1,\alpha} v_{1,\beta} \rangle + B_{in}] - [\langle a_{1,\alpha} a_{1,\beta} \rangle + B_{out}] \end{aligned}$$



The sign of the correlations is sensitive to the "direction" (in- or out-of-plane), the background is suppressed ($B_{in}-B_{out}$) at least by a factor of $v_2 < 10^{-1}$.

$$B_{in} - B_{out} \propto v_{2,clust} \langle \cos(\phi_{\alpha} + \phi_{\beta} - 2\phi_{clust}) \rangle$$

Azimuthal Charged-Particle Correlations and Possible Local Strong Parity Violation

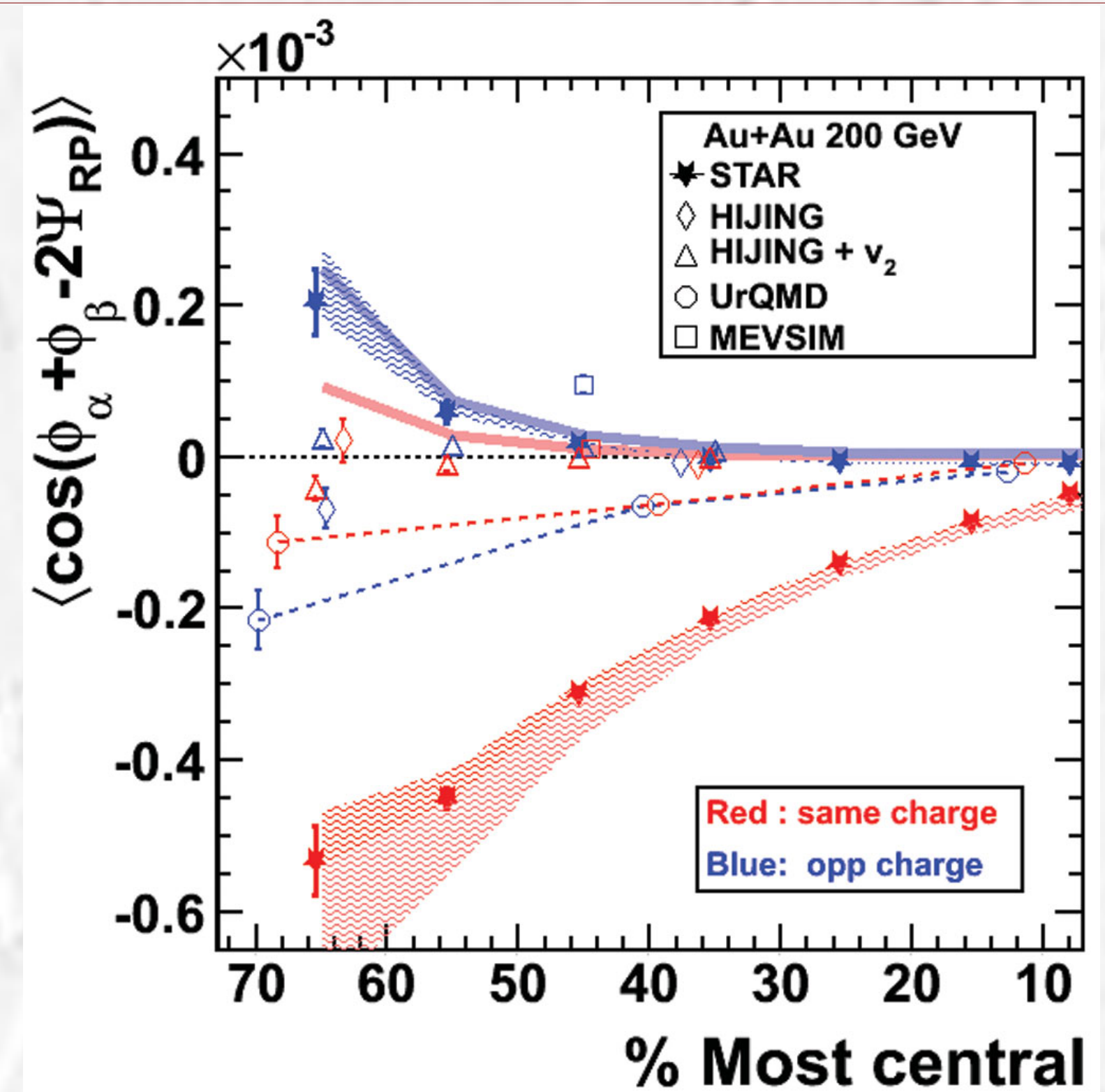


FIG. 4 (color). $\langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle$ results from 200 GeV Au + Au collisions are compared to calculations with event generators HIJING (with and without an “elliptic flow afterburner”), URQMD (connected by dashed lines), and MEVSIM. Thick lines represent HIJING reaction-plane-independent background

HIJING+v2 = added “afterburner” to generate flow
MEVSIM: flow as in experiment, maximum number of resonances that is consistent with experiment

No event generator could explain more than 1/3 of the signal
 $\Delta\gamma = \gamma_{\text{opp}} - \gamma_{\text{same}}$

Observation of charge-dependent azimuthal correlations and possible local strong parity violation in heavy-ion collisions

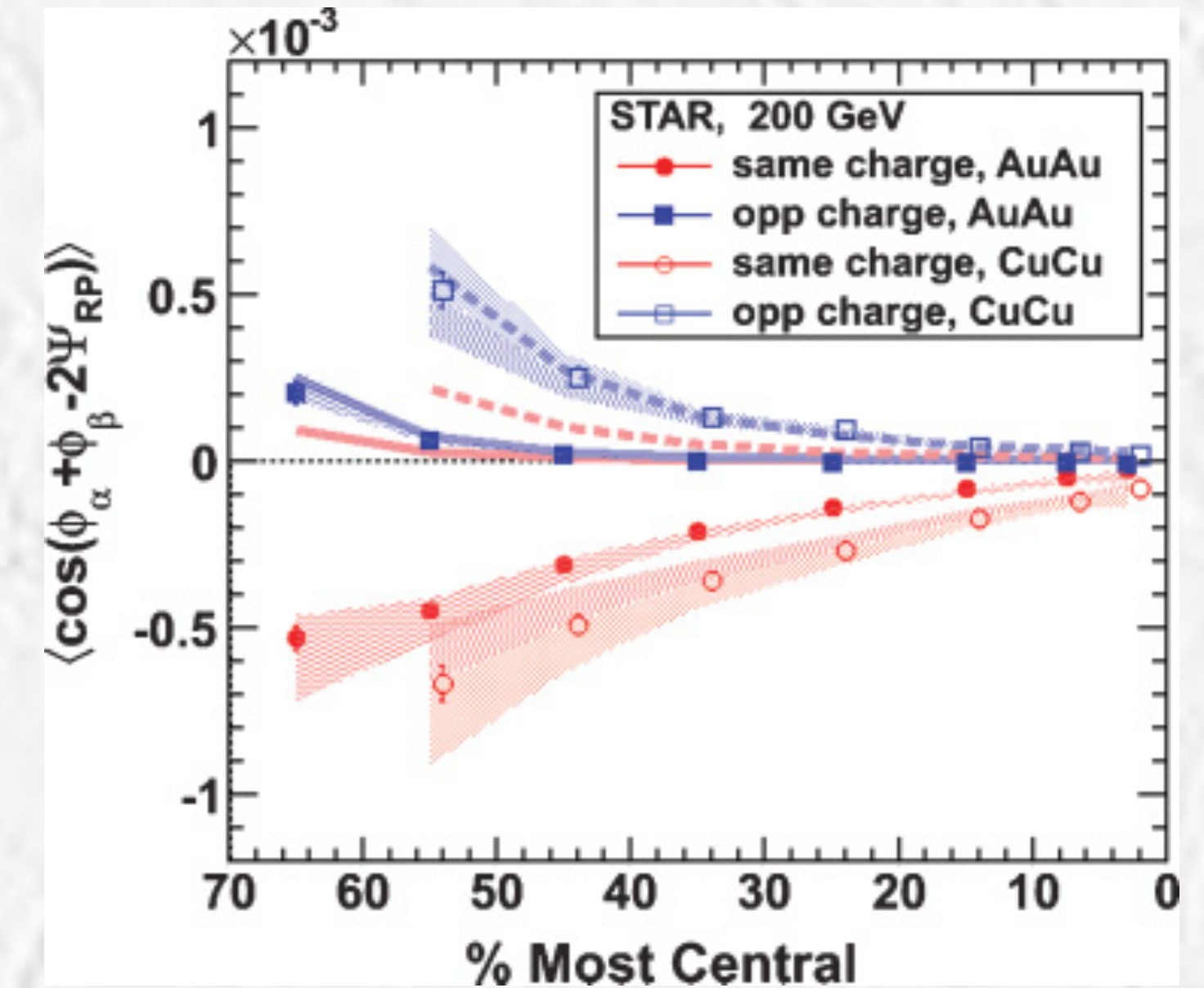


FIG. 7. (Color online) $\langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle$ in Au + Au and Cu + Cu collisions at $\sqrt{s_{NN}} = 200$ GeV calculated using Eq. (7). The error bars show the statistical errors. The shaded areas reflect the uncertainty in the elliptic flow values used in calculations, with lower (in magnitude) limit obtained with elliptic flow from two-particle correlations and upper limit from four-particle cumulants. For details, see Sec. IV. Thick solid (Au + Au) and dashed (Cu + Cu) lines represent possible non-reaction-plane-dependent contribution from many-particle clusters as estimated by HIJING (see Sec. VII A).

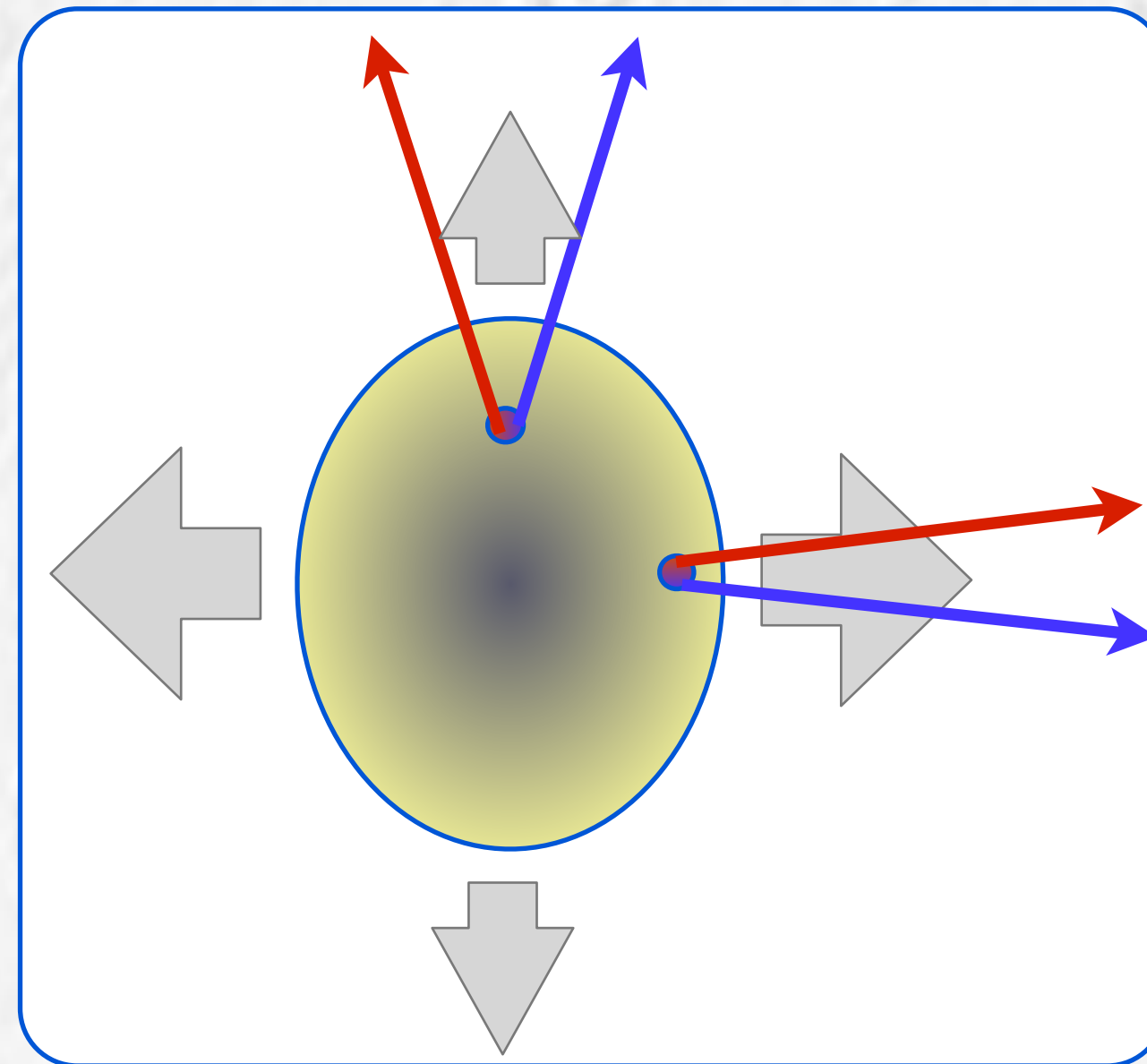
STAR ☆ Types of the background

I. Physics (RP dependent). (Can not be suppressed)

$$\begin{aligned} \gamma_{\alpha,\beta} &\equiv \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle \\ &= \langle \cos \Delta\phi_\alpha \cos \Delta\phi_\beta \rangle - \langle \sin \Delta\phi_\alpha \sin \Delta\phi_\beta \rangle \\ &= [\langle v_{1,\alpha} v_{1,\beta} \rangle + B_{in}] - [\langle a_{1,\alpha} a_{1,\beta} \rangle + B_{out}] \end{aligned}$$

“Flowing clusters” (including LCC)
charge dependent directed flow.

$$B_{in} - B_{out} \propto v_{2,clust} \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{clust}) \rangle$$



Local charge conservation (LCC)

Pratt, arXiv:1002.1758v1[nucl-th]
Schlichting and Pratt, PRC83 014913 (2011)

Note, LCC:

- Correlations only between opposite charges
- To be consistent with data must be combined with (negative) charge independent correlations (e.g. momentum conservation).
- No event generator exhibits such strong correlations as predicted by the Blast Wave model

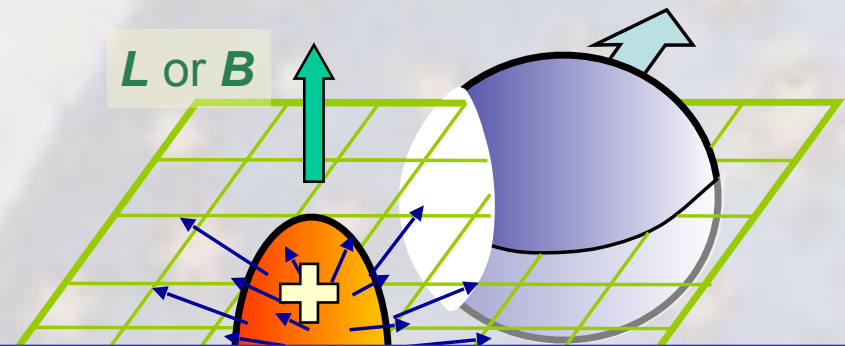
II. Measurements (RP independent).

(depends on method, in principle can be reduced)

$$\langle \cos(\phi_a + \phi_b - 2\phi_c) \rangle \stackrel{?}{\rightarrow} \langle \cos(\phi_a + \phi_b - 2\Psi_2) \rangle v_{2,c}$$

Observation of charge-dependent azimuthal correlations and possible local strong parity violation in heavy-ion collisions

Sergei A. Voloshin
 for the Collaboration



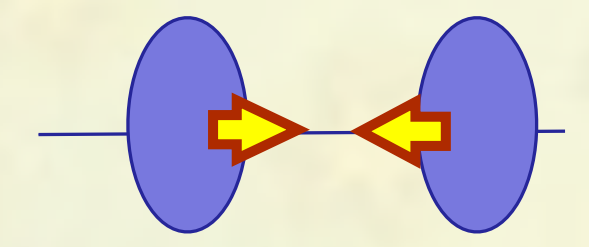
Outline:

- Chiral Magnetic Effect
- STAR results (PRL, arXiv: 0909.1717, ...)

Developing a program

Dedicated experimental and theoretical program focused on the local parity violation, and more generally on non-perturbative QCD: structure of the vacuum, hadronization, etc.

Experiment:

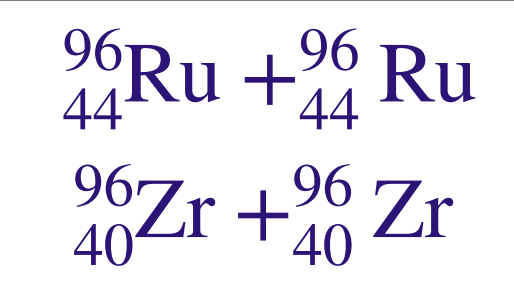


- U+U central body-body collisions
- Beam energy scan / Critical point search
- Isobaric beams
- High statistics PID studies / properties of the clusters

Such collisions ("easy" to trigger on) will have low magnetic field and large elliptic flow – clean test of the *LPV* effect.

Look for a critical behavior, as *LPV* predicted to depend strongly on **deconfinement and chiral symmetry restoration**

Colliding isobaric nuclei (the same mass number and different charge) and by that controlling the magnetic field



Note that such studies will be also very valuable for understanding the initial conditions, baryon stopping, origin of the directed flow, etc.

in particular with neutral particles; see also next slide

It is likely that the measurements are dominated by the “background” (LCC?).

Goal: identification of the presence or the lack of the CME signal at the level of ~5% of the measured gamma correlator value

$$f_{\text{CME}} = \Delta\gamma^{\text{CME}} / \Delta\gamma$$

At hand: Signal depends on the magnetic field, the background depends on anisotropic flow

Collision energy dependence, Beam energy scan (signal should disappear at lower energies)

✓ Isobar collisions (vary magnetic field while keeping the same background)

Event Shape Engineering (increase/decrease background)

✓ Gamma wrt different Event Planes (Participants, Spectators)

✓ Higher harmonic correlators (background, no signal)

Small system collisions (background, no signal)

Very central collisions (background, no signal)

U+U (body-body vs tip-tip)

Correlations with identified particles (e.g. for the next bullet)

Cross-correlation of different observables, CME X CMW X CVE

New ideas/observables (invariant mass, Signed BF, R-correlator)

Studies of the EM fields



Sergei A. Voloshin
Nuclear Physics A 827 (2009) 377c–382c

Suggestion of using isobar beams
 ${}^{96}_{44}\text{Ru} + {}^{96}_{44}\text{Ru}$ and ${}^{96}_{40}\text{Zr} + {}^{96}_{40}\text{Zr}$
to disentangle CME signal from BG

PRL 105, 172301 (2010)

PHYSICAL REVIEW LETTERS

week ending
22 OCTOBER 2010

Deng et al. PHYSICAL REVIEW C 94, 041901(R) (2016)

STAR 2018 Beam Use Request

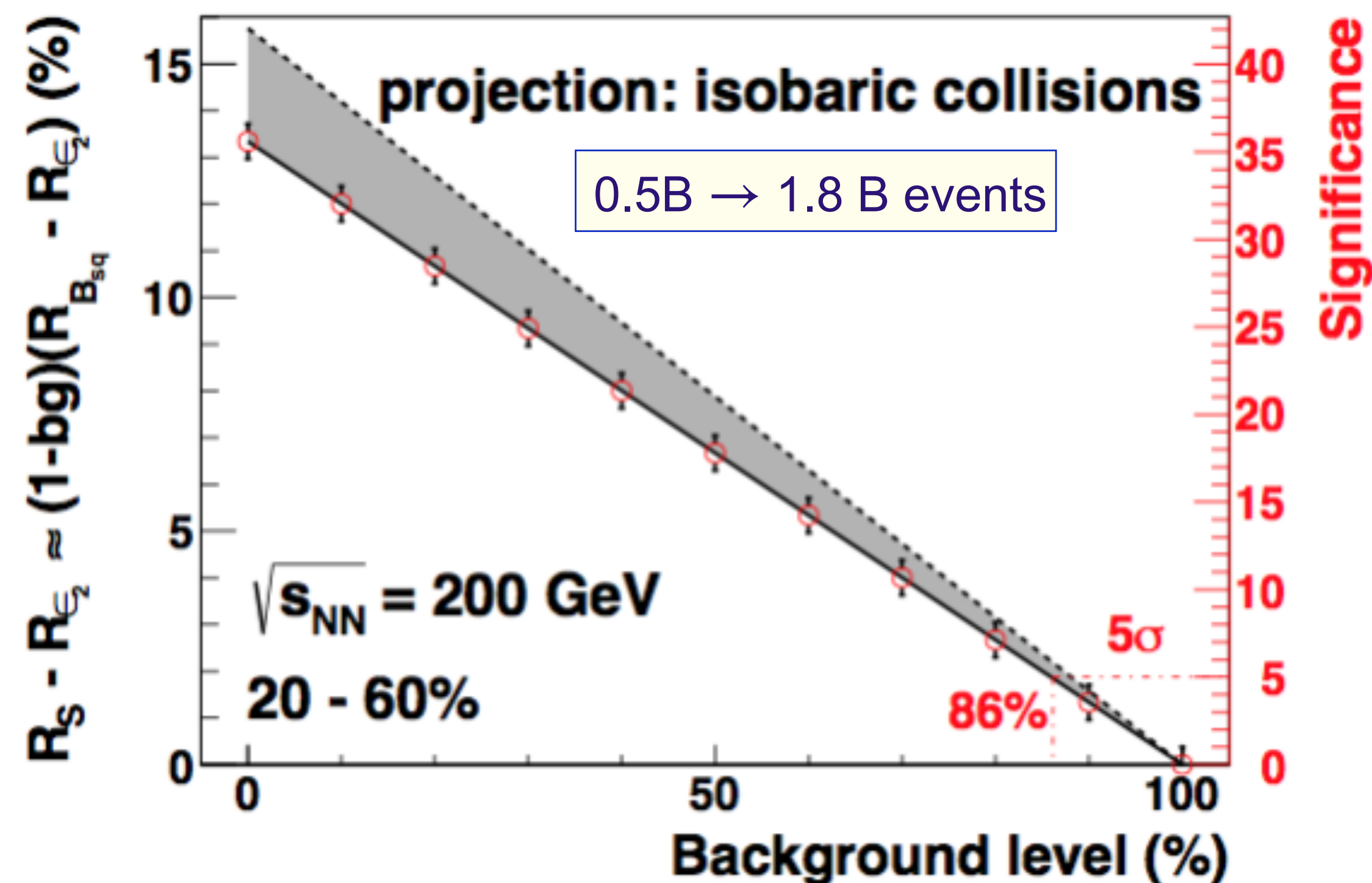
Testing the Chiral Magnetic Effect with Central $U + U$ Collisions

Sergei A. Voloshin

The charge separation dependence on the strength of the magnetic field can be further studied with the collision of isobaric nuclei, such as ${}^{96}_{44}\text{Ru}$ and ${}^{96}_{40}\text{Zr}$. These nuclei have

$$\frac{(\Delta\gamma/v_2)_{\text{Ru+Ru}}}{(\Delta\gamma/v_2)_{\text{Zr+Zr}}} \approx 1 + f_{\text{CME}}^{\text{Zr+Zr}} \underbrace{\left[\left(\frac{B_{\text{Ru+Ru}}}{B_{\text{Zr+Zr}}} \right)^2 - 1 \right]}_{\approx 0.18}$$

To measure f_{CME} at the level of 3% one has to measure the double ratio with accuracy 0.6%



$$f_{\text{CME}} = 1 - f_{\text{BG}}$$

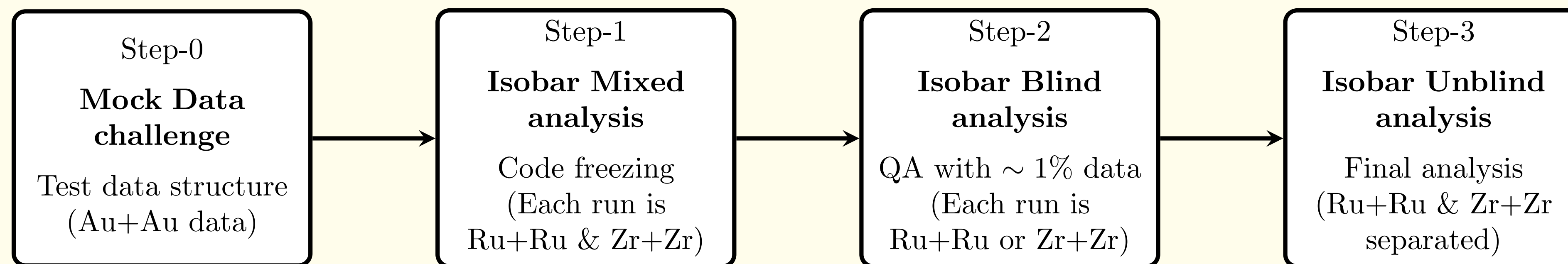
2017 BNL NPP Program Advisory Committee recommended *blind analyses* of *CME studies* of Run-18 isobar data

STAR develops

Methods for a blind analysis of isobar data collected by the STAR collaboration

NUCL SCI TECH (2021) 32:48

<https://doi.org/10.1007/s41365-021-00878-y>



Step-0 Mock Data challenge

Step-1 Provide output files composed of collision data from a mix of the two isobar species

Time-dependent QA tuned. **Codes, and Analysis Notes are frozen**

Step-2 Provide files that blind the isobar species but do not “mix” data from different data acquisition runs

Only allow “run-by-run” corrections

Step-3 Full unblinding. Codes run by analyzers from a different group.

Prior to the blind analysis, the CME signatures are predefined as a significant excess of the CME-sensitive observables in Ru+Ru collisions over those in Zr+Zr collisions, owing to a larger magnetic field in the former.

- Five institutional groups within the collaboration perform blind analyses of the isobar data
- Each group focuses on a specific analysis method
- Substantial overlap of some analyses helps to cross-check results
- All analyses have a common set of variations for the purpose of systematic uncertainty determination
- Set of common and analysis-specific variables for data QA and selection of the data with stable detector performance

The paper includes results ONLY for predefined observables described in the Analyses Notes and a very limited (~1/2 page) post-blinding section

- ◆ Large data set needed to hit *small statistical uncertainty* target
- ◆ *Systematic uncertainties* between species need to be controlled *below that level*

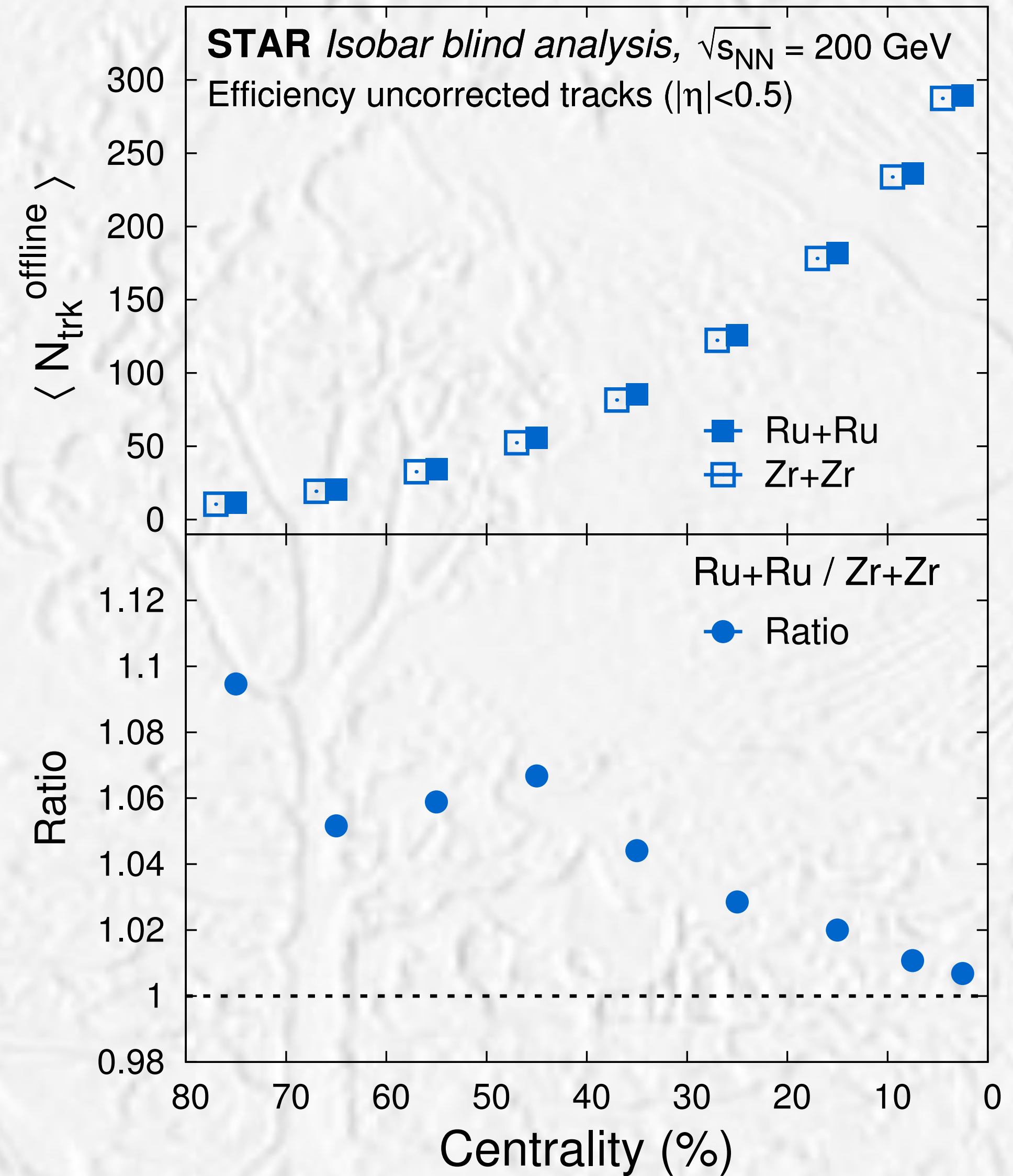
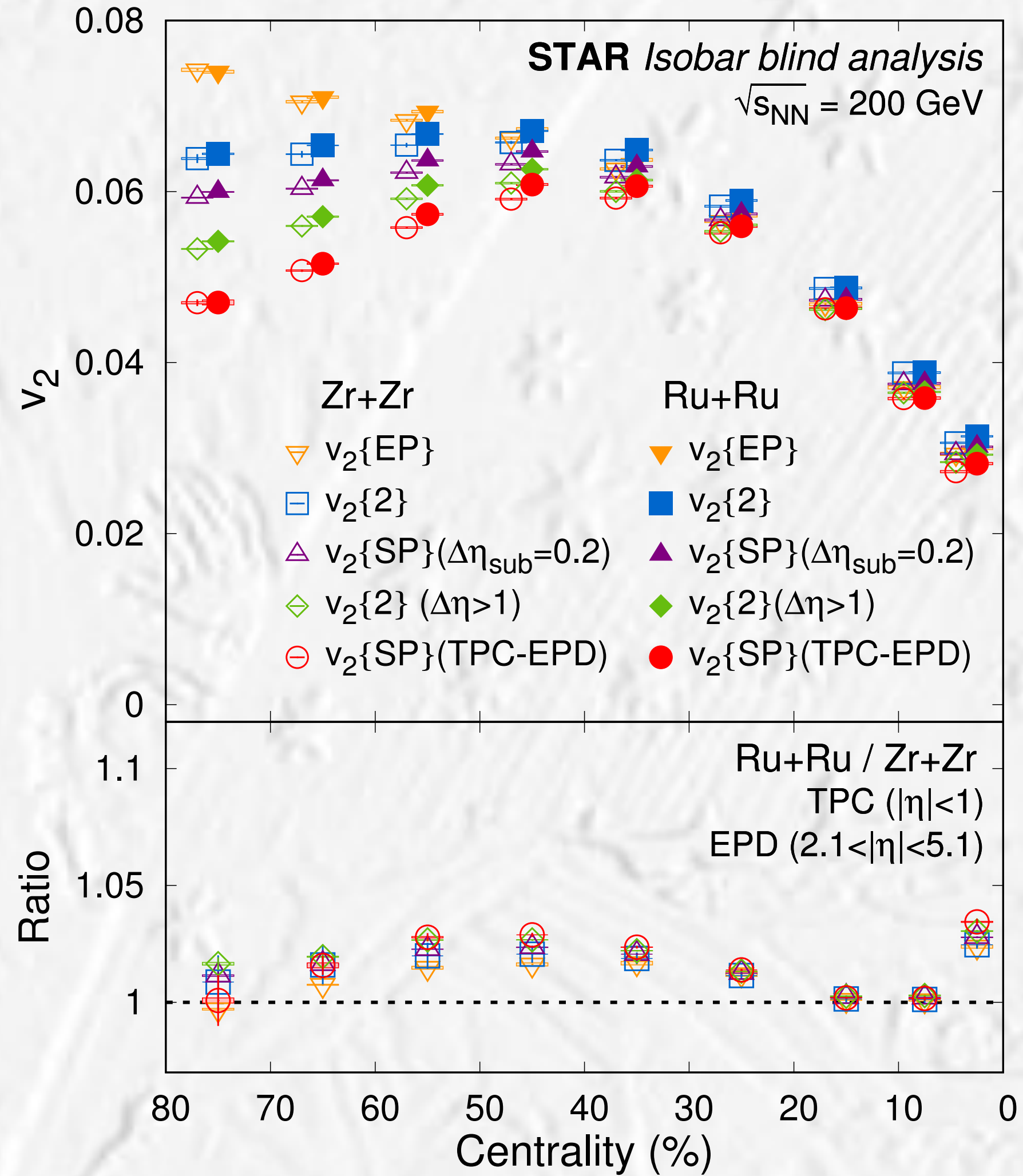
Special RHIC conditions (*G. Marr et al., 10th International Particle Accelerator Conference (2019) pp. 28–32*)

- ◆ Alternate the isobar species between each store of beam in RHIC
- ◆ Keep long stores with constant beam luminosity
- ◆ Match luminosities between the species
- ◆ Adjust the luminosity in such a way that the hadronic interaction rate at STAR is close to 10 kHz.

Precision target achieved:

A precision down to 0.4% is achieved, as anticipated, in the relative magnitudes of the pertinent observables between the two isobar systems

v_2 and multiplicity

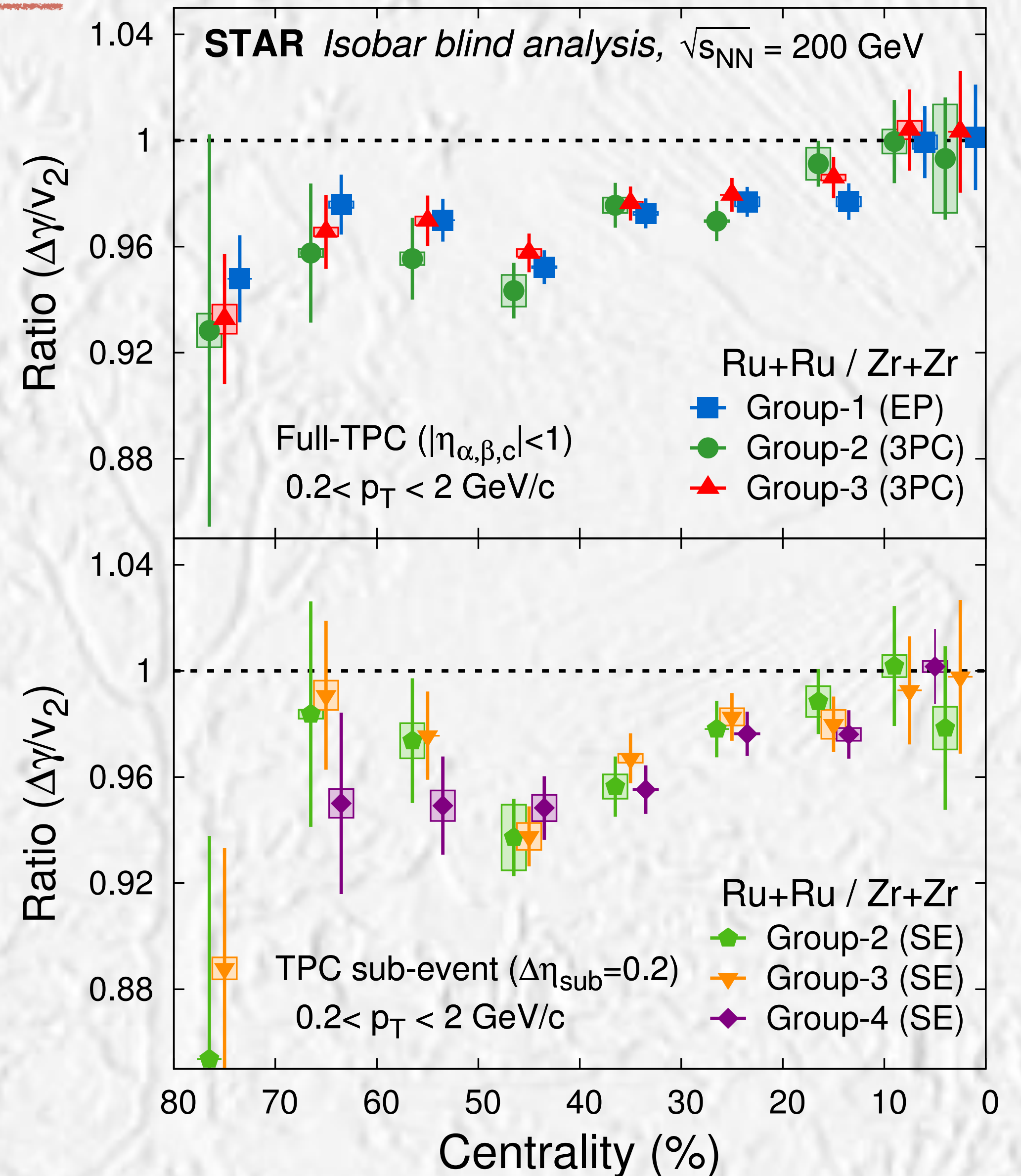


Groups 1 - 4

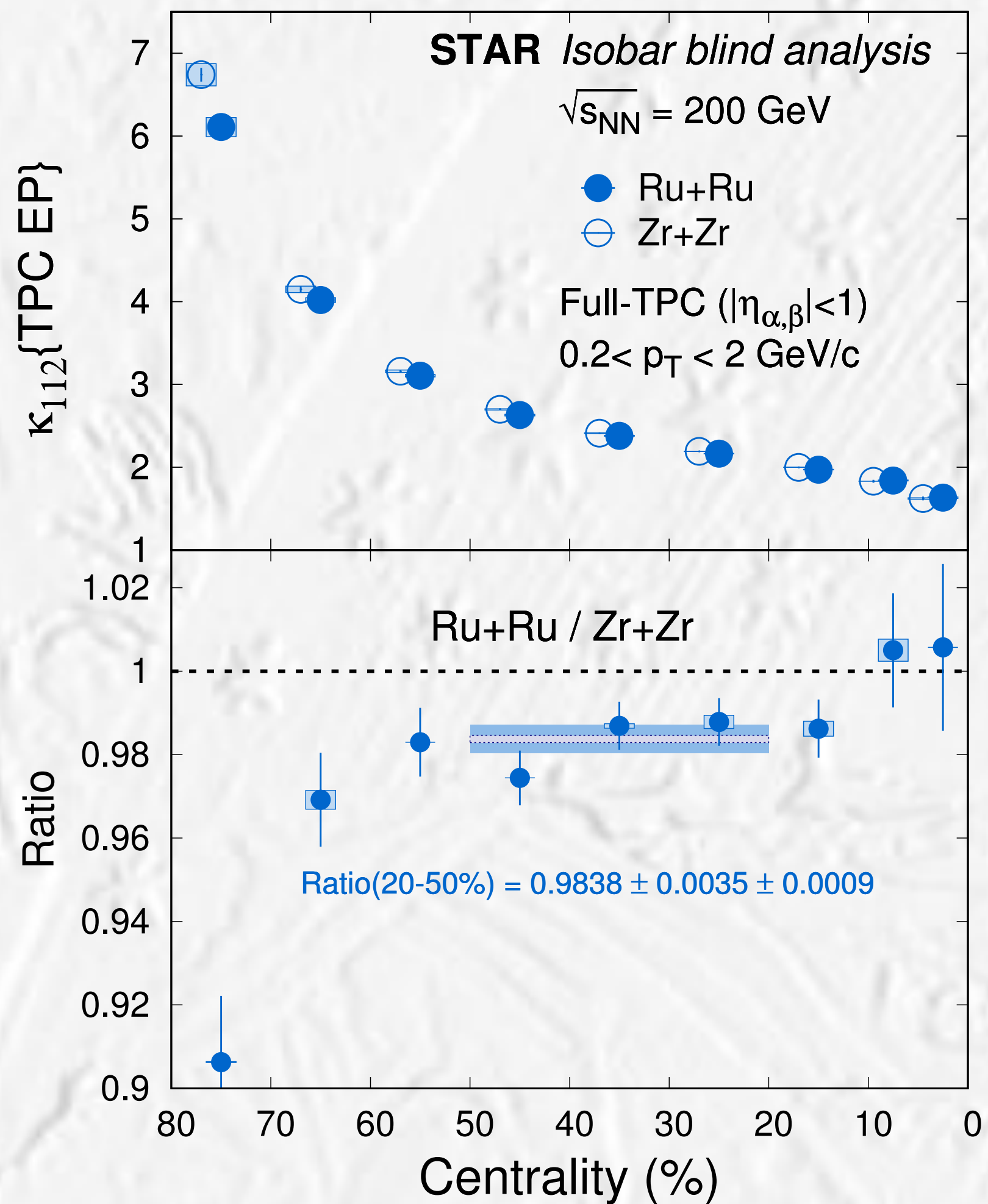
$$\frac{(\Delta\gamma/v_2)_{\text{Ru+Ru}}}{(\Delta\gamma/v_2)_{\text{Zr+Zr}}} = 1 + f_{\text{CME}}^{\text{Zr+Zr}} \left[\left(\frac{B_{\text{Ru+Ru}}}{B_{\text{Zr+Zr}}} \right)^2 - 1 \right]$$

Note:

- In most cases the calculation of double ratio does not require knowledge of the Reaction Plane resolution.
- SE (subevent) — η gap between subevents, $\Delta\gamma$ calculation in a narrower η window



Group 1



$$\delta = \langle \cos(\phi_\alpha - \phi_\beta) \rangle$$

$$= (\langle v_{1,\alpha} v_{1,\beta} \rangle + B_{IN}) + (\langle a_{1,\alpha} a_{1,\beta} \rangle + B_{OUT})$$

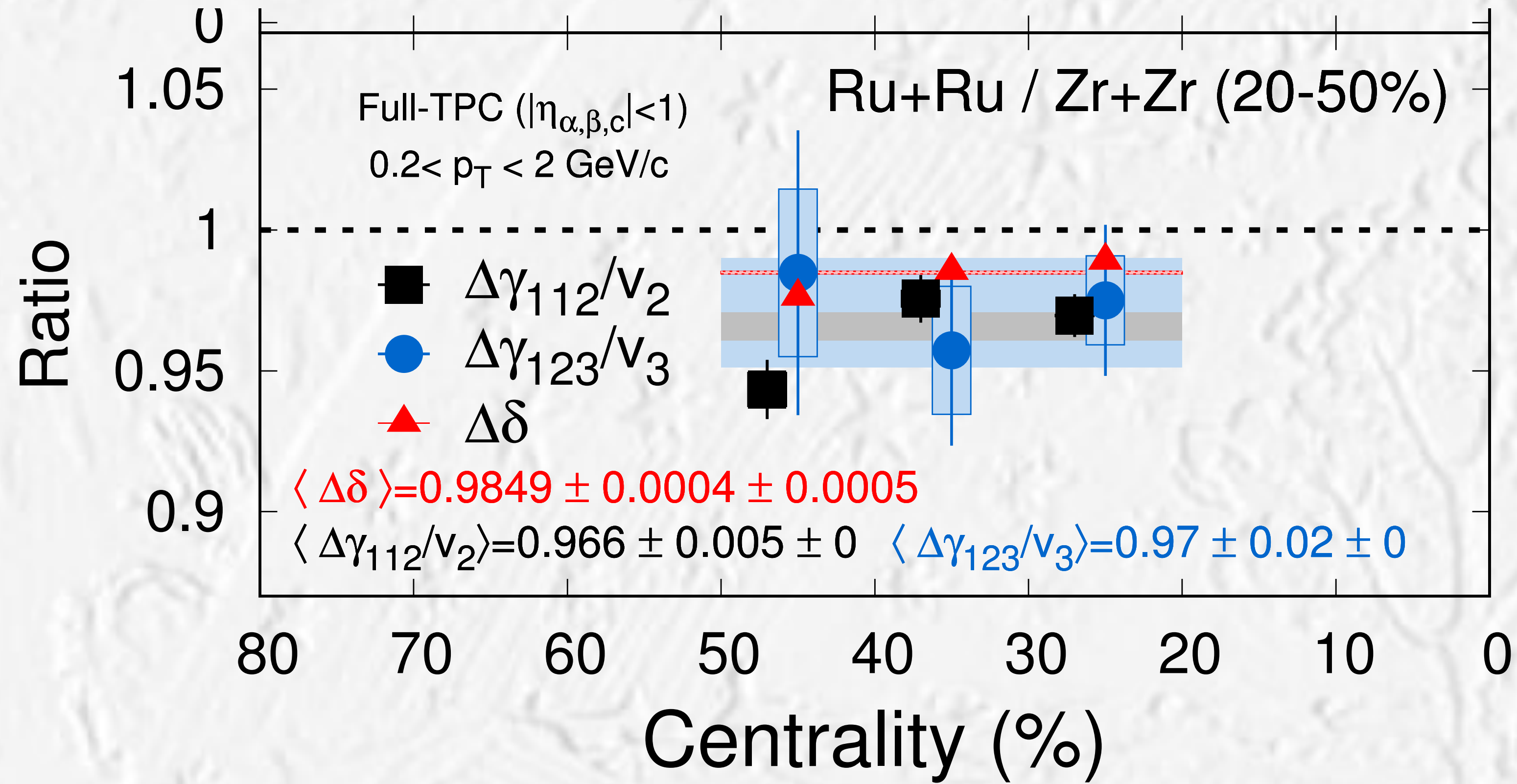
$$\kappa_{112} \equiv \frac{\Delta \gamma_{112}}{v_2 \Delta \delta}$$

$$\frac{\kappa_{112}^{Ru+Ru}}{\kappa_{112}^{Zr+Zr}} > 1$$

$$\frac{k_2^{Ru+Ru}}{k_2^{Zr+Zr}} > \frac{k_3^{Ru+Ru}}{k_3^{Zr+Zr}}$$

Expectations are not met

Mixed harmonics



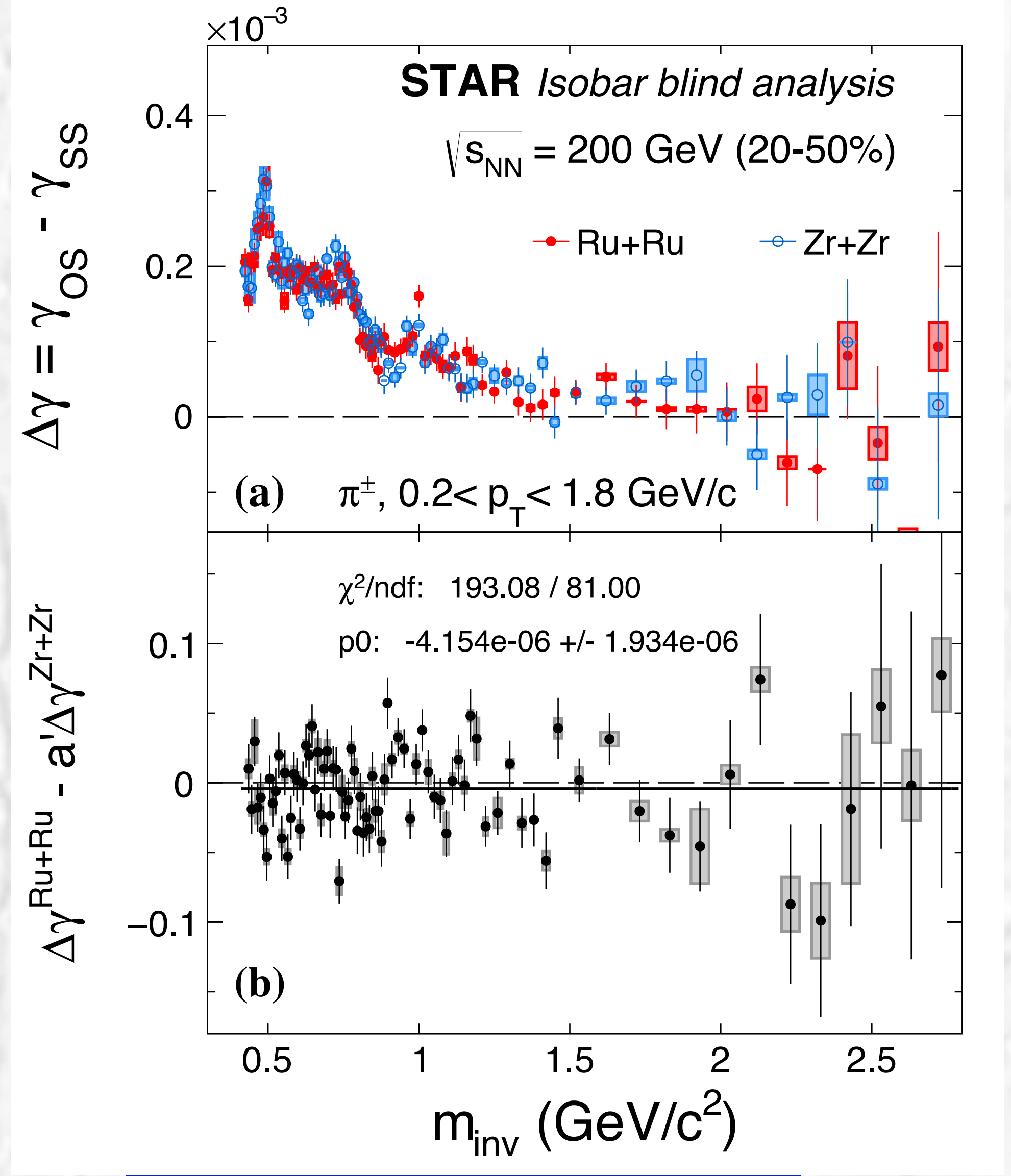
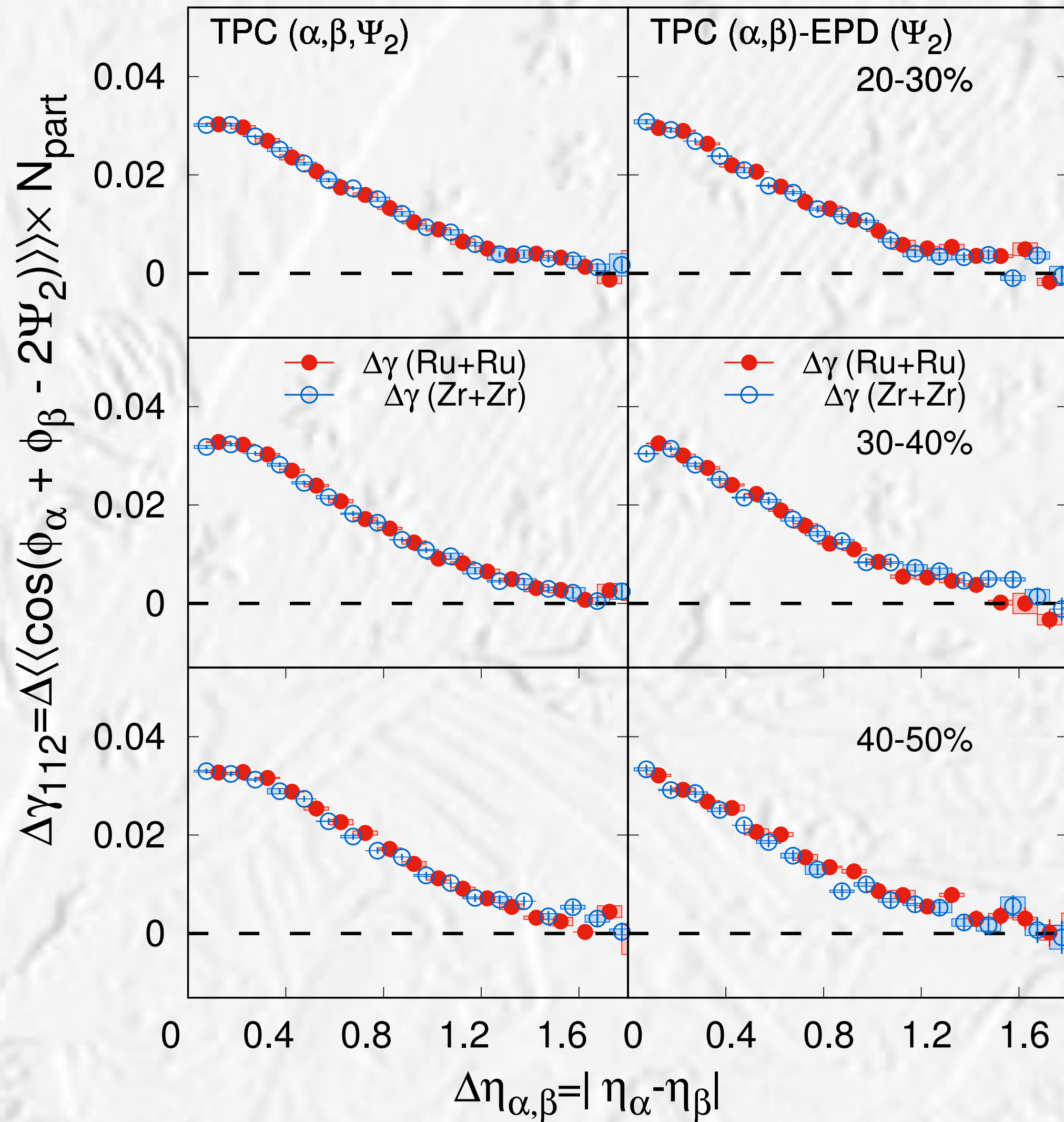
$$\gamma_{112} = \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_2) \rangle \quad \text{CME + background}$$

$$\gamma_{123} = \langle \cos(\phi_\alpha + 2\phi_\beta - 3\Psi_3) \rangle \quad \text{Only background}$$

Group 3

Group 2

STAR Isobar blind analysis, $\sqrt{s_{NN}} = 200$ GeV



No obvious enhancement anywhere

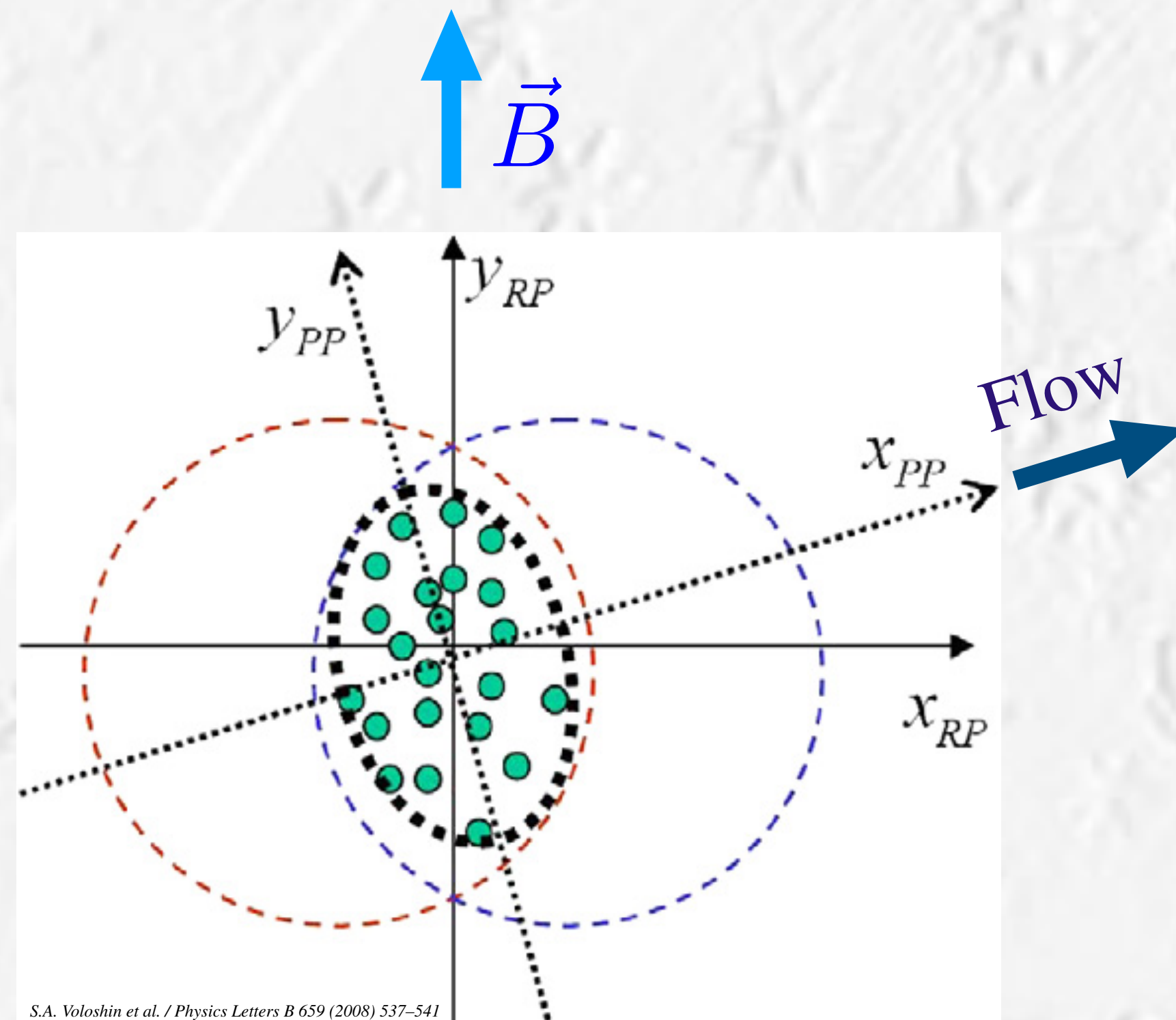


Fig. 1. The definitions of the *RP* and *PP* coordinate systems.

Assumption: spectator plane defines the magnetic field direction

Estimate of the signal from the chiral magnetic effect in heavy-ion collisions from measurements relative to the participant and spectator flow planes

Sergei A. Voloshin

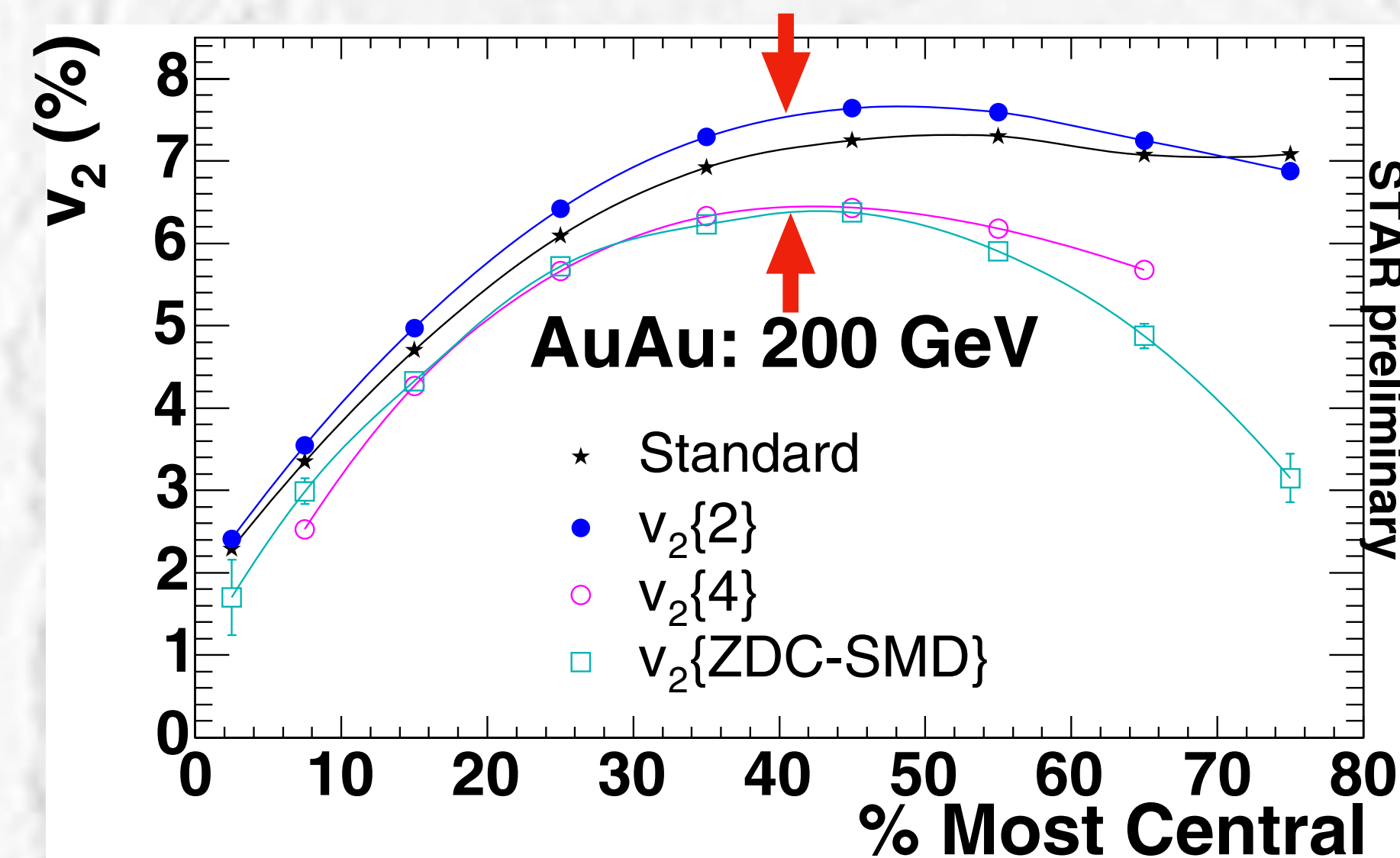
Wayne State University, 666 West Hancock, Detroit, Michigan 48201, USA

Varying the chiral magnetic effect relative to flow in a single nucleus-nucleus collision*

Hao-Jie Xu (徐浩浩)¹, Jie Zhao (赵杰)², Xiao-Bao Wang (王小保)¹, Han-Lin Li (李汉林)³, Zi-Wei Lin (林子威)^{4,5}, Cai-Wan Shen (沈彩万)¹ and Fu-Qiang Wang (王福强)^{1,2}

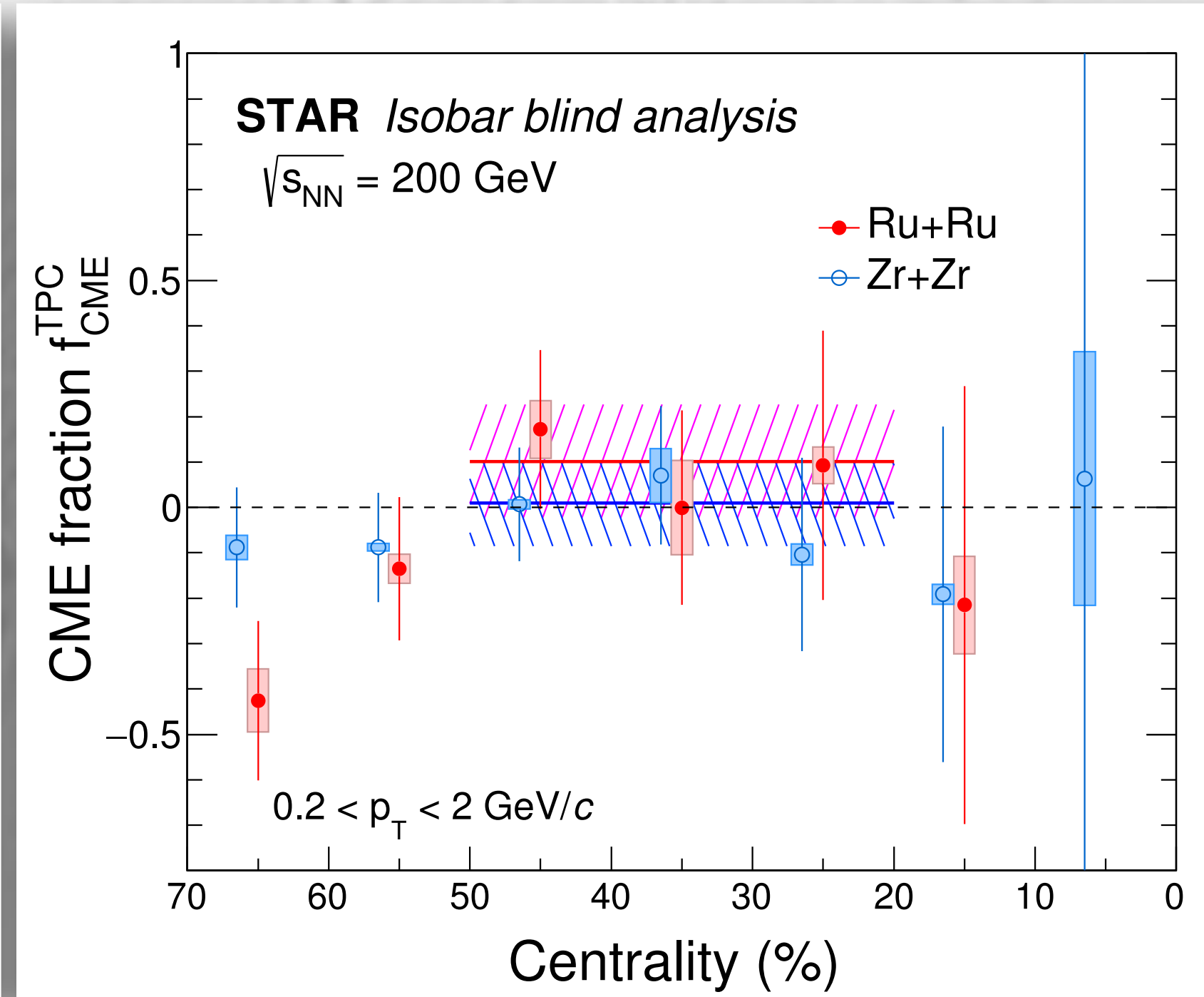
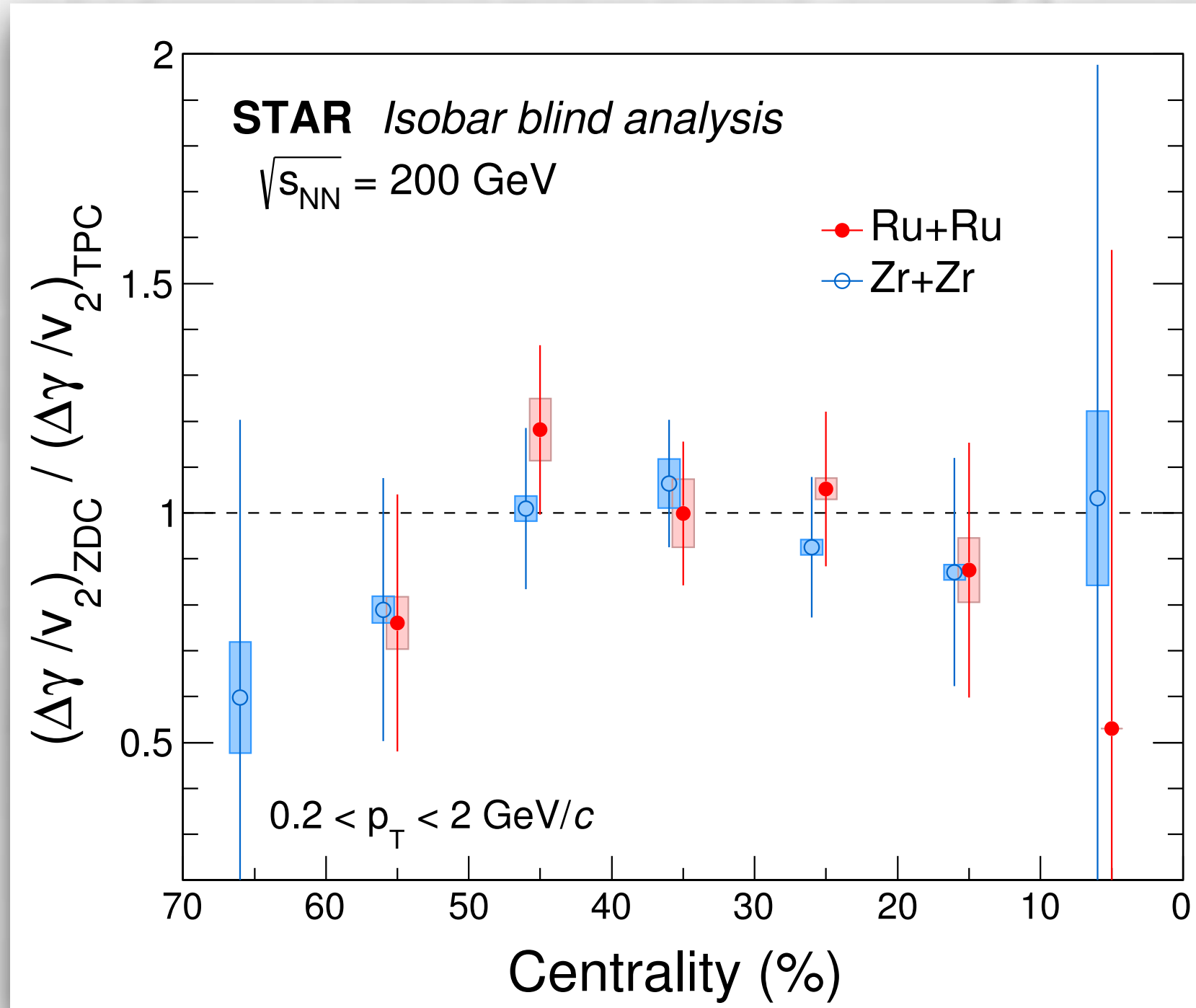
Published 1 July 2018 • © 2018 Chinese Physical Society and the Institute of High Energy

Chinese Physics C, Volume 42, Number 8



Decorrelation is strong enough to measure the difference in the CME signal

Group 4 (also done by group 3)



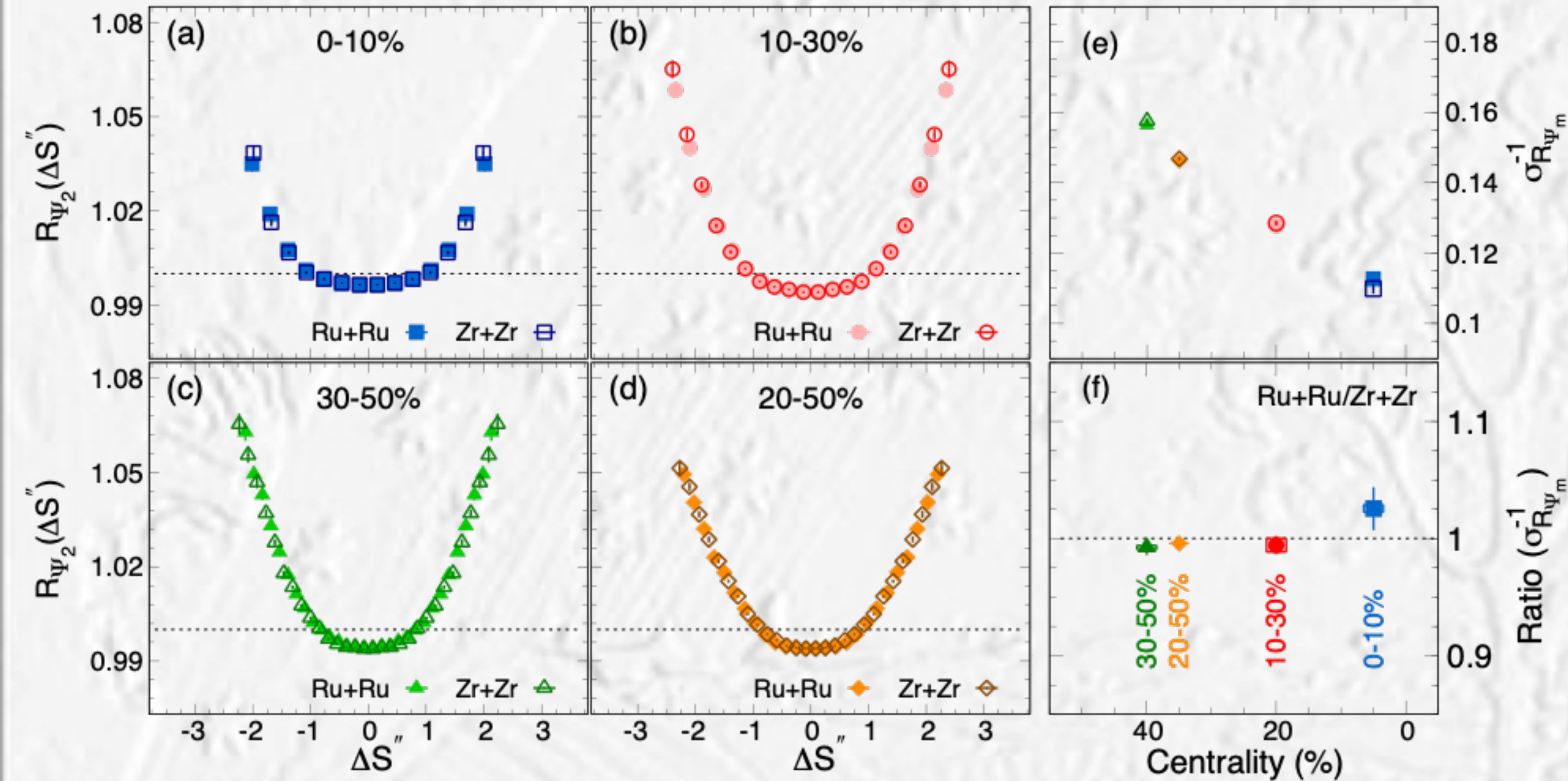
Ru:
 0.101 ± 0.123 (stat.) ± 0.023 (syst.)
 Zr:
 0.009 ± 0.088 (stat.) ± 0.033 (syst.)

$$\frac{(\Delta\gamma/v_2)_{\text{spectator}}}{(\Delta\gamma/v_2)_{\text{participant}}} = \frac{(\Delta\gamma/v_2)_{\text{ZDC}}}{(\Delta\gamma/v_2)_{\text{TPC}}} = \frac{\Delta\langle \cos(\phi_\alpha + \phi_\beta - \Psi_1^W - \Psi_1^E) \rangle / \langle \cos(2\phi - \Psi_1^W - \Psi_1^E) \rangle}{\Delta\langle \cos(\phi_\alpha + \phi_\beta - 2\phi_c) \rangle / \langle \cos(2\phi_\alpha - 2\phi_c) \rangle}.$$

$$\frac{(\Delta\gamma/v_2)_{\text{ZDC}}}{(\Delta\gamma/v_2)_{\text{TPC}}} = 1 + f_{\text{CME}}^{\text{TPC}} \left(\frac{v_2^2\{\text{TPC}\}}{v_2^2\{\text{ZDC}\}} - 1 \right)$$

Group 5

STAR Isobar blind analysis, $\sqrt{s_{NN}}=200$ GeV



$$R_{\Psi_2}(\Delta S) = C_{\Psi_2}(\Delta S) / C_{\Psi_2}^{\perp}(\Delta S),$$

$$C_{\Psi_2}(\Delta S) = \frac{N_{\text{real}}(\Delta S)}{N_{\text{shuffled}}(\Delta S)},$$

$$\Delta S = \frac{\sum_1^{n^+} w_i^+ \sin(\Delta\varphi_2)}{\sum_1^{n^+} w_i^+} - \frac{\sum_1^{n^-} w_i^- \sin(\Delta\varphi_2)}{\sum_1^{n^-} w_i^-},$$

Predefined CME signature:

$$1/\sigma_{R_{\Psi_2}}(\text{Ru} + \text{Ru}) > 1/\sigma_{R_{\Psi_2}}(\text{Zr} + \text{Zr})$$

No CME signature that satisfies the predefined criteria observed

For the relation to $\Delta\gamma$ see S. Choudhury (Fudan U.), et al. e-Print: 2105.06044

Blind analysis summary plot

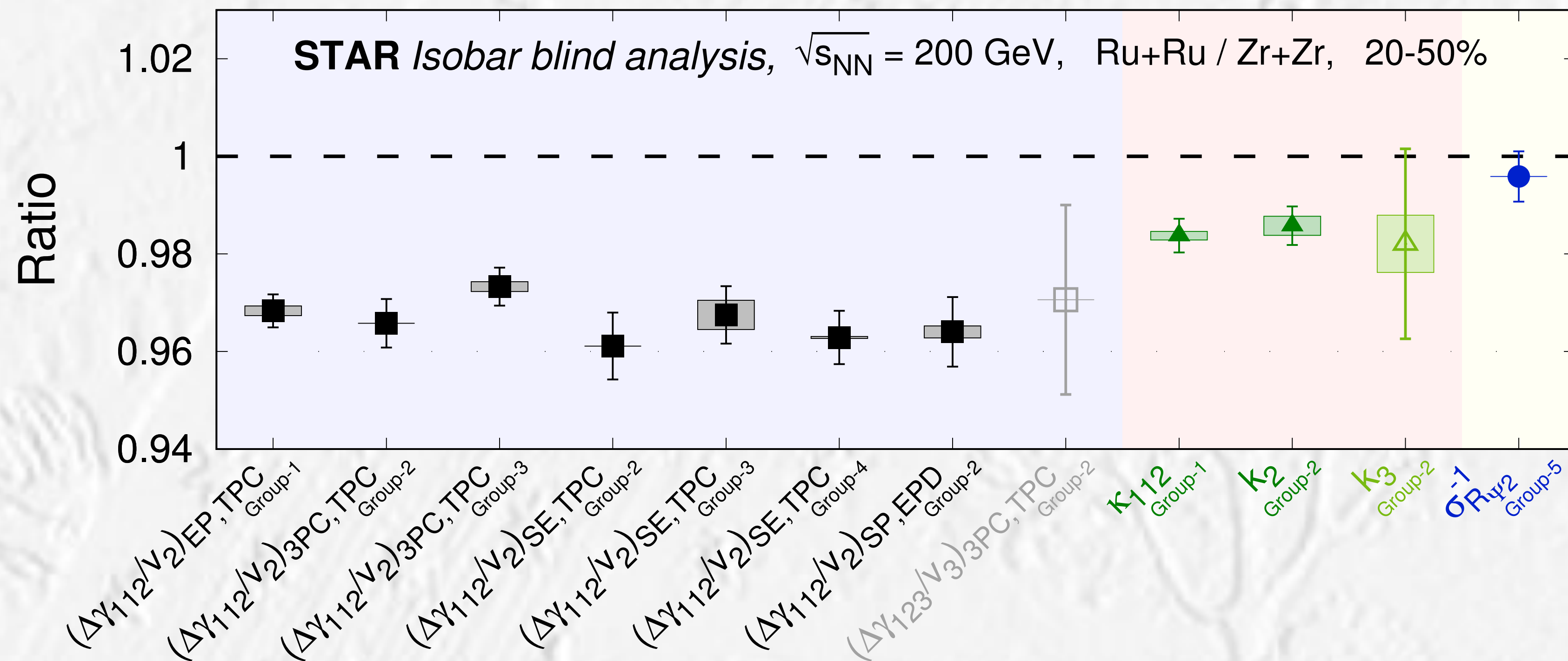
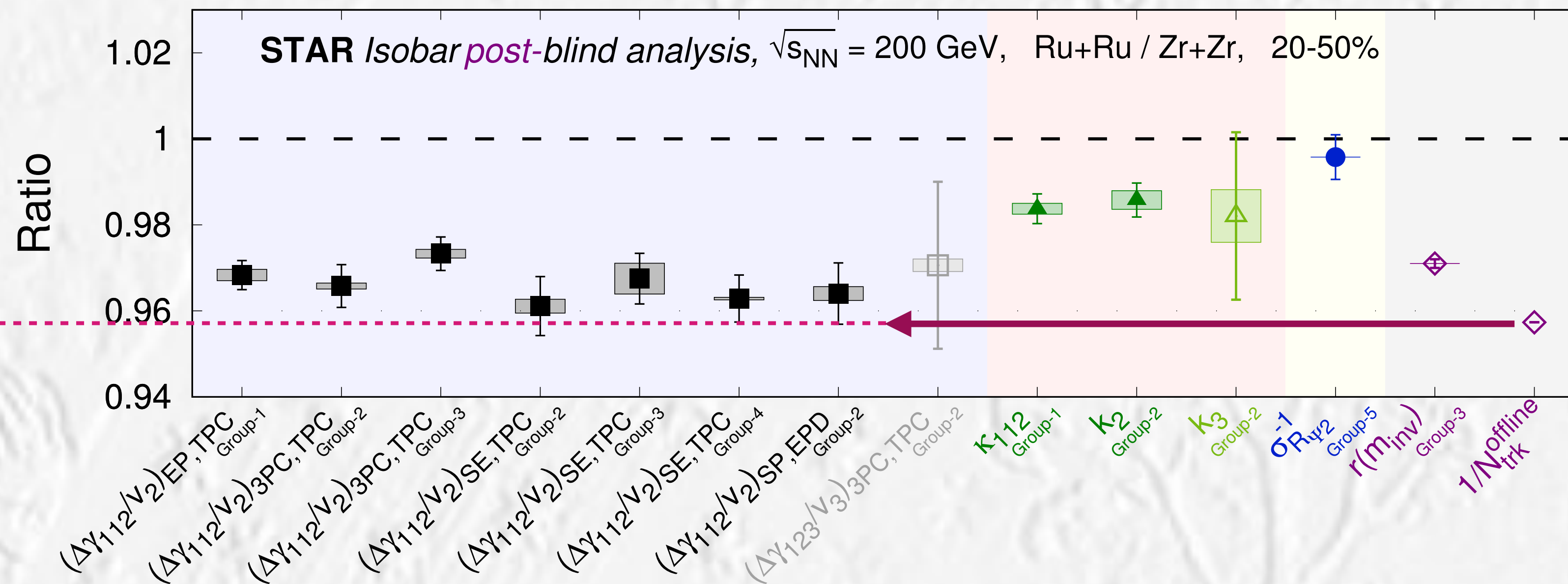


FIG. 26. Compilation of results from the blind analysis. Only results contrasting between the two isobar systems are shown. Results are shown in terms of the ratio of measures in Ru+Ru collisions over Zr+Zr collisions. Solid dark symbols show CME-sensitive measures whereas open light symbols show counterpart measures that are supposed to be insensitive to CME. The vertical lines indicate statistical uncertainties whereas boxes indicate systematic uncertainties. The colors in the background are intended to separate different types of measures. The fact that CME-sensitive observable ratios lie below unity leads to the conclusion that no predefined CME signatures are observed in this blind analysis.

No CME signature that satisfies the predefined criteria observed

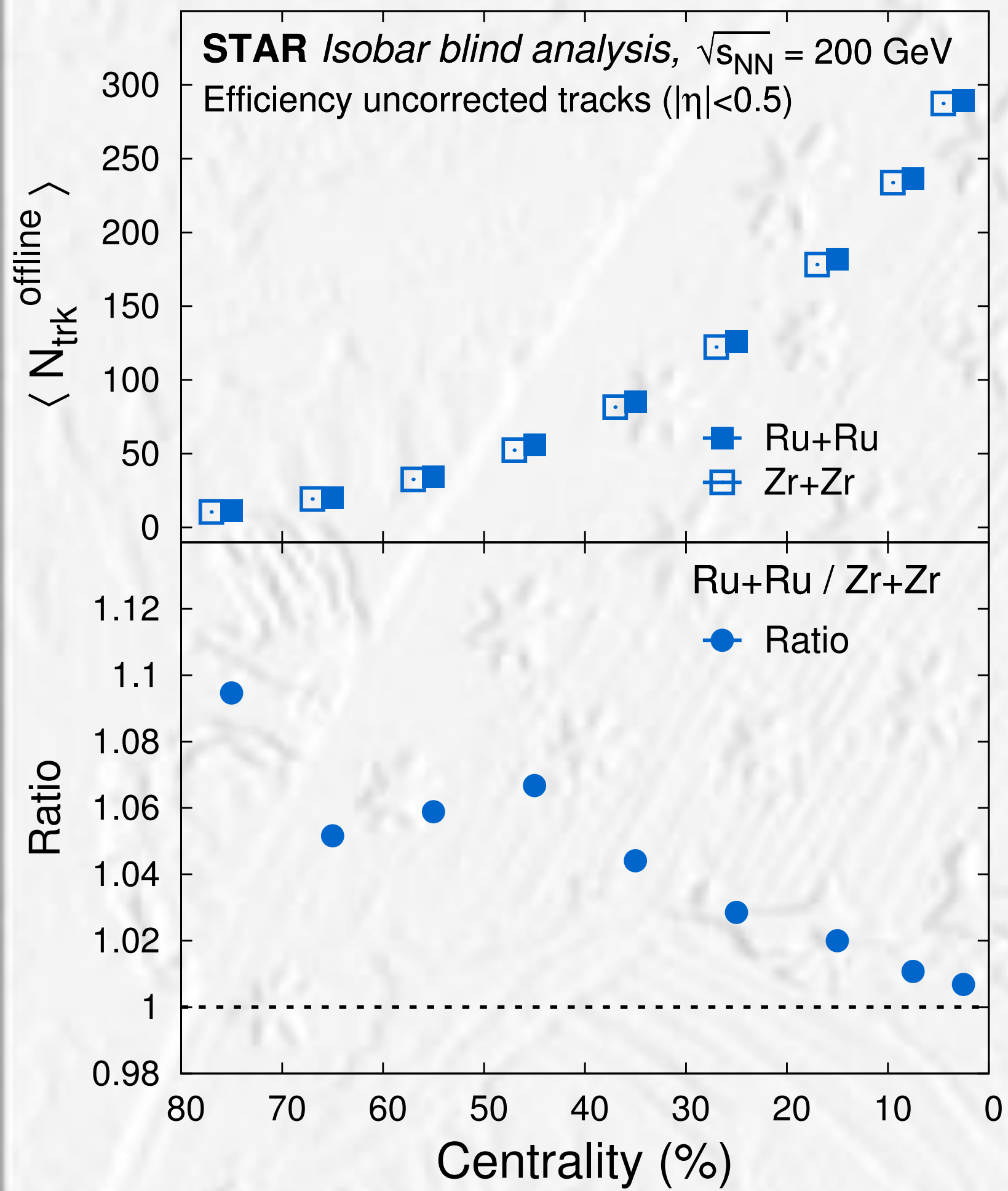
Summary plot (post-blinding)



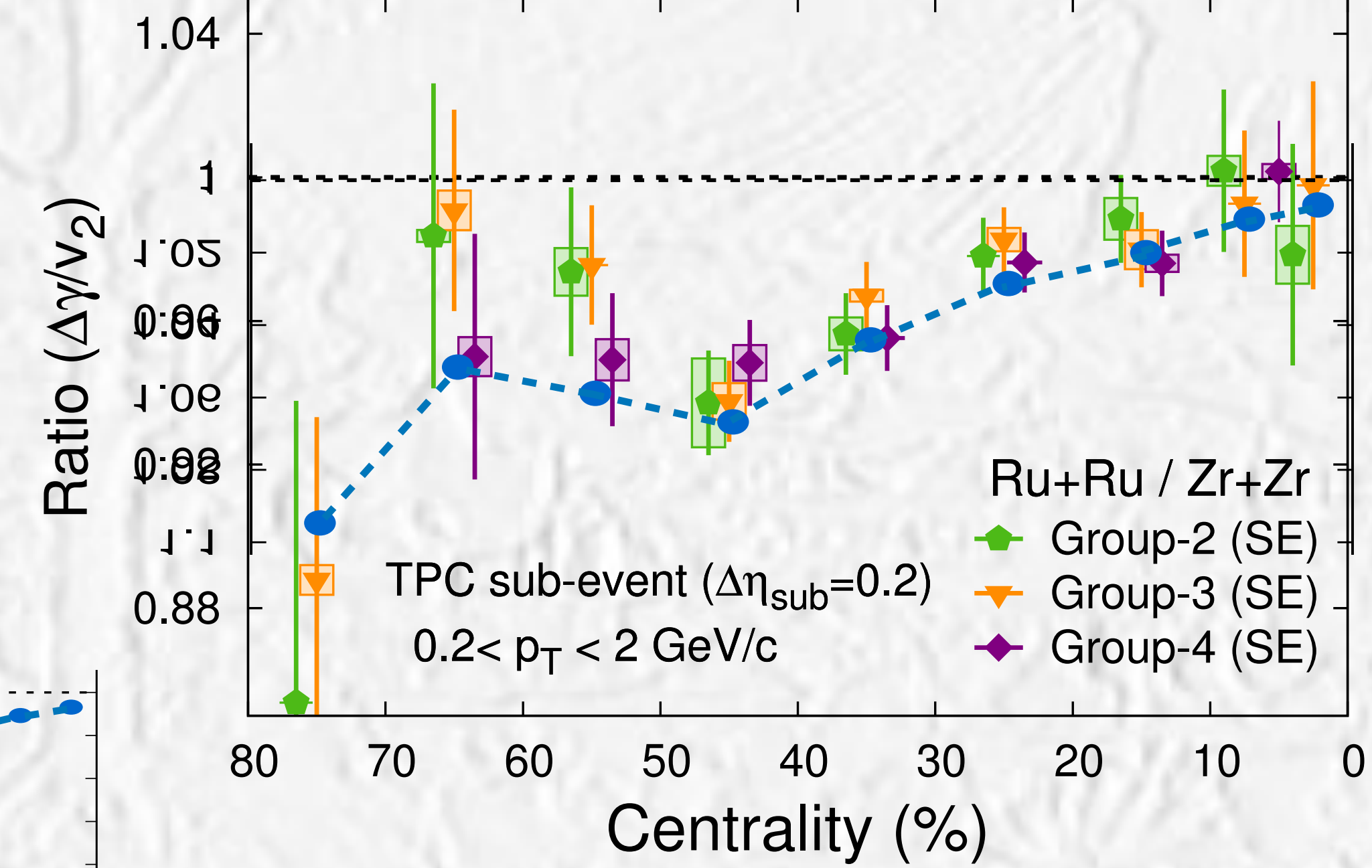
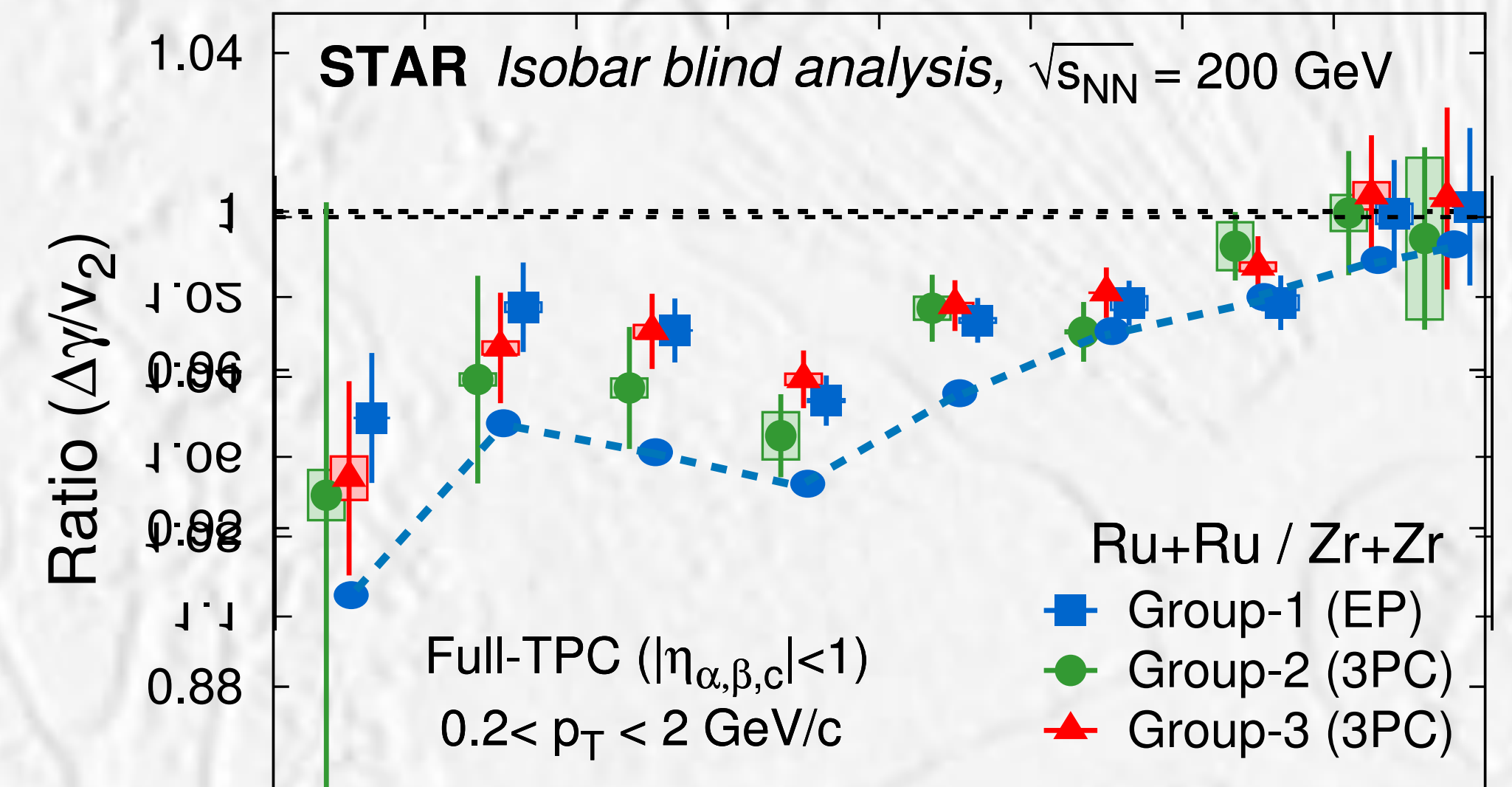
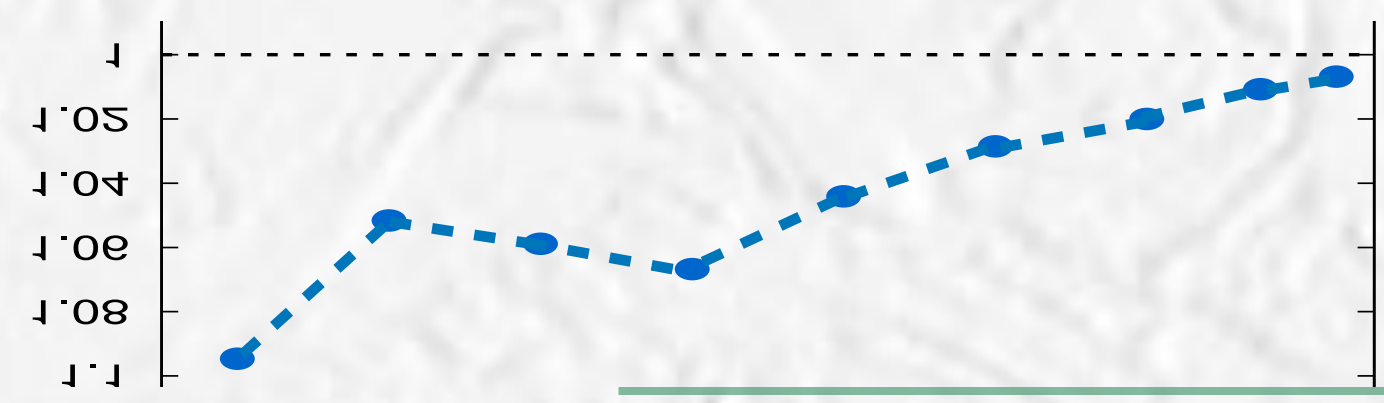
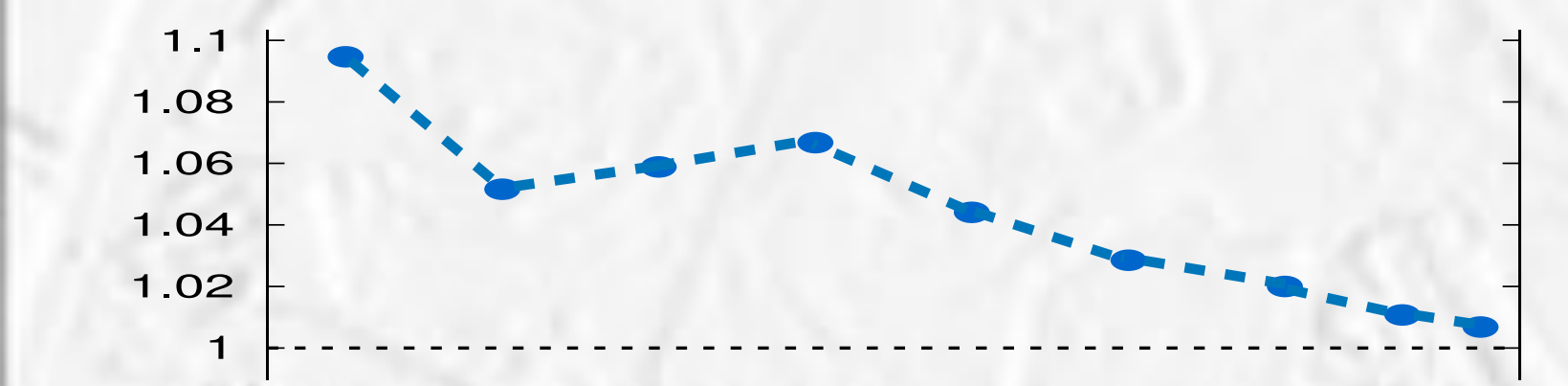
Two most right points added for post-blinding discussion

FIG. 27. Compilation of post-blinding results. This figure is largely the same as Fig. 26 with the following differences: numerical changes in the results from the new run-by-run QA algorithm are treated as an additional systematic uncertainty added in quadrature, and two data points (open markers) have been added on the right to indicate the ratio of inverse multiplicities ($N_{trk}^{offline}$) and the ratio of relative pair multiplicity difference (r) as explained in the text.

Any two particle correlation due to small clusters scale as 1/multiplicity !
 A better comparison might be for $[N_{ch}^{Ru}(\Delta\gamma/v_2)_{Ru}] / [N_{ch}^{Zr}(\Delta\gamma/v_2)_{Zr}]$



Cut,
Flip,
Overlay



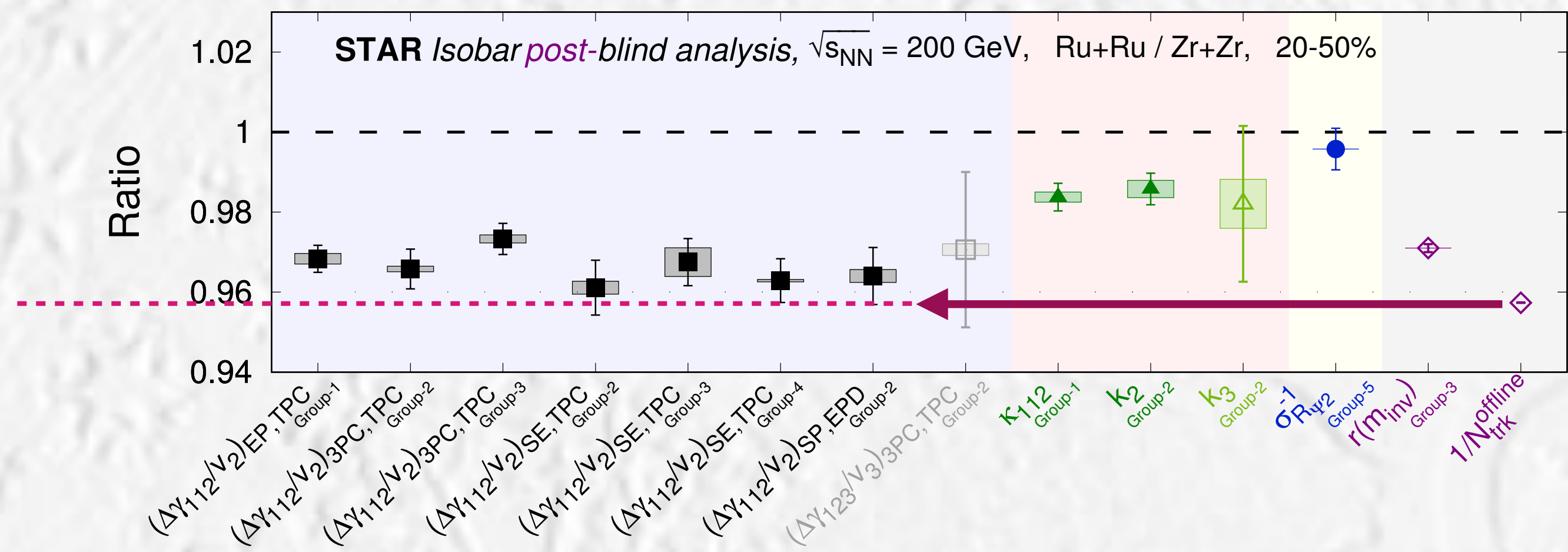
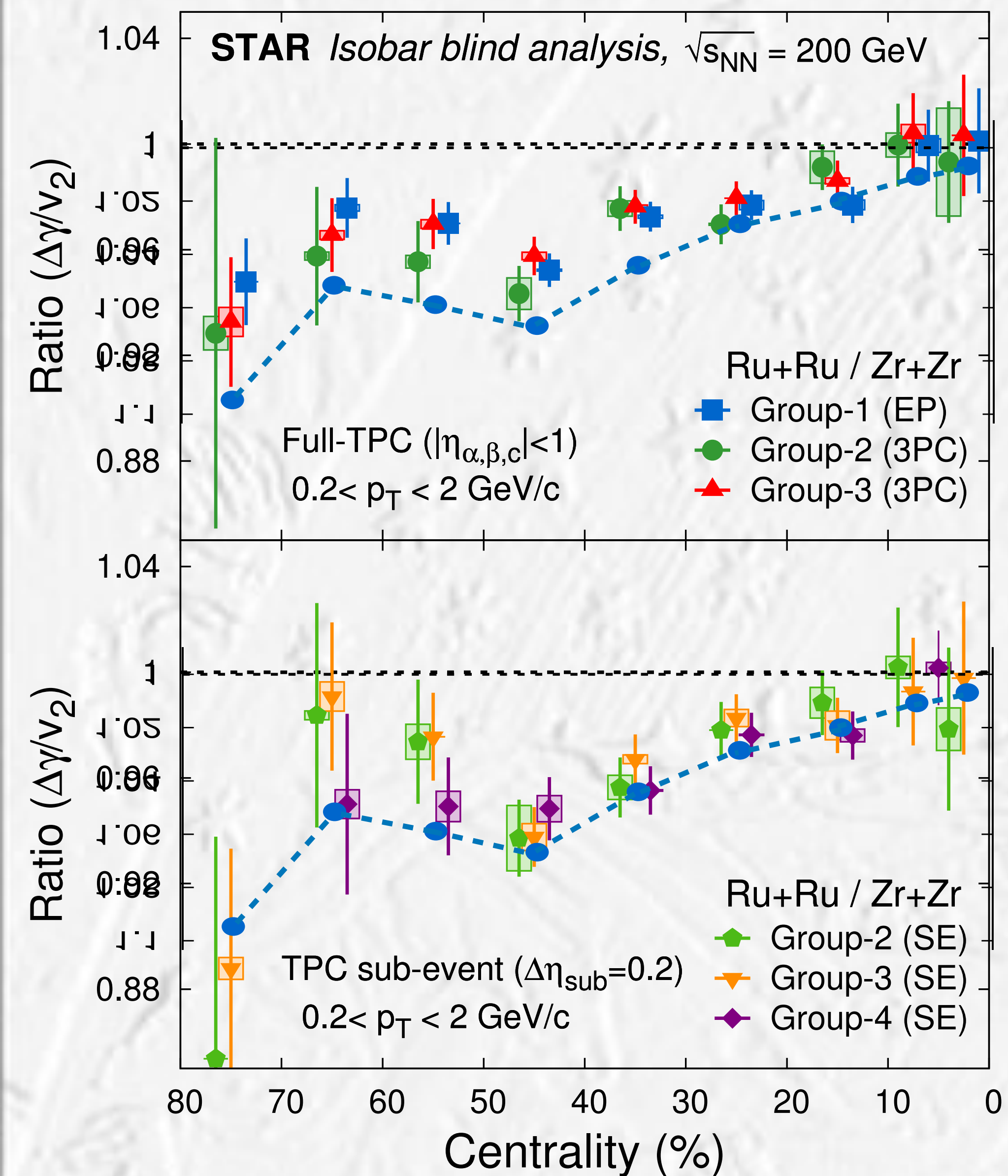


FIG. 27. Compilation of post-blinding results. This figure is largely the same as Fig. 26 with the following differences: numerical changes in the results from the new run-by-run QA algorithm are treated as an additional systematic uncertainty added in quadrature, and two data points (open markers) have been added on the right to indicate the ratio of inverse multiplicities ($N_{trk}^{offline}$) and the ratio of relative pair multiplicity difference (r) as explained in the text.

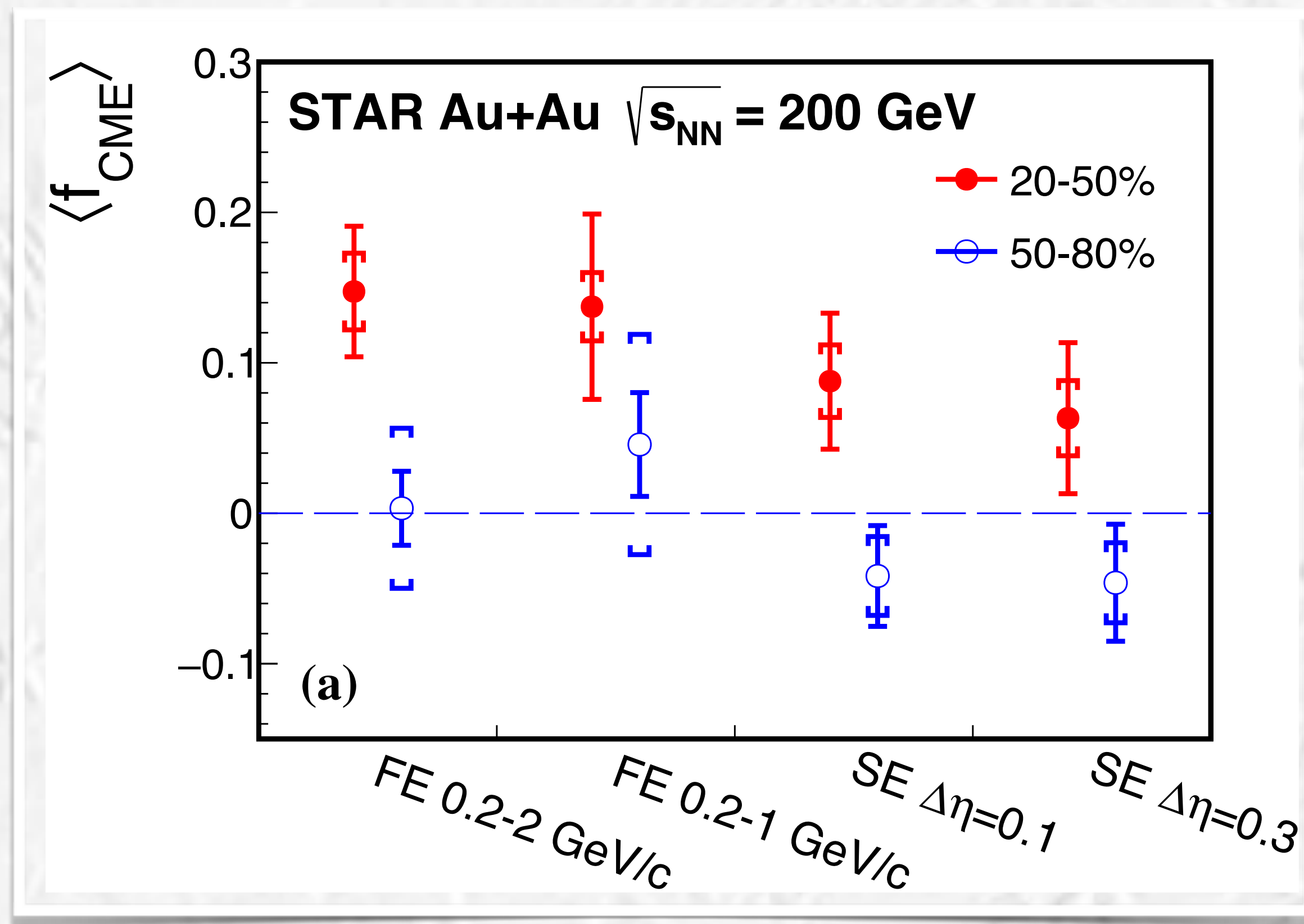
To establish exact limits, one need to resolve/understand systematics in the ratio

$$\frac{(\Delta\gamma/v_2)_{Ru+Ru}}{(\Delta\gamma/v_2)_{Zr+Zr}}$$

up to a (sub)percent level (note difference between results from different groups, “Full” vs “SE”).

Search for the chiral magnetic effect via charge-dependent azimuthal correlations relative to spectator and participant planes in Au+Au collisions at $\sqrt{s_{\text{NN}}} = 200$ GeV

STAR Collaboration arXiv:2106.09243v1 [nucl-ex] 17 Jun 2021



Note:

-SE (subevent) — η gap between subevents,
 $\Delta\gamma$ calculation in a narrower η window

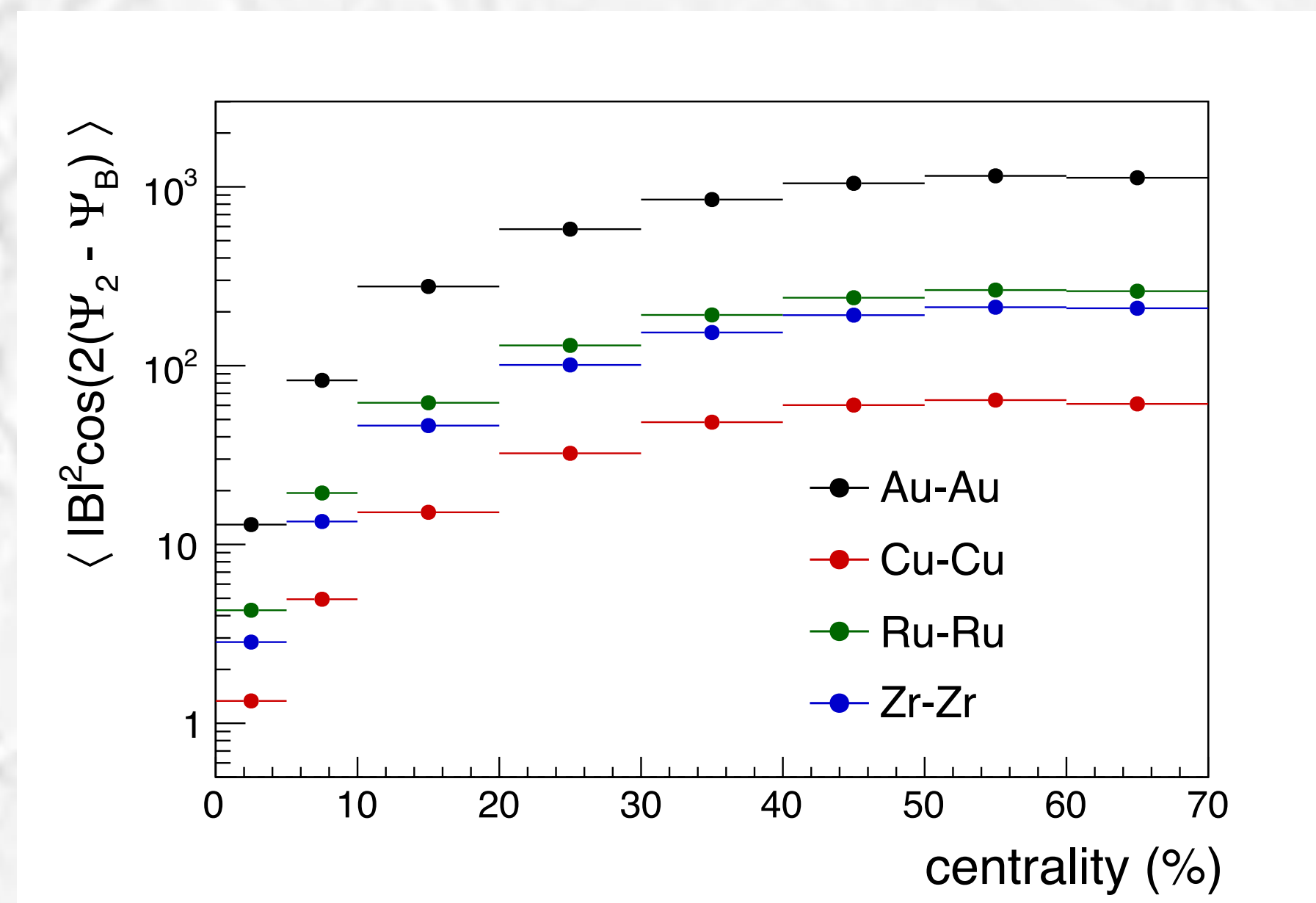
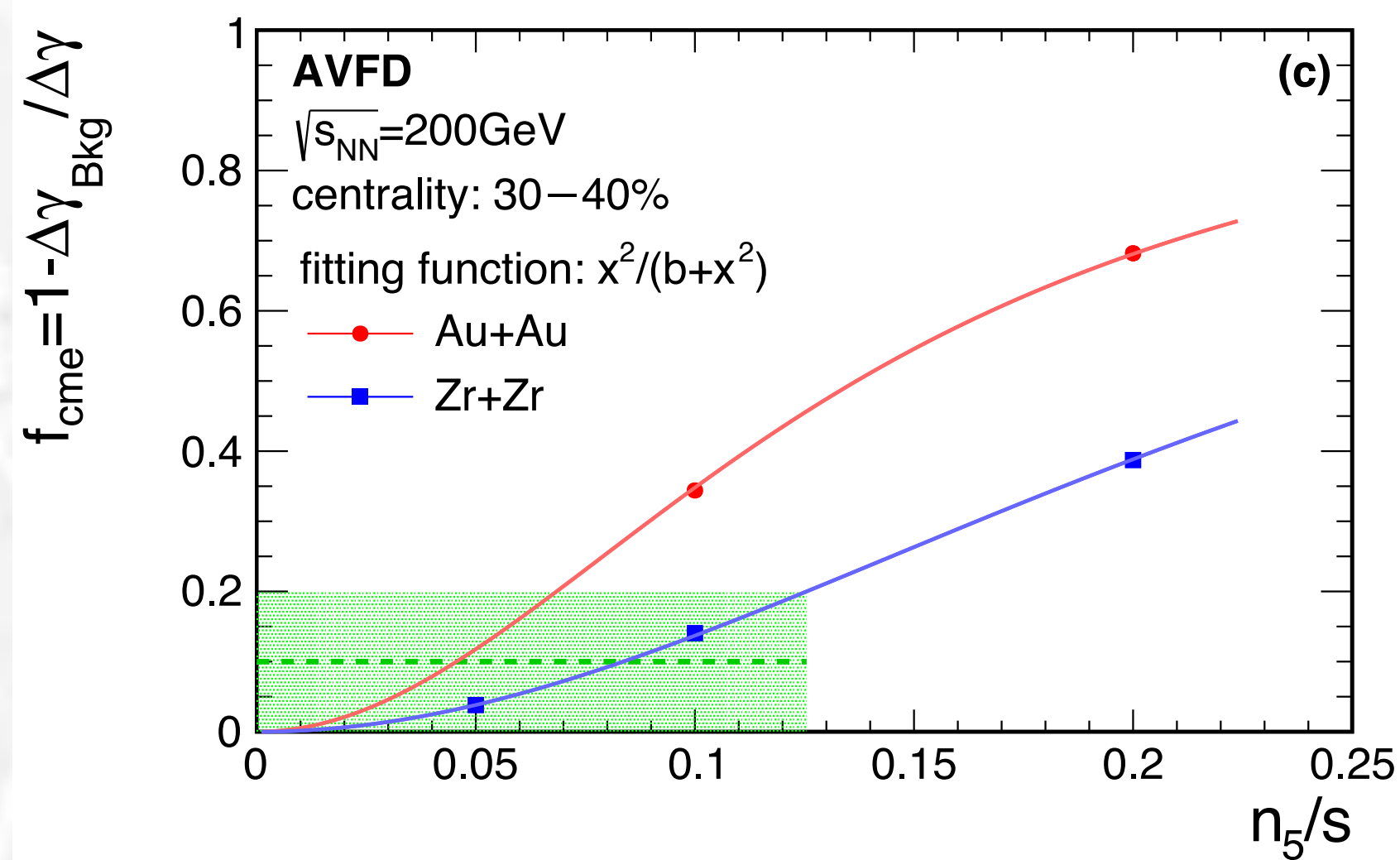
Error bars show statistical uncertainties; the caps indicate the systematic uncertainties.

STAR isobars:

- No CME signature that satisfies the predefined criteria has been observed in isobar collisions in this blind analysis.
- Accurate upper limits for f_{CME} are being evaluated.

Isobar results do not exclude a bigger signal in AuAu.
The signal could be significantly smaller in such (relatively small nuclei) collisions

Y. Feng, Y. Lin, J. Zhao, and F. Wang, Phys. Lett. B **820**, 136549 (2021),



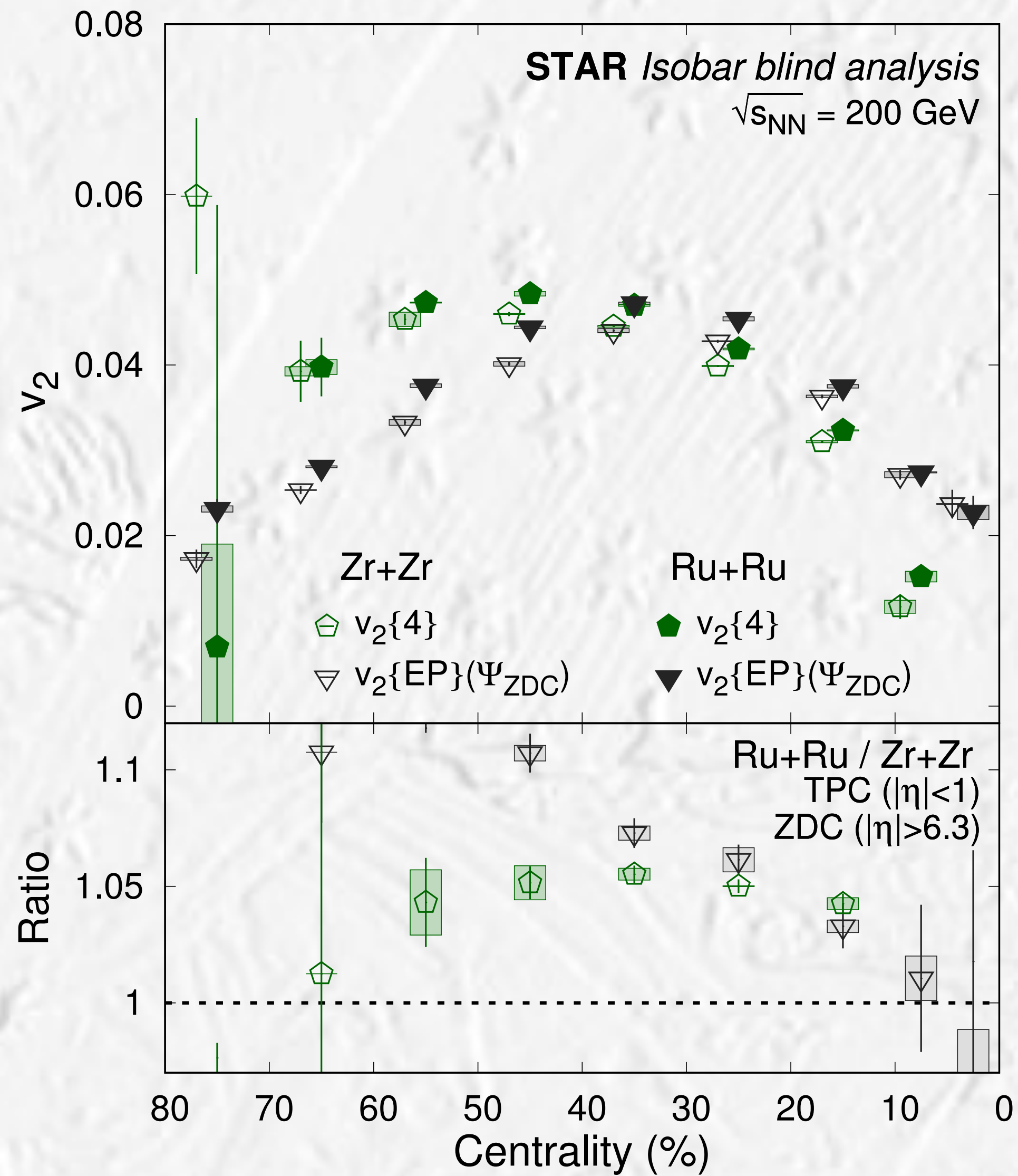
The signal could depend strongly on the system size.
Calculations by A. Dobrin (private communication)

Isobar run was a real success (not only for the CME search)
Should we request for more? $^{136}_{54}\text{Ce}$, $^{136}_{50}\text{Xe}$?

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EXTRA SLIDES

$v_2\{2\}, v_2\{\text{ZDC}\}$



Search for the Chiral Magnetic Effect with Isobar Collisions at $\sqrt{s_{\text{NN}}} = 200$ GeV by the STAR Collaboration at RHIC

STAR Collaboration arXiv:2109.00131v1 [nucl-ex] 1 Sep 2021

The chiral magnetic effect (CME) is predicted to occur as a consequence of a local violation of \mathcal{P} and \mathcal{CP} symmetries of the strong interaction amidst a strong electro-magnetic field generated in relativistic heavy-ion collisions. Experimental manifestation of the CME involves a separation of positively and negatively charged hadrons along the direction of the magnetic field. Previous measurements of the CME-sensitive charge-separation observables remain inconclusive because of large background contributions. In order to better control the influence of signal and backgrounds, the STAR Collaboration performed a blind analysis of a large data sample of approximately 3.8 billion isobar collisions of ${}^{96}_{44}\text{Ru}+{}^{96}_{44}\text{Ru}$ and ${}^{96}_{40}\text{Zr}+{}^{96}_{40}\text{Zr}$ at $\sqrt{s_{\text{NN}}} = 200$ GeV. Prior to the blind analysis, the CME signatures are predefined as a significant excess of the CME-sensitive observables in Ru+Ru collisions over those in Zr+Zr collisions, owing to a larger magnetic field in the former. A precision down to 0.4% is achieved, as anticipated, in the relative magnitudes of the pertinent observables between the two isobar systems. Observed differences in the multiplicity and flow harmonics at the matching centrality indicate that the magnitude of the CME background is different between the two species. No CME signature that satisfies the predefined criteria has been observed in isobar collisions in this blind analysis.