

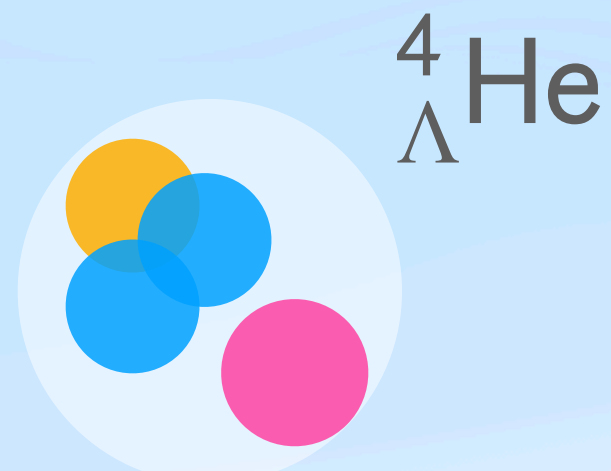
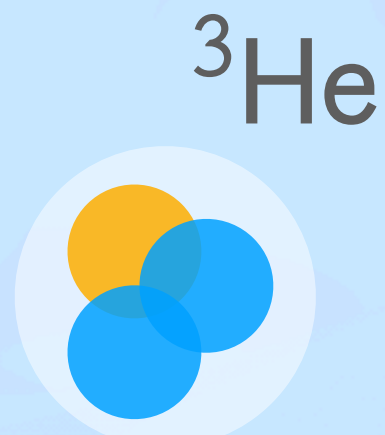
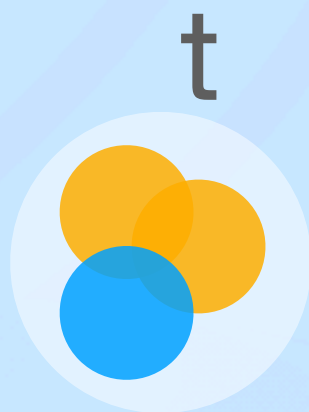


# New Hypernuclei Measurements from STAR

n

p

$\Lambda$

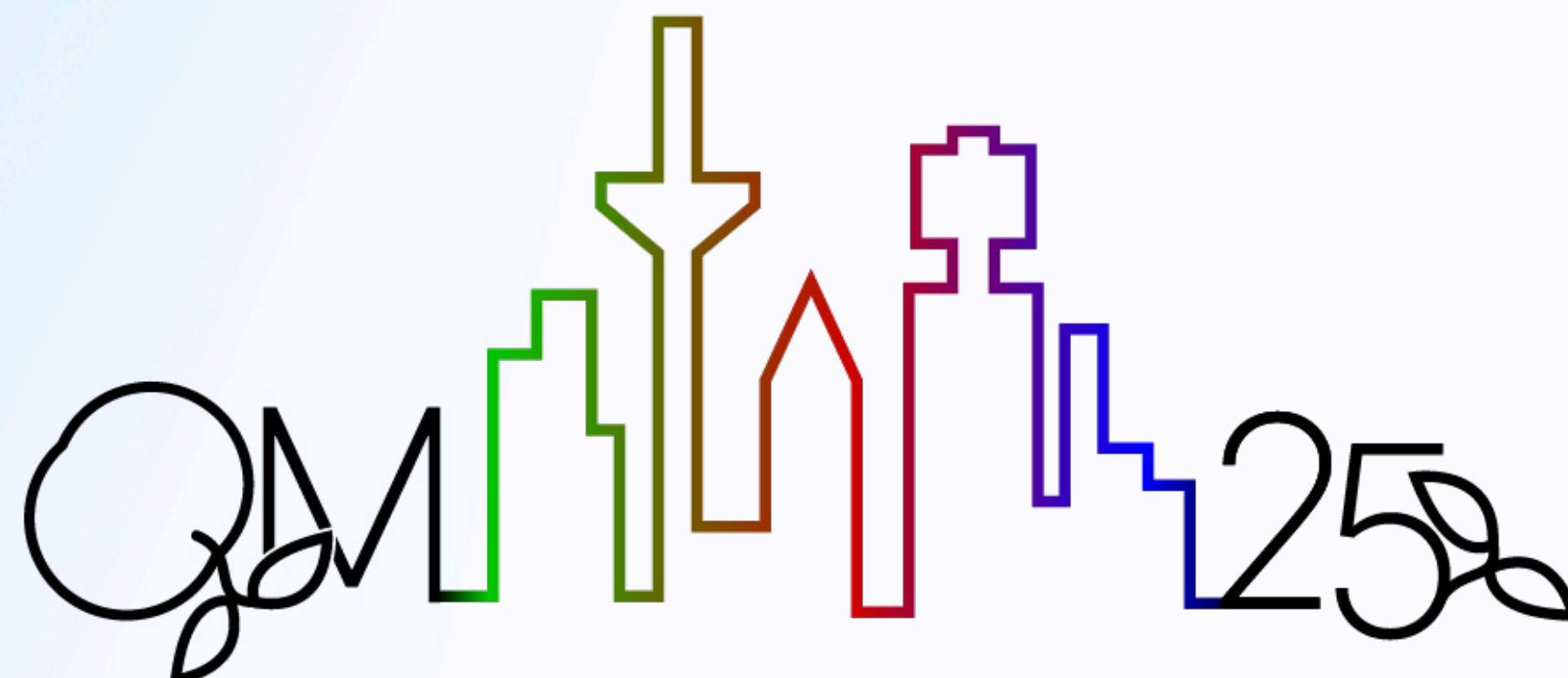


## Outline

1. Introduction
2. Particle Yields
3. Transverse Momentum Distribution
4. Collective Flow
5. Summary and Outlook

*Yingjie Zhou for the STAR Collaboration  
CCNU, GSI*

*Quark Matter 2025, Frankfurt, Germany*



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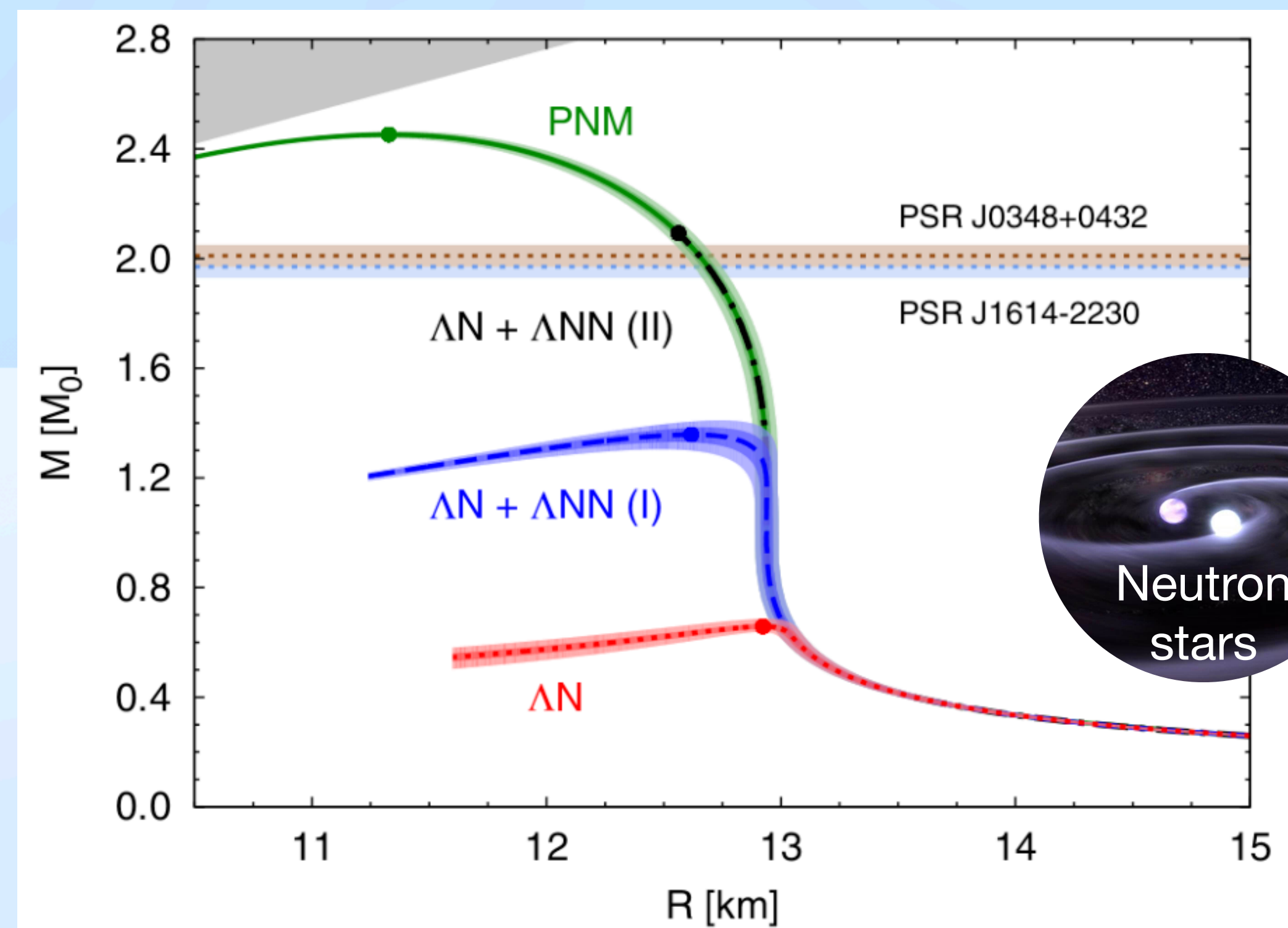




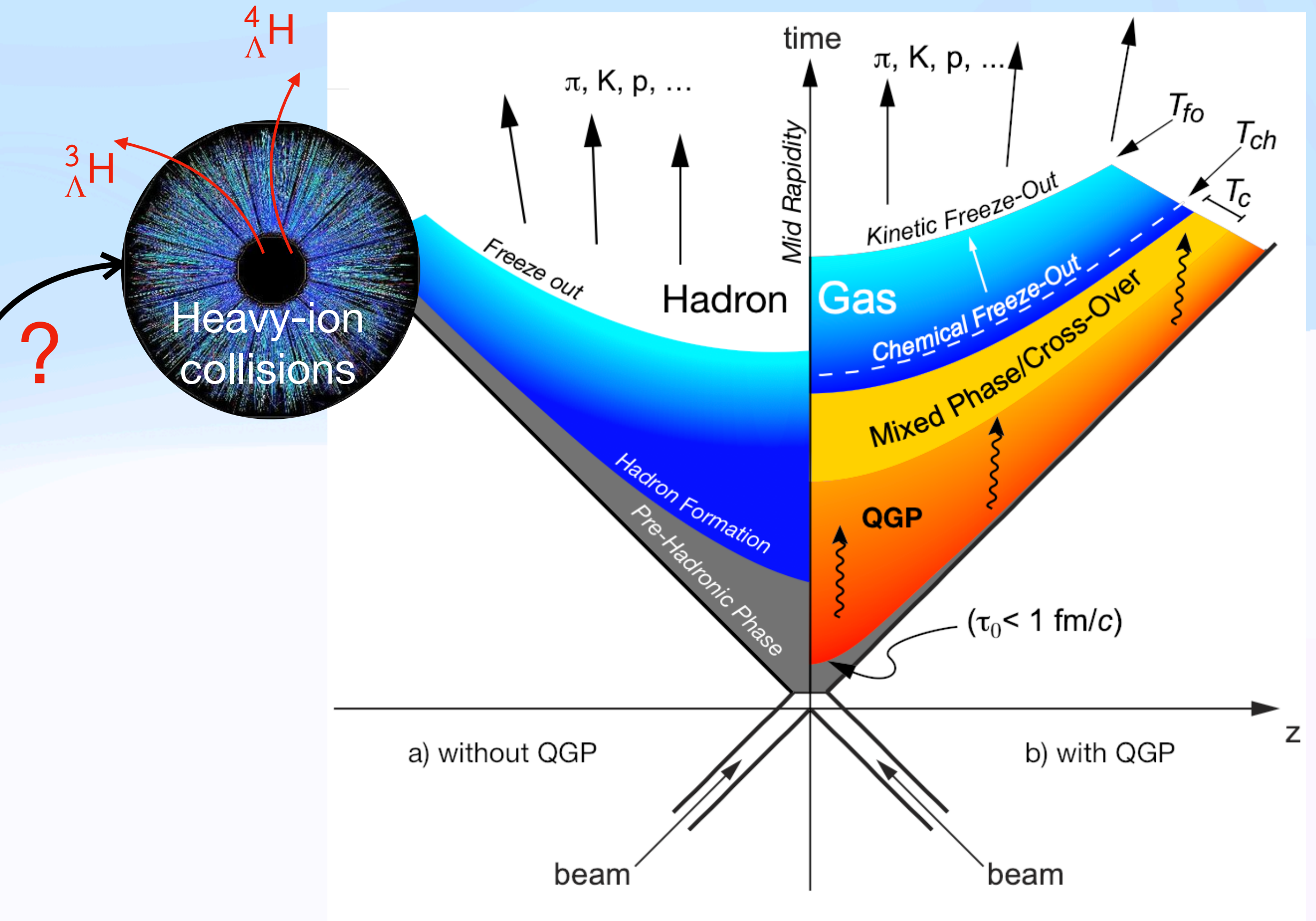
# Hypernuclei and Hyperon-Nucleon (Y-N) Interaction

- **Hyperon Puzzle:** difficulty to reconcile the measured masses of neutron stars with the presence of hyperons in their interiors

D. Lonardoni et al., PRL 114, 092301 (2015)



- Hypernuclei have been measured in heavy-ion collisions over a broad range of baryon densities



- Density dependent YN, YNN interactions are essential for solving the hyperon puzzle

- Can hypernuclei production be used to constrain the in-medium Y-N interaction?



# Hypernuclei Production Mechanisms

When are nuclei produced in a heavy-ion collision?

## 1. Thermal models

- Hadrons and (hyper-)nuclei are treated equally
- Yields are predicted with thermal equilibrium assumptions

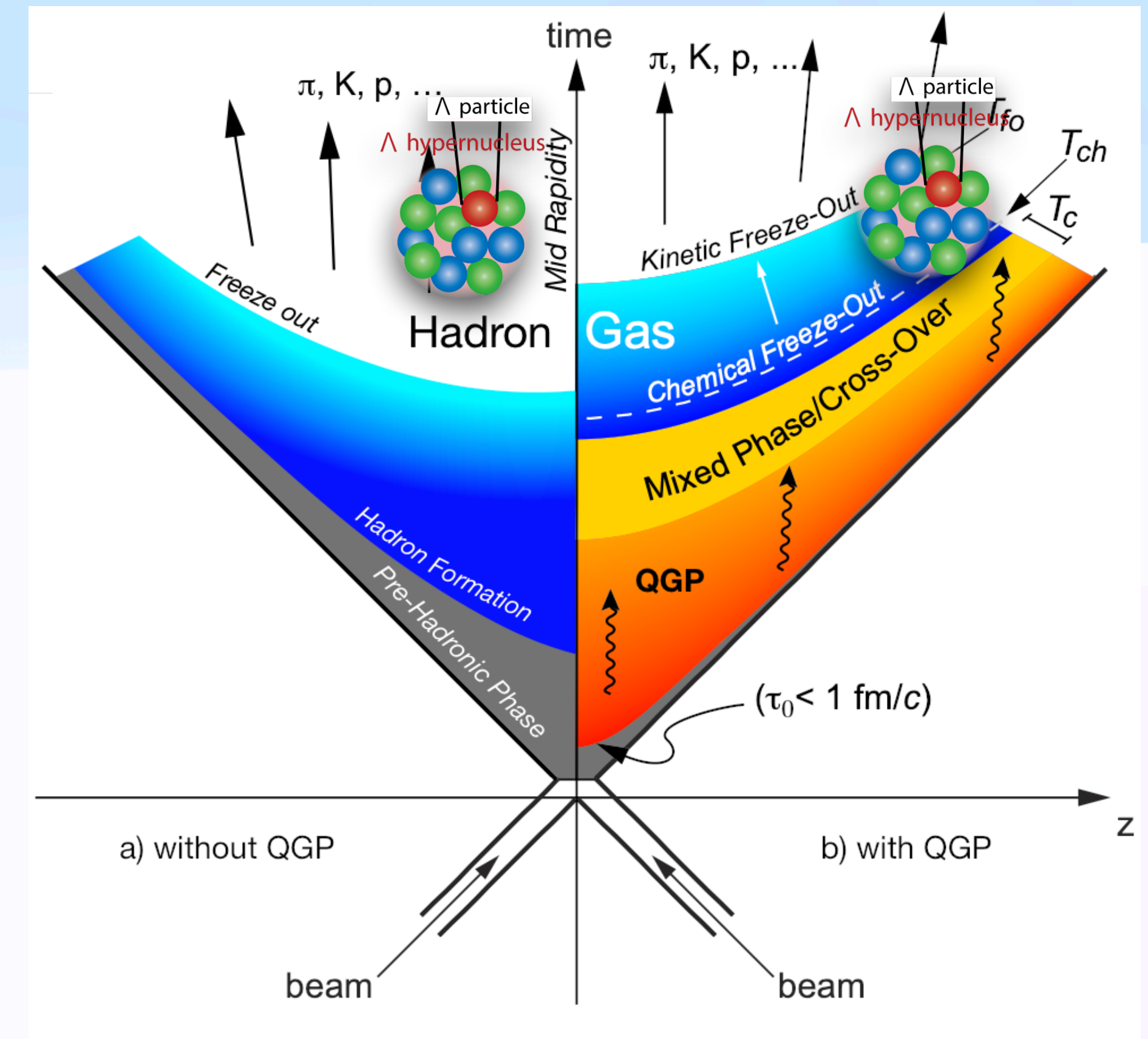
✗ Y–N dynamics have minimal impact on final yield

## 2. Coalescence model

- (Hyper-)nuclei formation after kinetic freeze-out
- Nucleon coalescence
  - Wigner function
  - Emission source size and nuclear radius

✓ In-medium Y–N interactions modify freeze-out phase space, affecting coalescence and hypernuclei yields

*Need a solid understanding in **hypernuclei production mechanisms** before we can use them to probe the in-medium Y–N interaction*



# Hypernuclei Production Mechanisms

When are nuclei produced in a heavy-ion collision?

## 1. Thermal models

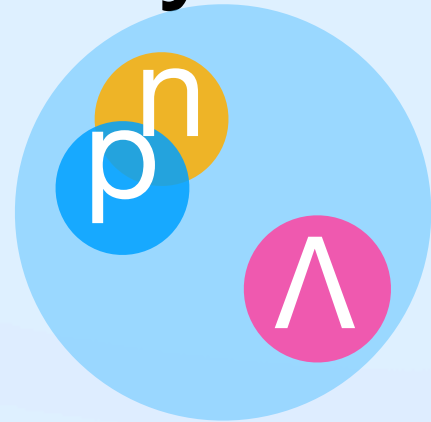
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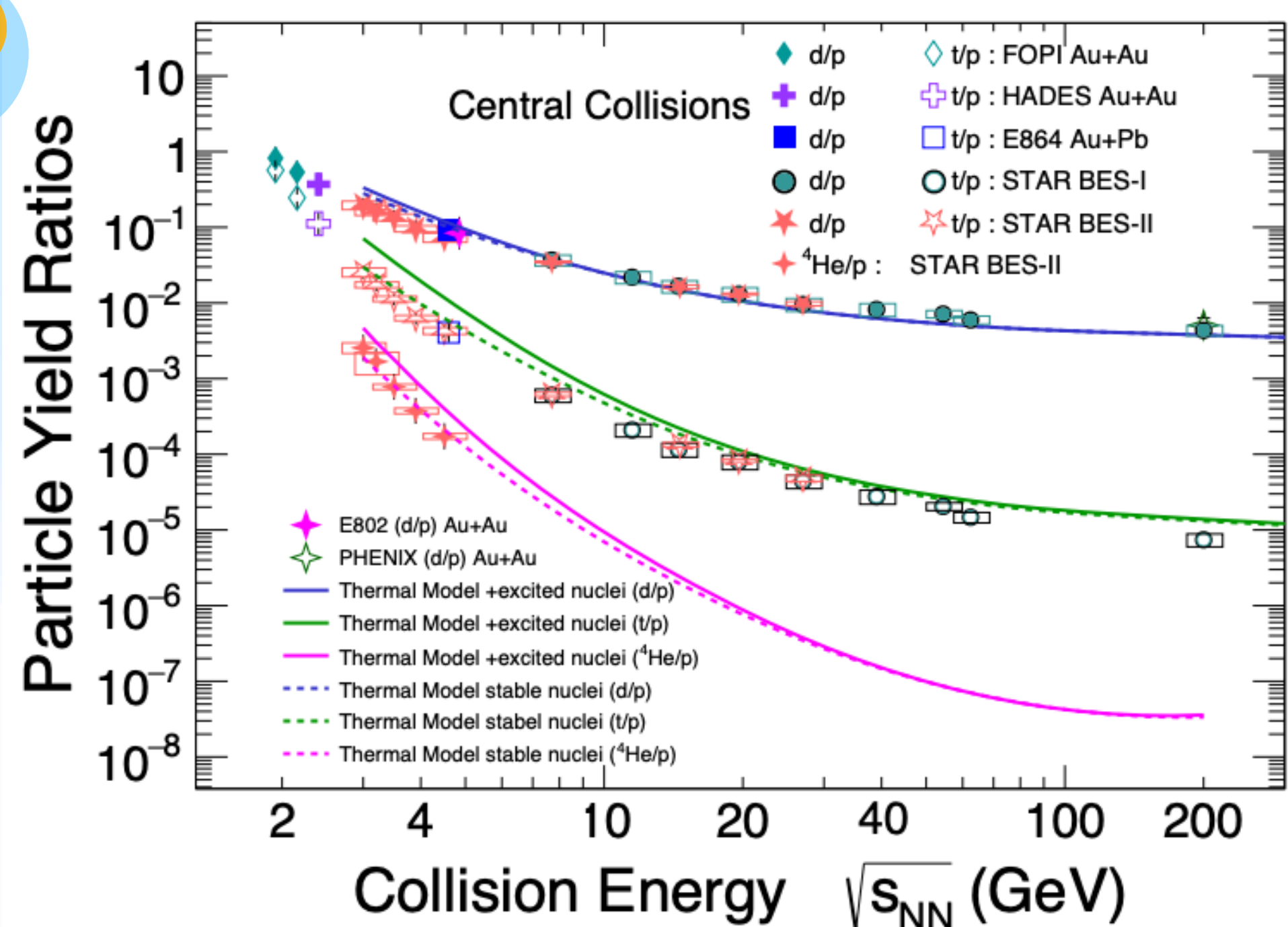
- (Hyper-)nuclei formation after kinetic freeze-out
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*✓* In-medium Y–N interactions modify freeze-out phase space, affecting coalescence and hypernuclei yields



What have we learnt from light nuclei production?

See poster by Liubing Chen ( xx/xx)

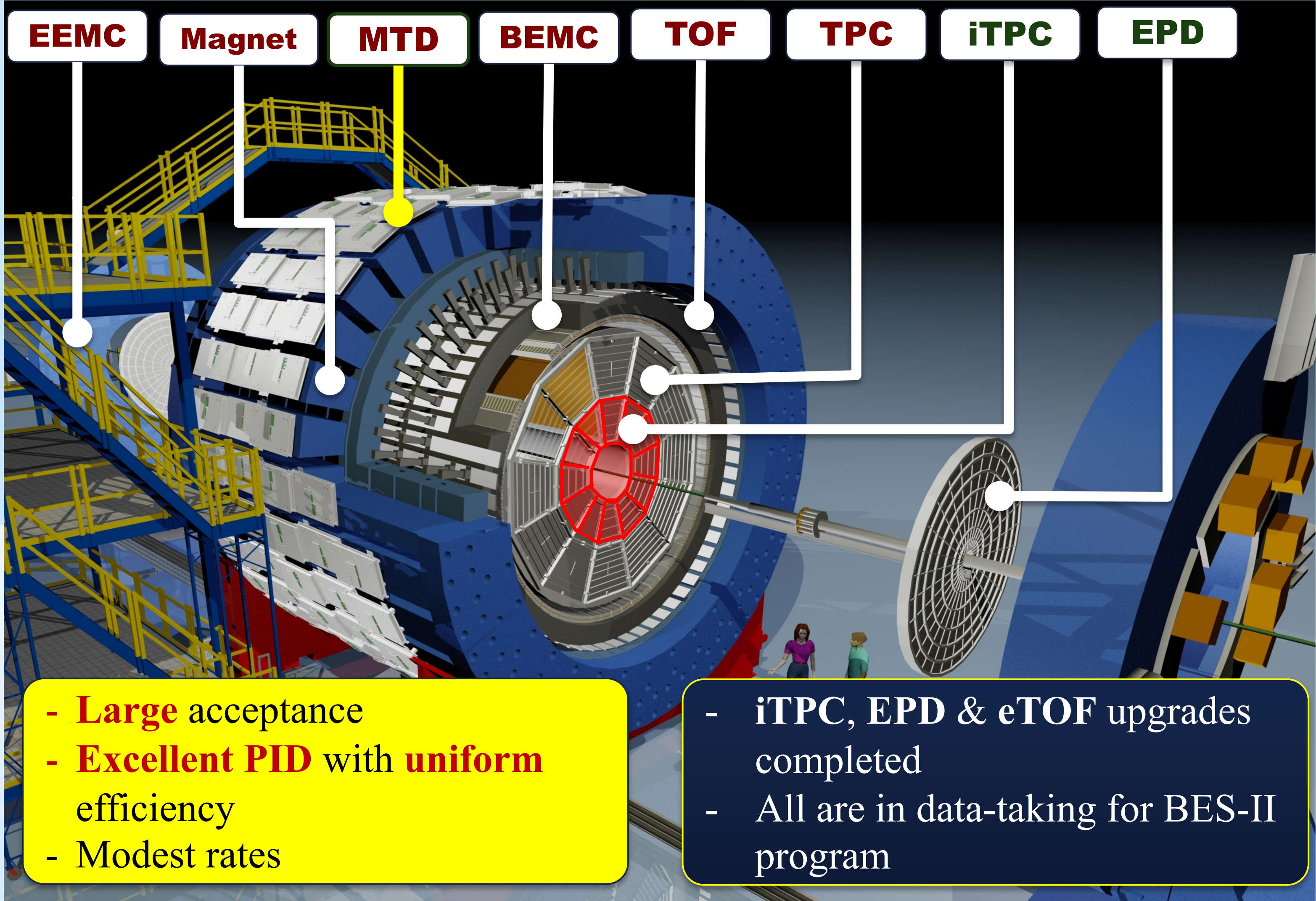


- $d/p$  is fairly well described by thermal model, but  $t/p$ ,  $^4\text{He}/p$  is overestimated

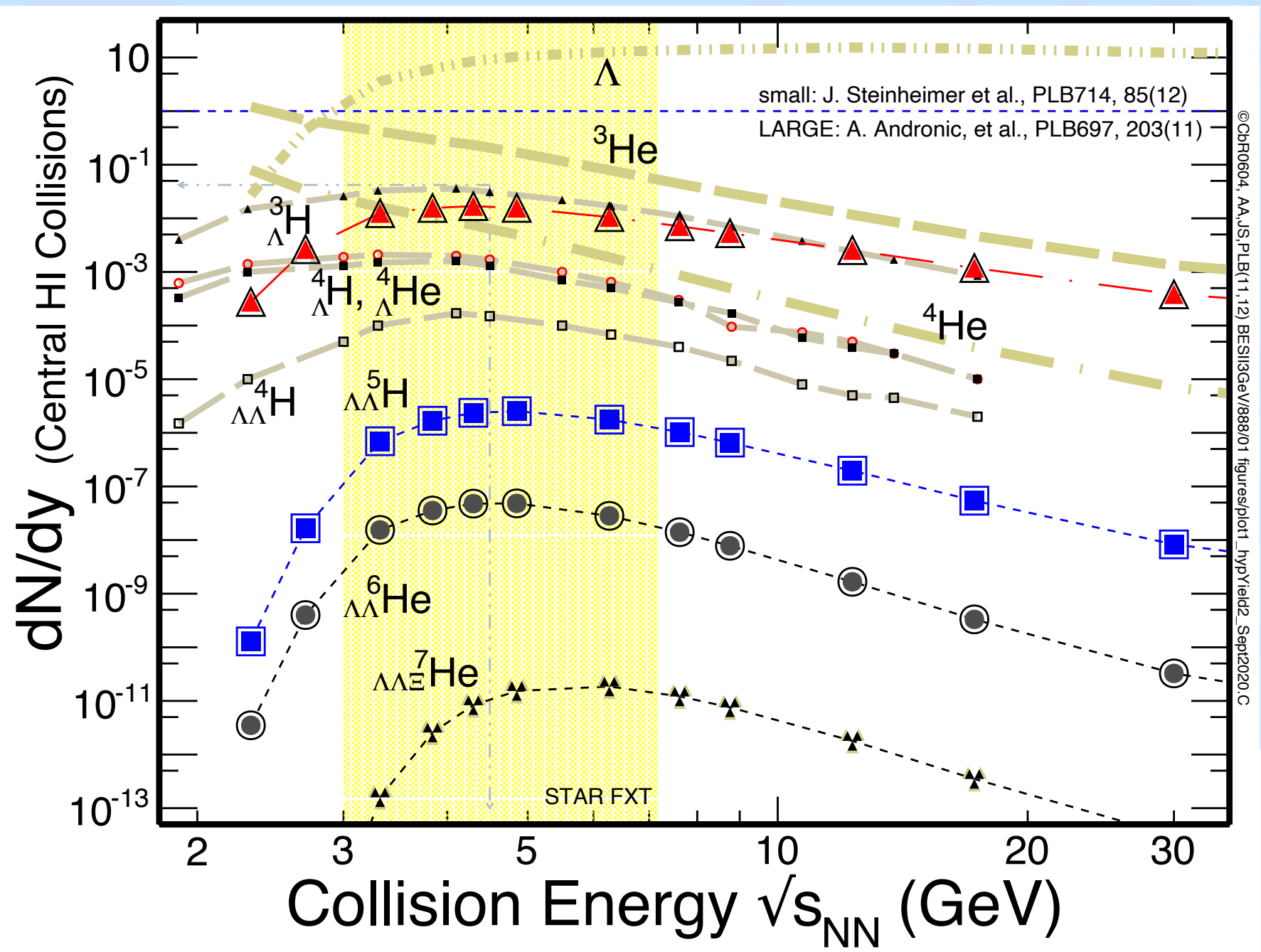
*Recent nuclei measurements poses challenges for thermal model*



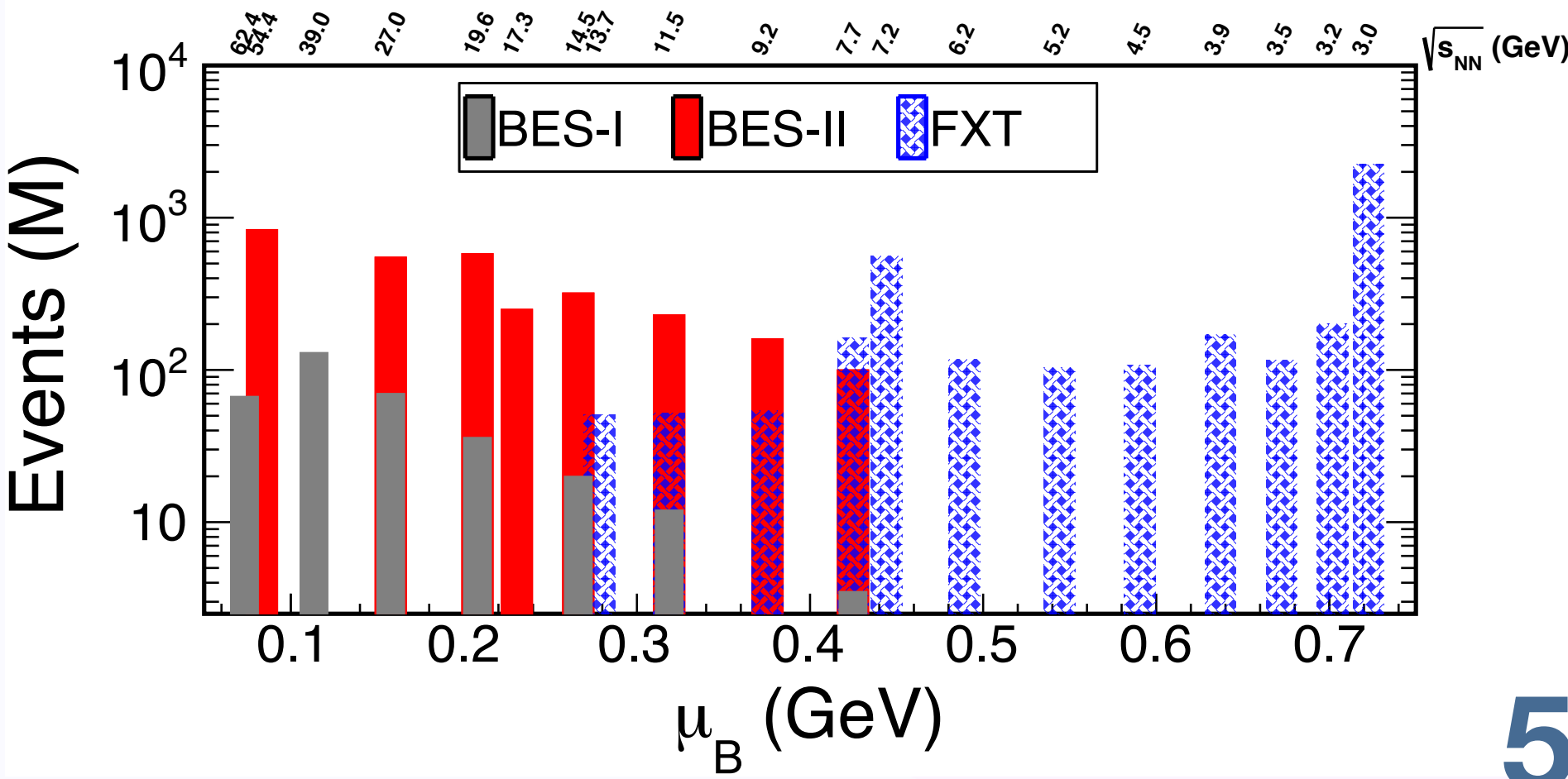
# STAR and Beam Energy Scan



- RHIC BES-II offers great opportunity for hypernuclei measurements

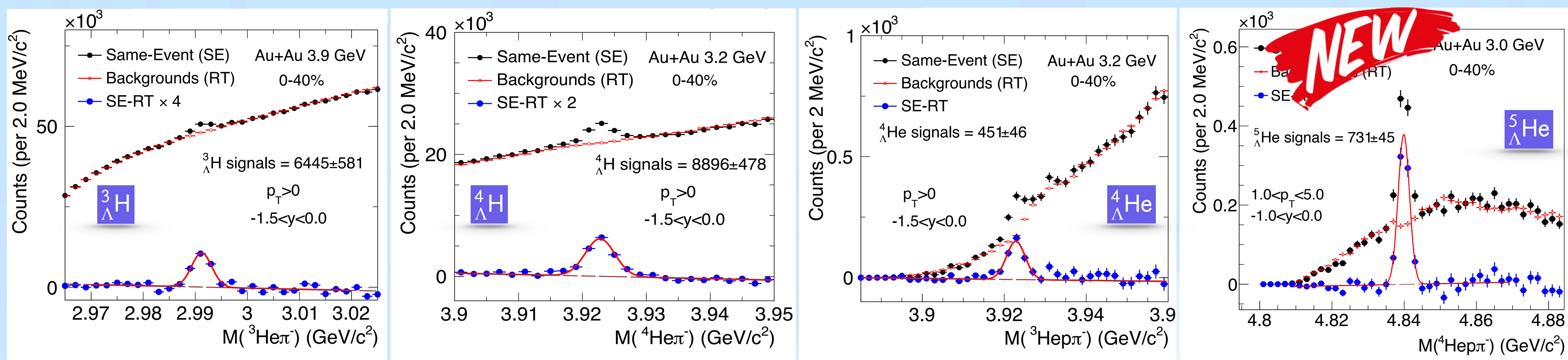


A. Andronic *et al.* Phys.Lett.B 697, 203 (2011)  
J. Steinheimer *et al.* Phys.Lett.B 714, 85 (2012)

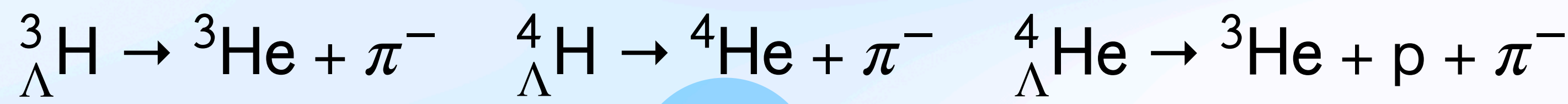




# Hypernuclei Reconstruction



• Hypernuclei are reconstructed using the following decay channels:



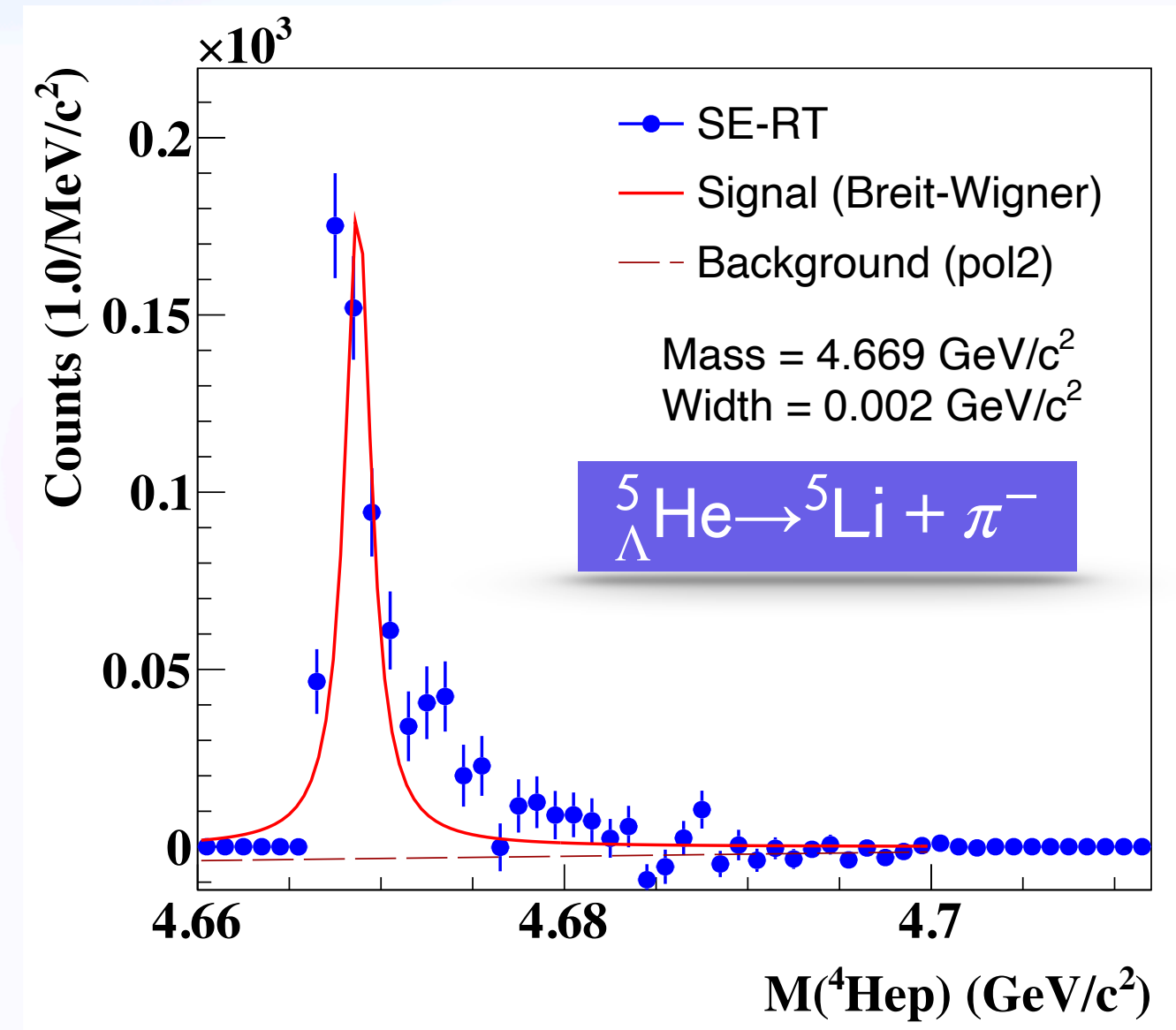
Helps to constrain  $\Lambda$ -N and  $\Lambda$ NN

H. Le et al., PRL 134, 072502 (2025)  
A. Jinno et al., PRC 110, 014001 (2024)

• Combinatorial background estimated via rotating fragments tracks or event mixing

• Efficiency correction using a **data-driven** GEANT simulation

- To account for the decay kinematics of  ${}^4_\Lambda\text{He}$ ,  ${}^5_\Lambda\text{He}$ , the three-body decay phase space is weighted according to the Dalitz plot from data



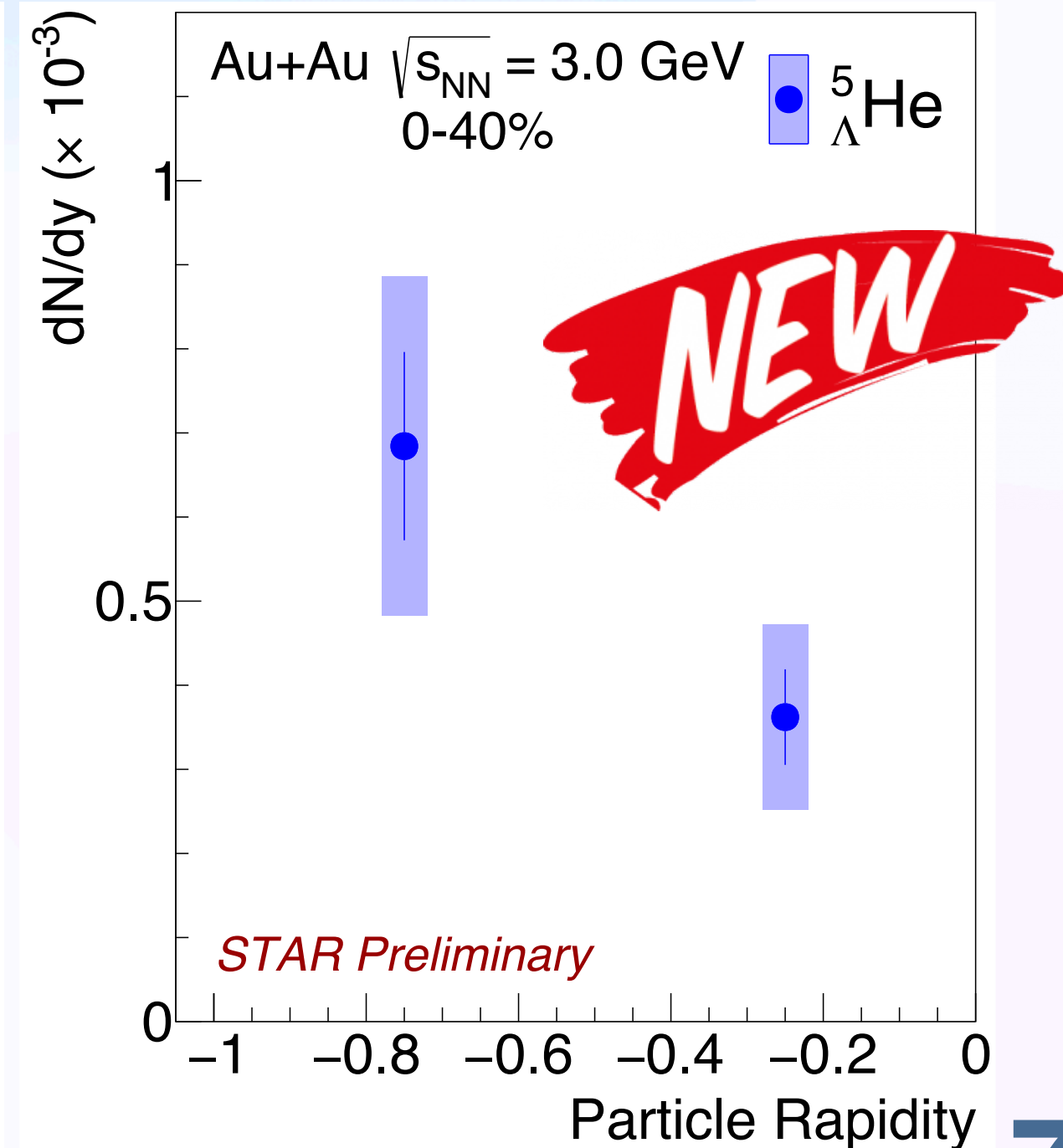
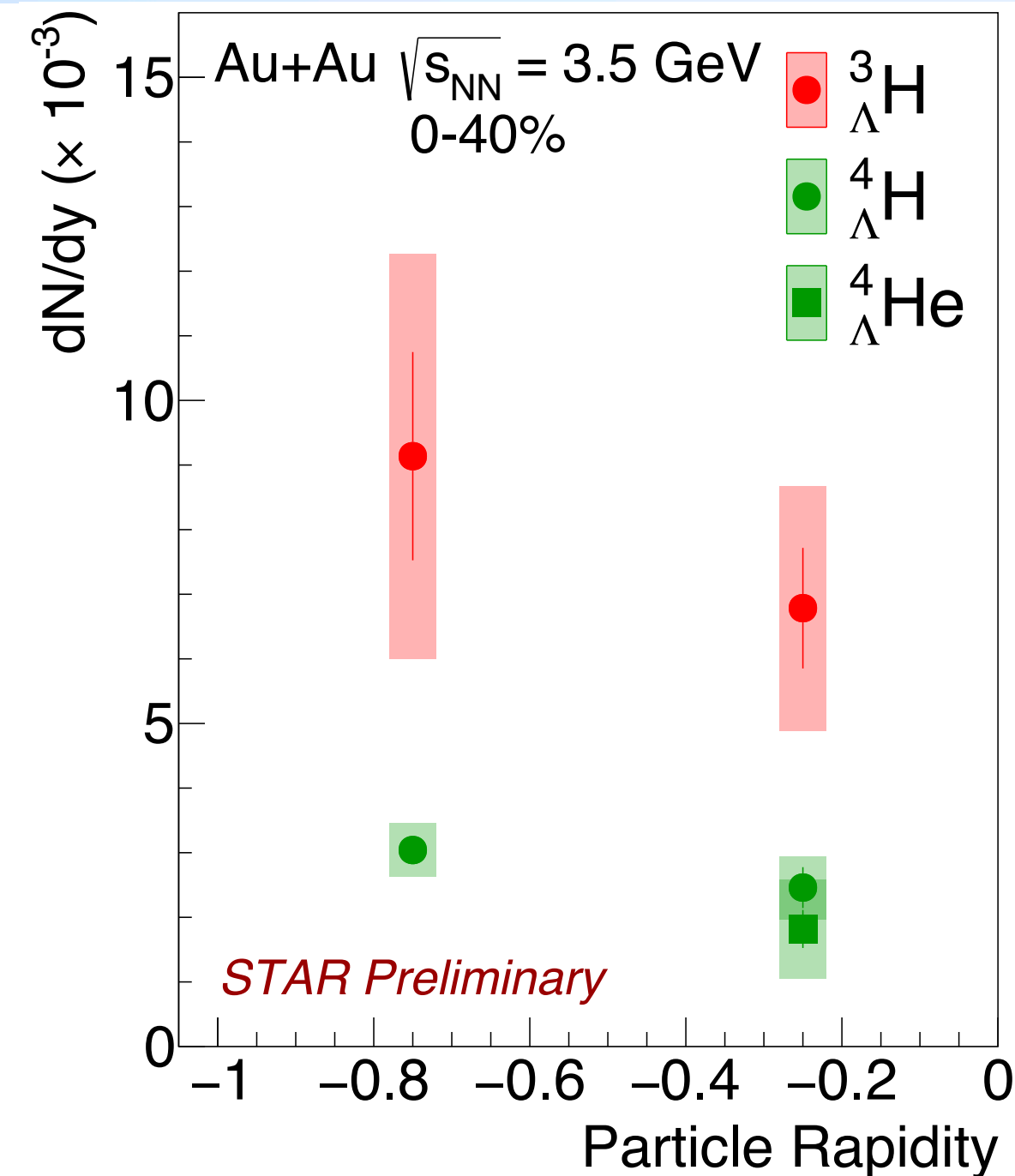
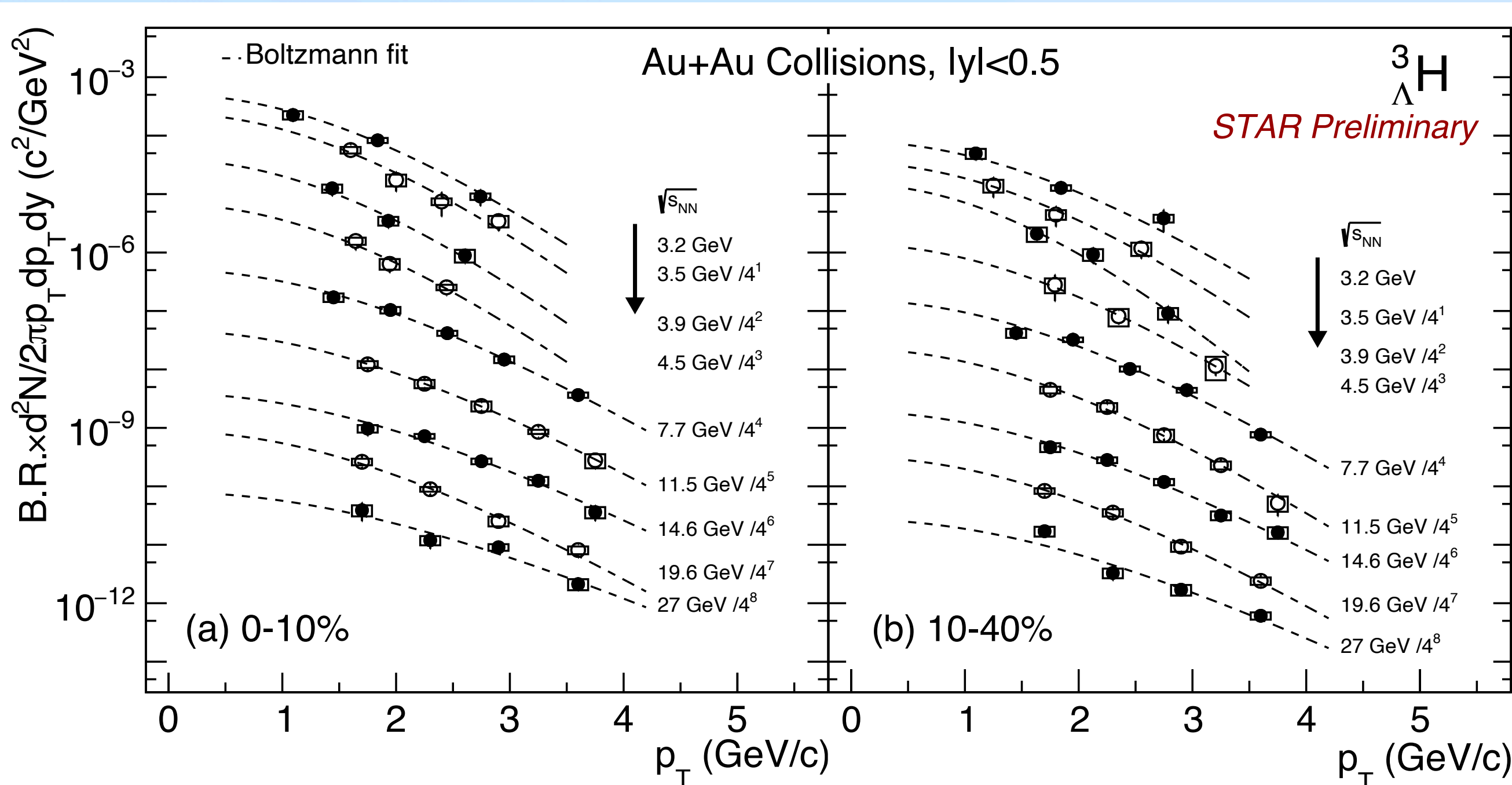


# Hypernuclei $p_T$ Spectra and Rapidity

- Measurements cover different energies
  - $^3_{\Lambda}\text{H}$  in Au+Au collisions at **3-27 GeV**, Au+Au collisions at **200 GeV**
  - $^4_{\Lambda}\text{H}$ ,  $^4_{\Lambda}\text{He}$  in Au+Au collisions at **3-3.5 GeV**
  - $^5_{\Lambda}\text{He}$  in Au+Au collisions at **3 GeV**
- Significant hypernuclei production at target rapidity, more pronounced for heavier hypernuclei  
*Spectator matter matters at target rapidity*

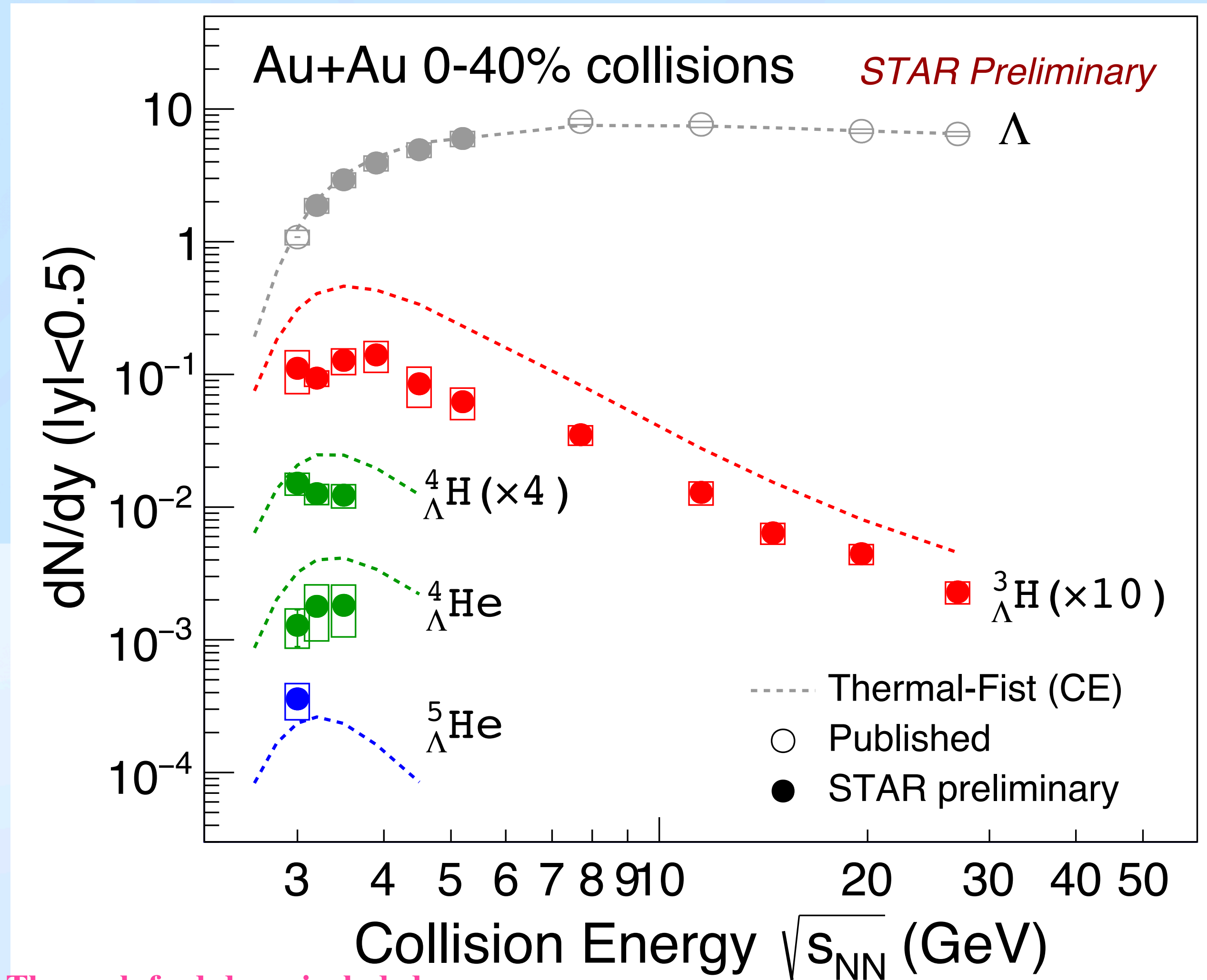
$^3_{\Lambda}\text{H}$  Au+Au  $\sqrt{s_{\text{NN}}} = 3\text{-}27\text{ GeV}$

A. S. Botvina et al., PRC 84, 064904 (2011)





# Excitation Function



Thermal: feed-down included

- ${}^3_{\Lambda}\text{H}$  plateaus at  $\sqrt{s_{NN}} = 3\text{-}4$  GeV
- Similar trend for  ${}^4_{\Lambda}\text{H}$  and  ${}^4_{\Lambda}\text{He}$

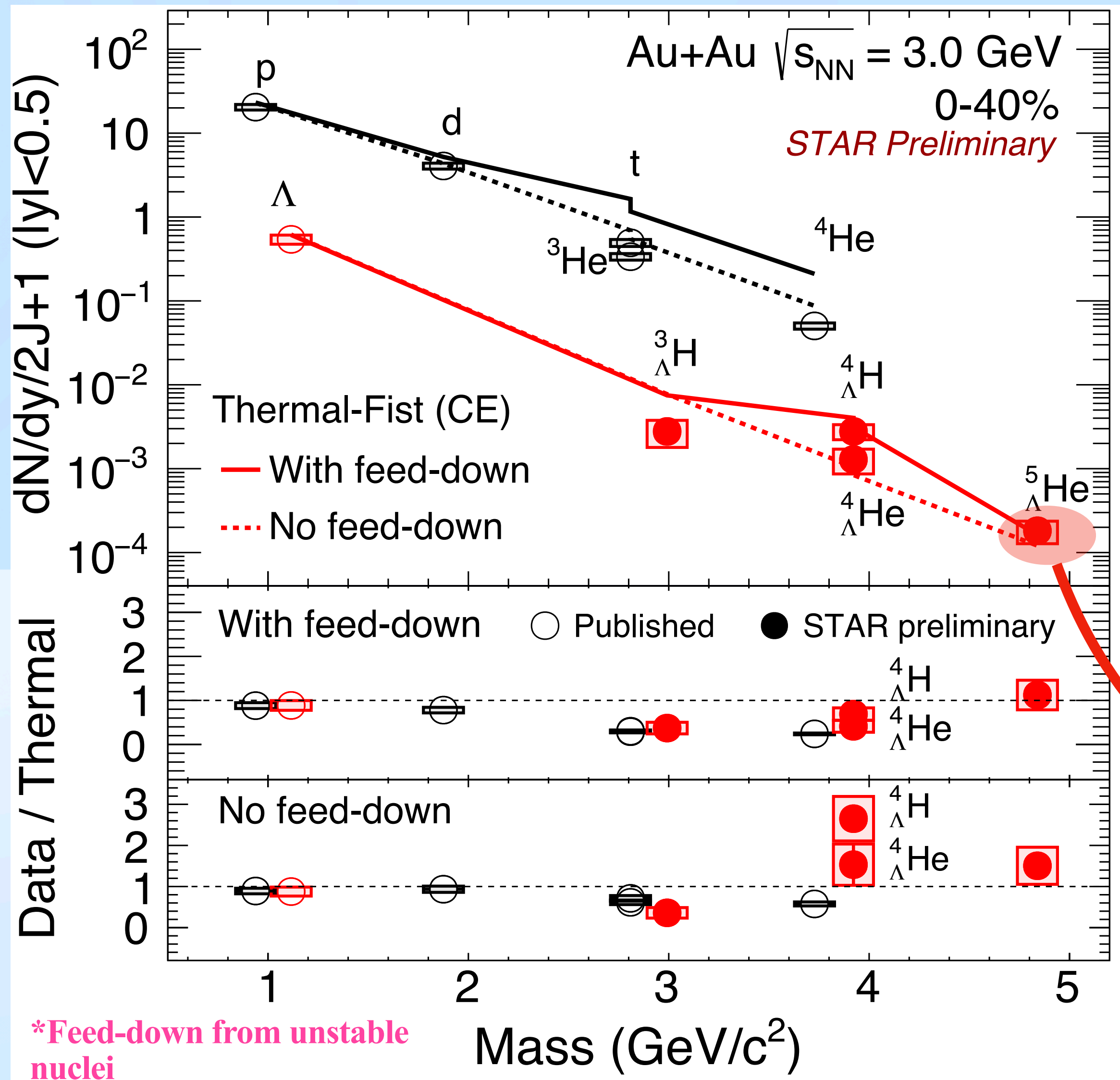
Interplay between increasing baryon production and stronger strangeness canonical suppression towards low energies

*Establishes **low energy** collision experiments as a promising tool to study **exotic strange matter***

- Thermal describes  $\Lambda$ , **over-estimate**  ${}^3_{\Lambda}\text{H}$ ,  ${}^4_{\Lambda}\text{H}$ , and  ${}^4_{\Lambda}\text{He}$ , slightly **under-estimate**  ${}^5_{\Lambda}\text{He}$

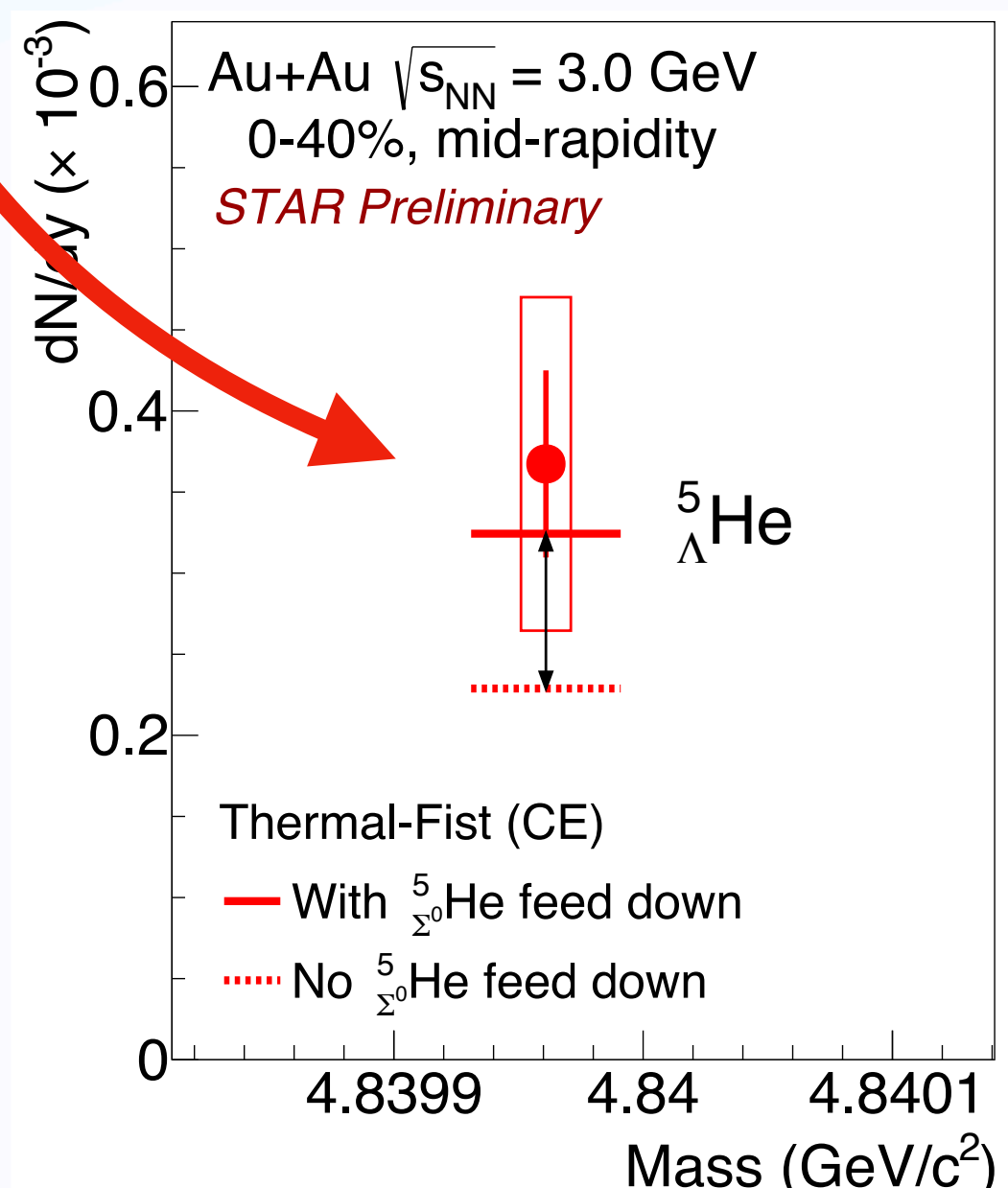
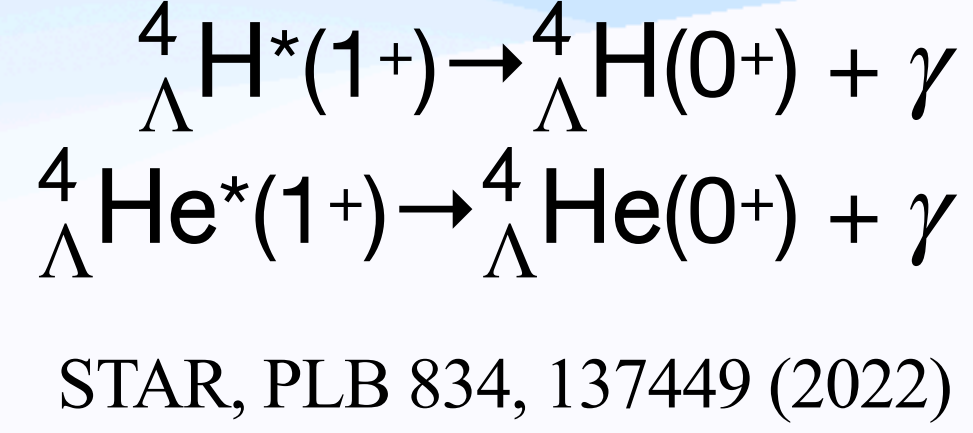


# Comparison to Thermal Model at 3 GeV



STAR, PRC 110, 054911 (2024)  
STAR, PRL 128, 202301 (2022)  
STAR, JHEP **2024** (2024)139

- Thermal model predicts approx. exponential dependence of yields/(2J+1) vs A
- Light nuclei overestimated by thermal with feed-down from unstable nuclei
- Evidence of the formation of  $^4_{\Lambda}H$  and  $^4_{\Lambda}He$  excited states

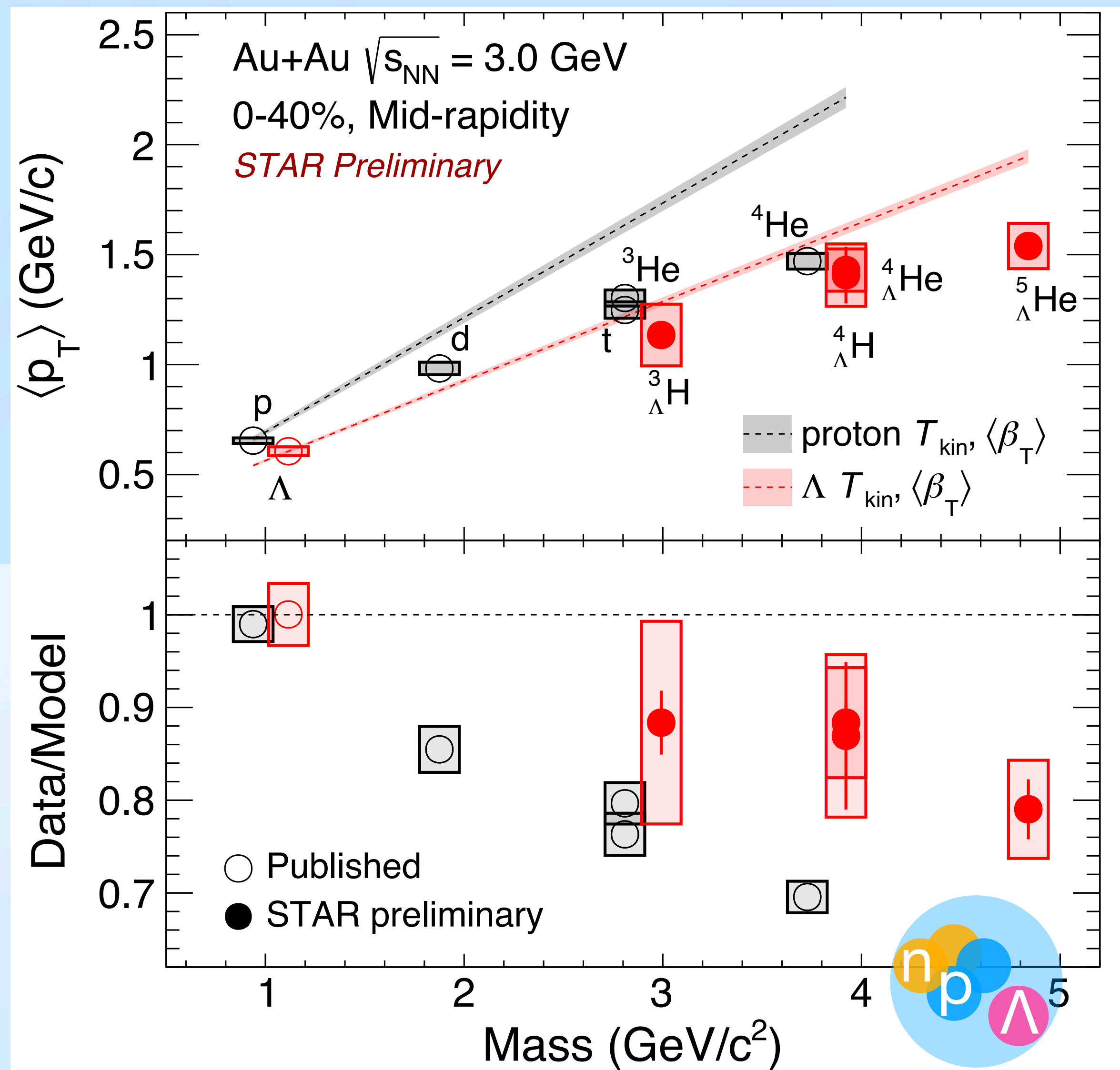


- Possible feed down from  $^5_{\Sigma^0}He \rightarrow ^5_{\Lambda}He + \gamma$

J. Johnstone et al., J. Phys. G 8, L105 (1982)



# Mean Transverse Momentum at 3 GeV



Blast wave function fit

$$\frac{d^2N}{2\pi p_T dp_T dy} = A \int_0^R r dr m_T \times I_0\left(\frac{p_T \sinh \rho(r)}{T_{kin}}\right) K_1\left(\frac{m_T p \cosh \rho(r)}{T_{kin}}\right)$$

$T_{kin}$ : kinetic freeze-out temperature

$\langle \beta_T \rangle$ : average transverse radial flow velocity,  $\rho = \tanh^{-1} \beta_r$

$n$ : the exponent of flow velocity profile,  $n=1$

Hydrodynamic-inspired Blast-Wave model: assumes particles are emitted thermally from an expanding source with a common  $\langle \beta_T \rangle$  and  $T_{kin}$

- Vary the mass to construct the blast-wave prediction using  $p(\Lambda)$  freeze-out parameters

E. Schnedermann et al., PRC 48, 2462 (1993)  
PLB 794, 50–63 (2019)

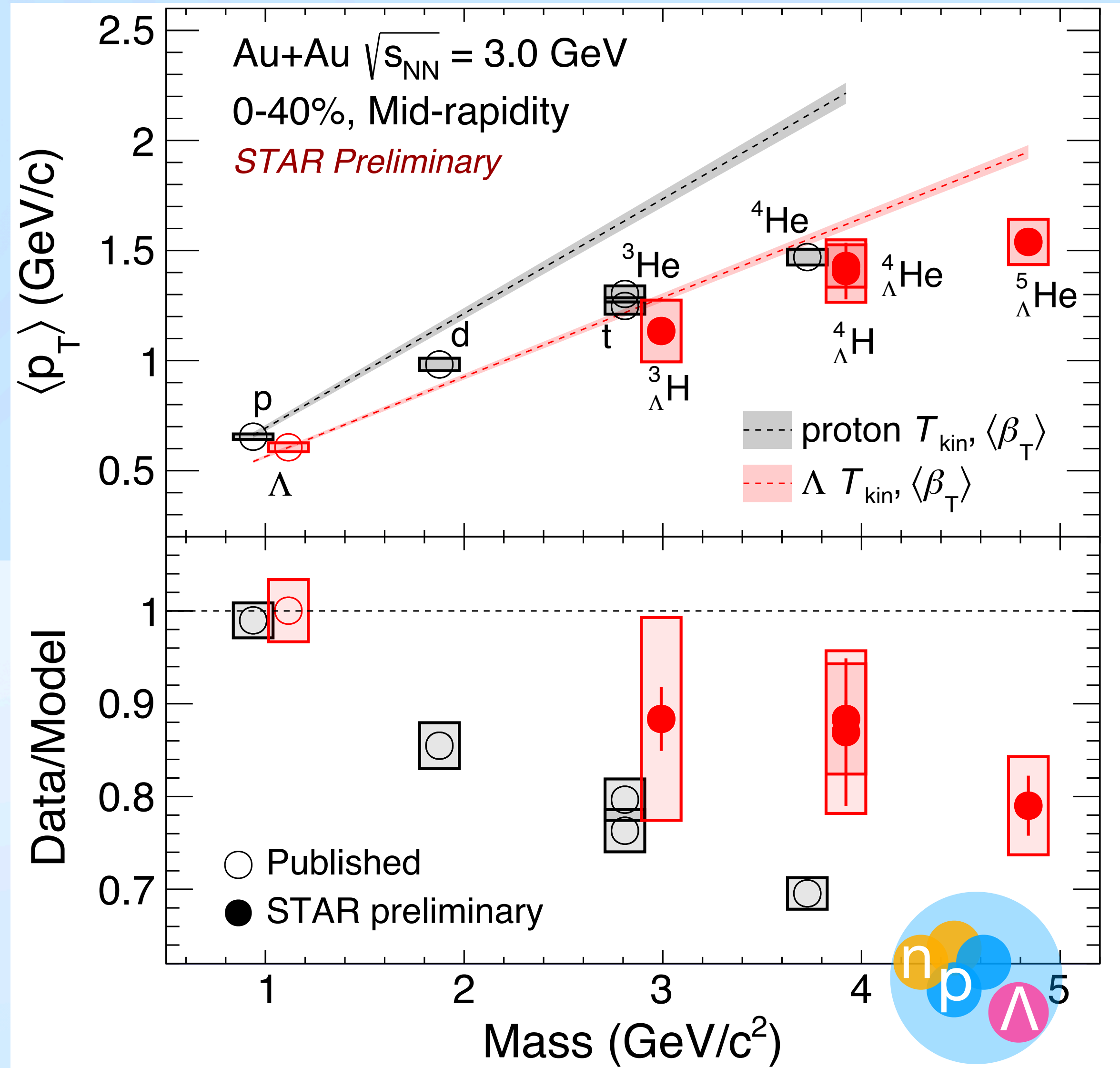
- Light nuclei and hypernuclei deviate from the full hydrodynamic picture

Coalescence scenario?

STAR, PRC 110, 054911 (2024)  
STAR, PRL 128, 202301 (2022)  
STAR, JHEP **2024** (2024)139



# Mean Transverse Momentum at 3 GeV

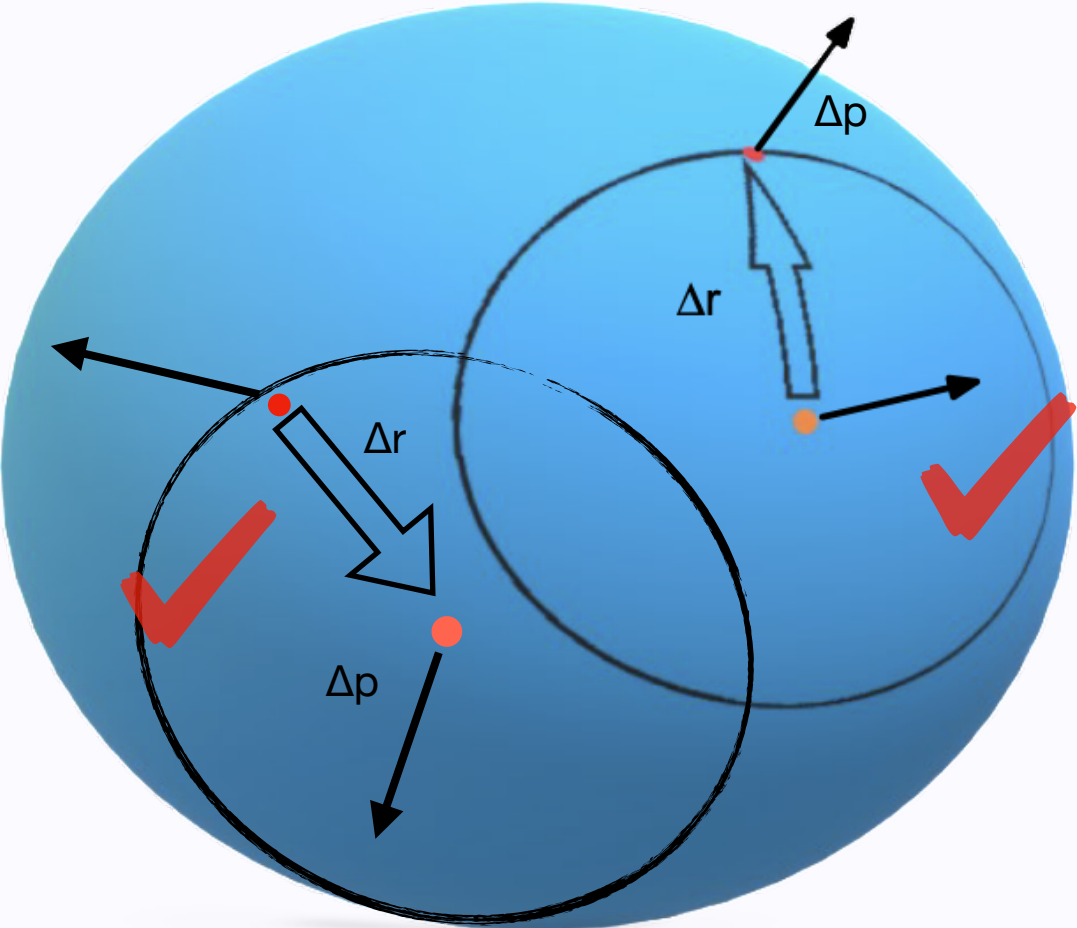


STAR, PRC 110, 054911 (2024)  
 STAR, PRL 128, 202301 (2022)  
 STAR, JHEP **2024** (2024)139

Coalescence scenario: nuclei formed at a later stage after kinetic freeze-out

- Light nuclei and hypernuclei deviate from the full hydrodynamic picture
  1. **Less correlated nucleons** coalesce to nuclei, leading to **smaller  $\langle p_T \rangle$**  than if perfectly aligned
  2. Heavier (hyper)nuclei  $\rightarrow$  Large deviation from blast wave ansatz

A. I. Sheikh et al., PRC 106, 054907 (2022)

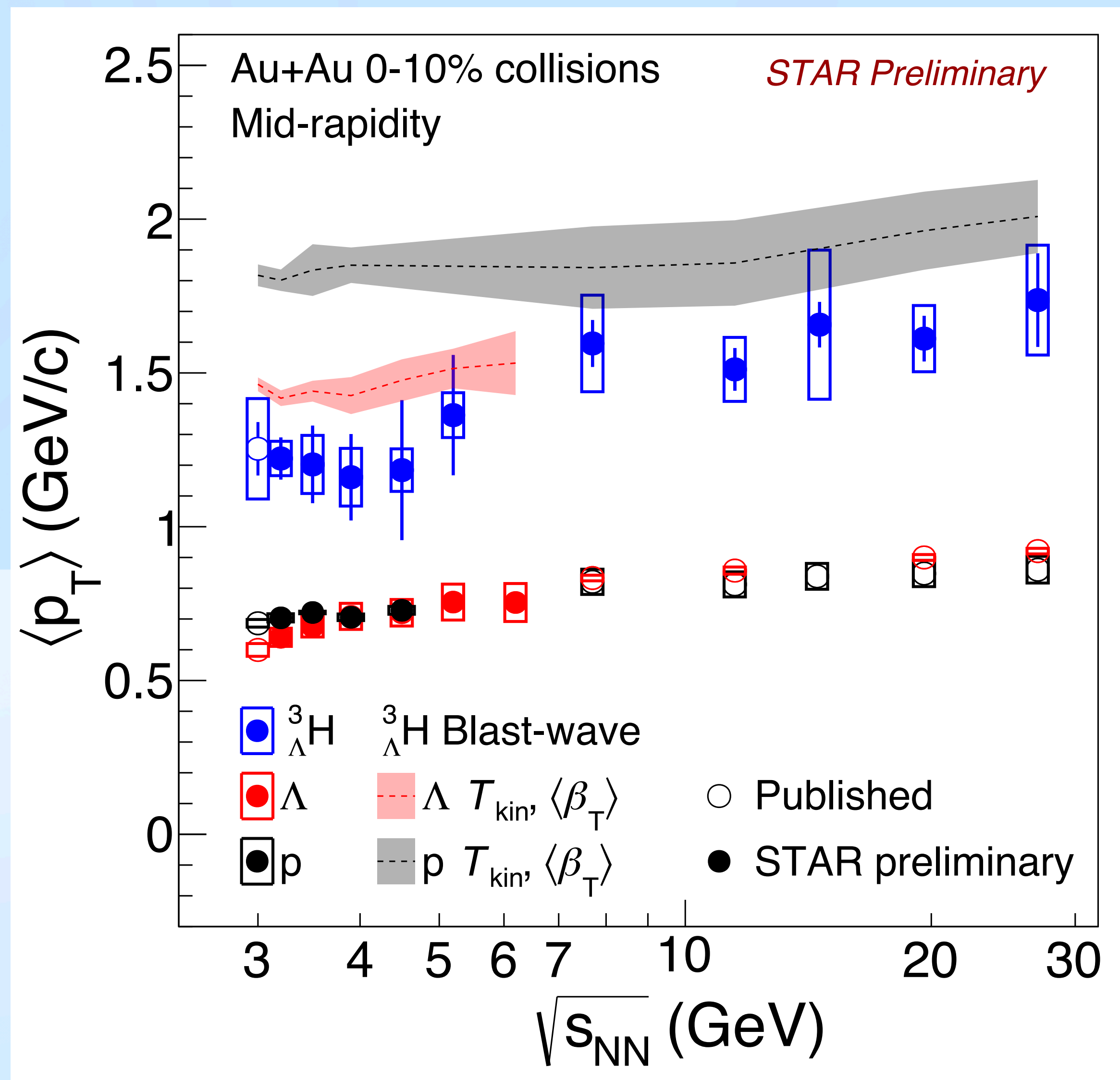


Small effective volume  
 nucleon's  $r$  and  $p$  are less correlated

T. Reichert, SQM 2022, slides (2022)

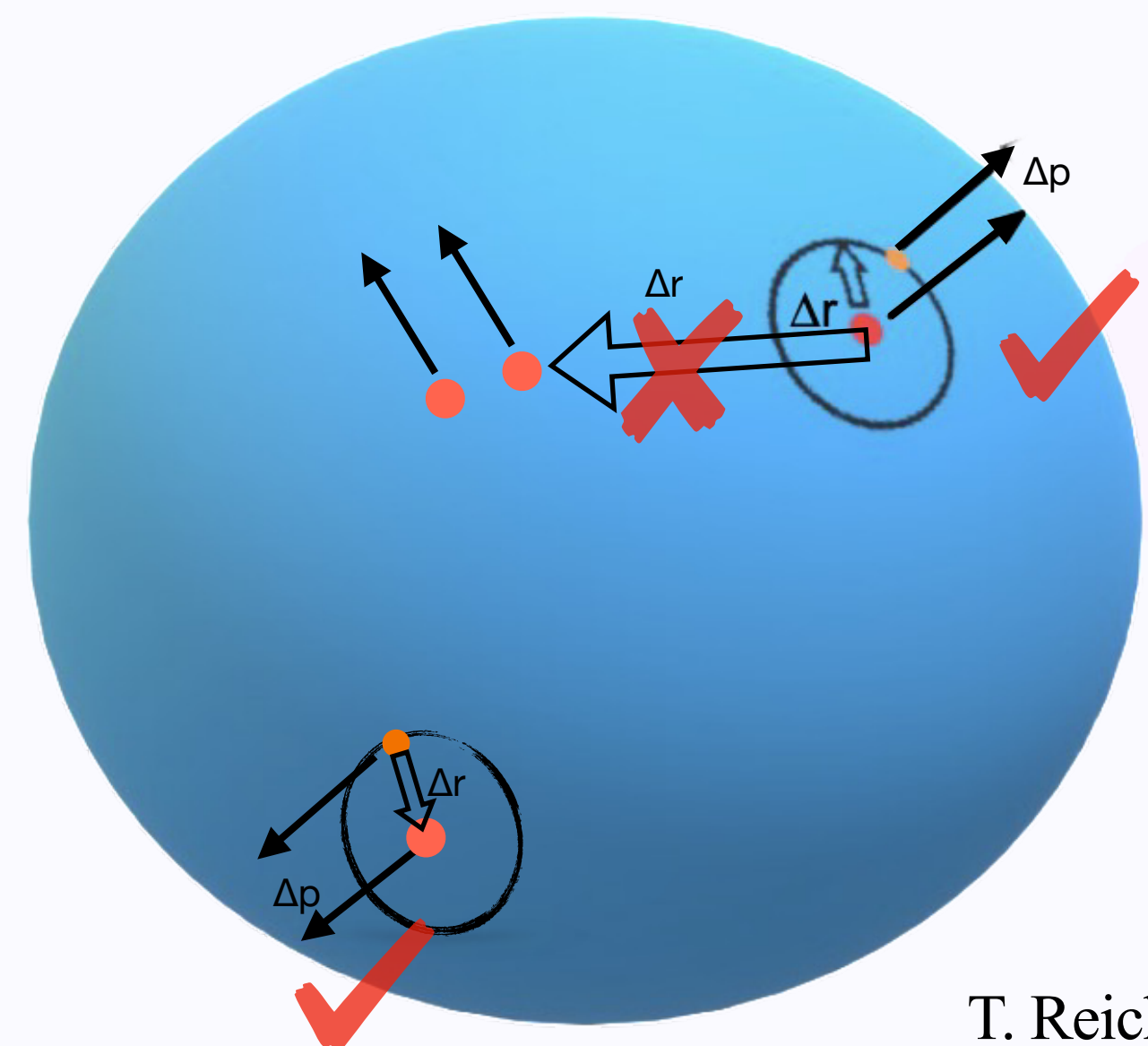


# Mean Transverse Momentum v.s. Collision Energy



- At  $\sqrt{s_{NN}} \geq 7.7$  GeV,  ${}^3_{\Lambda}\text{H}$   $\langle p_T \rangle$  tends to approach the blast-wave prediction with proton freeze-out parameters
- Likely due to larger effective volume, where nucleons that eventually coalesce are more likely to be aligned in space and momentum

*Suggests a different effective volume at  $\sqrt{s_{NN}} \geq 7.7$  GeV, with an noticeable change between 4.5 and 7.7 GeV*



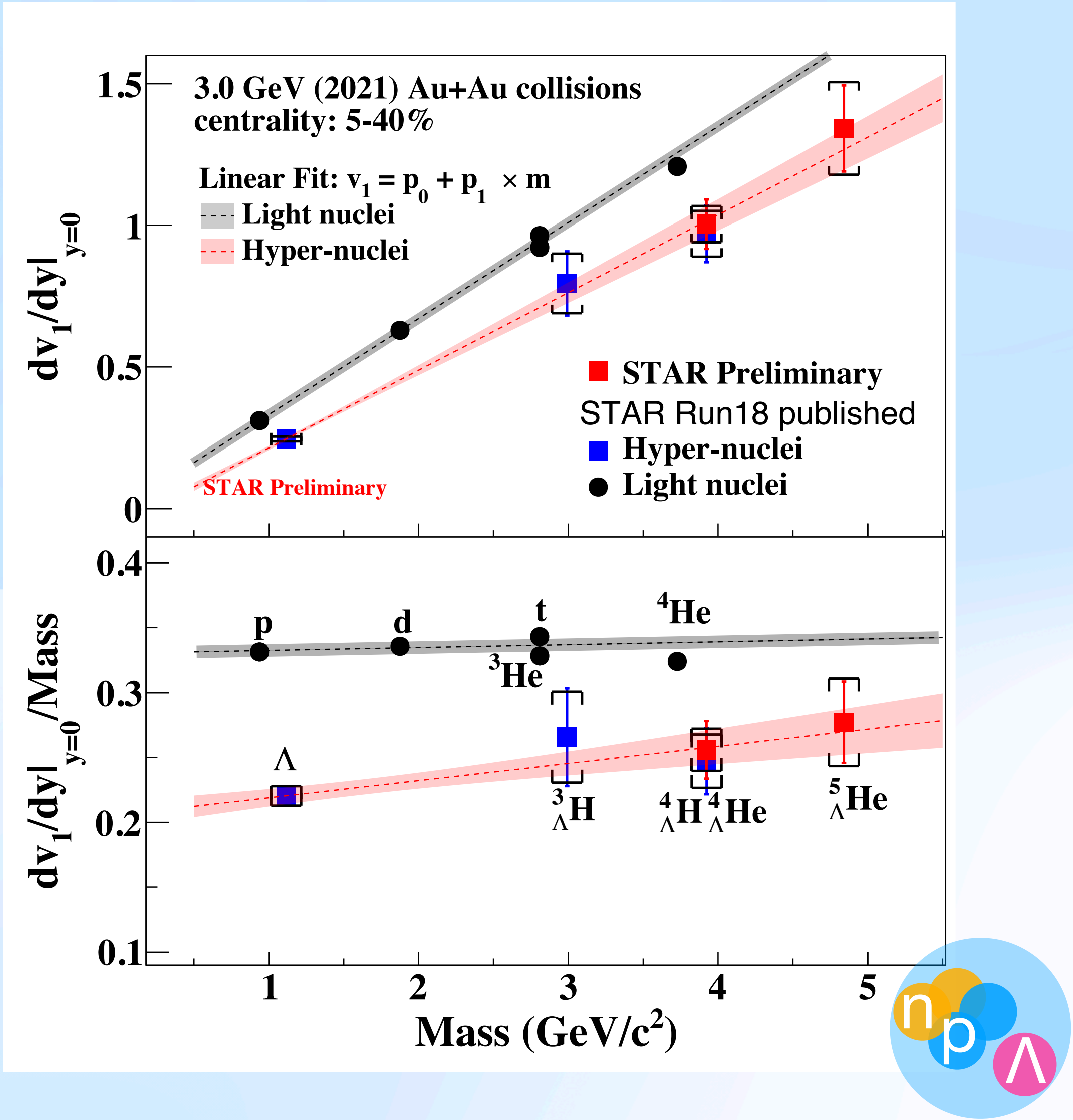
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STAR, PRC 110, 054911 (2024)  
STAR, PRL 128, 202301 (2022)  
STAR, JHEP **2024** (2024)139

T. Reichert, SQM 2022, slides (2022)

# Collective Flow at 3 GeV

See poster by Junyi Han ( xx/xx)



- Directed flow of hypernuclei follows mass scaling
- Qualitatively consistent with coalescence formation of hypernuclei



# Summary and Outlook

- Hypernuclei measurement from STAR BES-II at  $\sqrt{s_{NN}} = 3-27$  GeV
  1. **First measurement of  $A = 5$  hypernuclei yield and directed flow in Au+Au collisions at  $\sqrt{s_{NN}} = 3$  GeV**
  2. Thermal model overestimates  ${}^3_{\Lambda}\text{H}$ ,  ${}^4_{\Lambda}\text{H}$ , and  ${}^4_{\Lambda}\text{He}$ , and slightly underestimates  ${}^5_{\Lambda}\text{He}$
  3. Mean transverse momentum tends to be lower than hydrodynamic-inspired blast-wave model at  $\sqrt{s_{NN}} < 7.7$  GeV

Consistent with coalescence picture: weaker space and momentum correlation among coalescing nucleons in a smaller effective volume at low collision energies
  4. Collective flow qualitatively consistent with coalescence model

## Outlook:

- High statistics 3 GeV FXT data: more precise measurement of  $A < 5$  hypernuclei,  ${}^5_{\Lambda}\text{He}$  intrinsic properties ( $B_{\Lambda}$ , dalitz plot, lifetime), search for heavier hypernuclei ( $A > 5$ ), double- $\Lambda$  hypernuclei
- ➔ Further constrain production mechanism, structure of hypernuclei, and YN, YY, YNN interactions

**Backup Slides Follow**

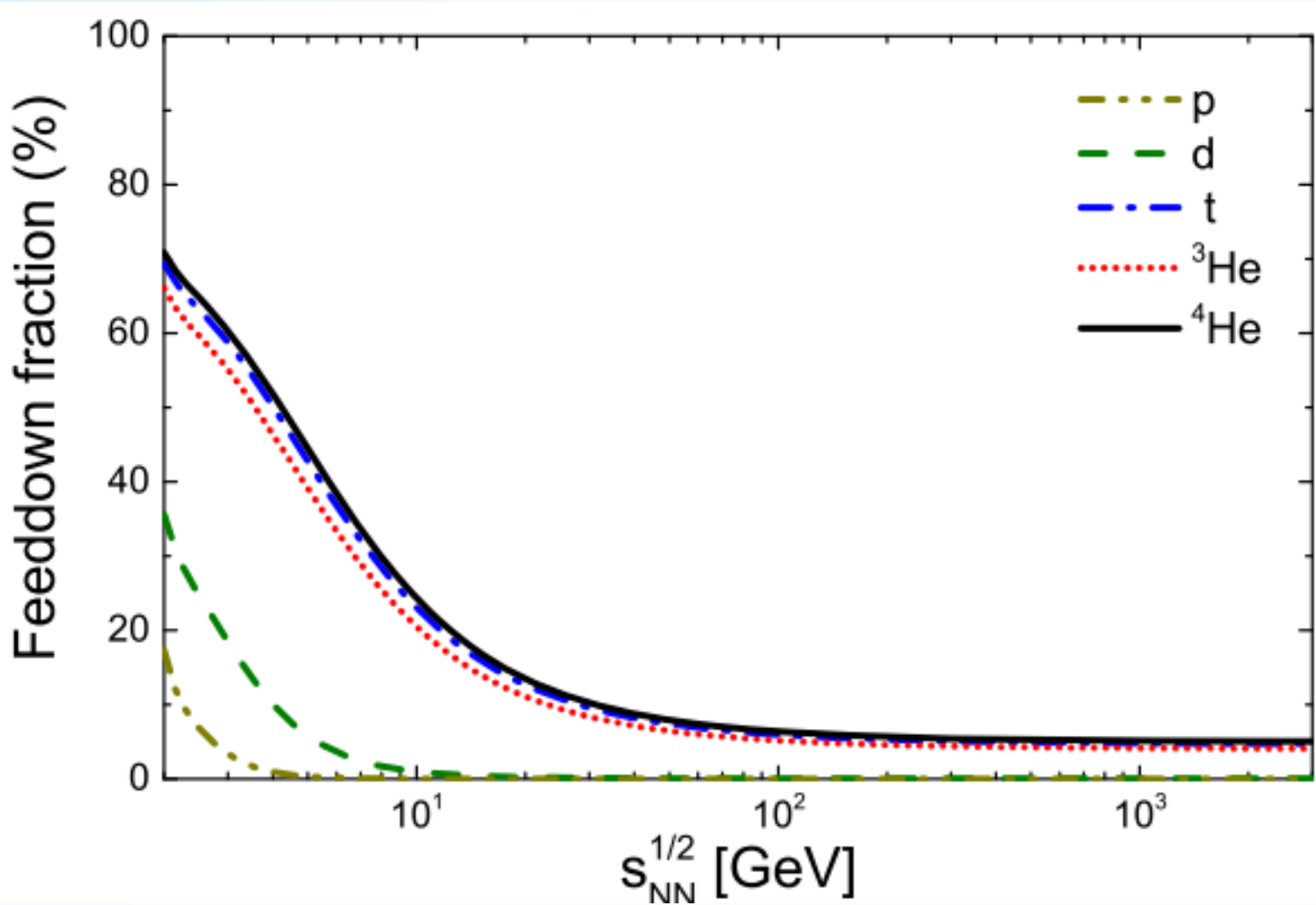


# Feed-down from Unstable Nuclei

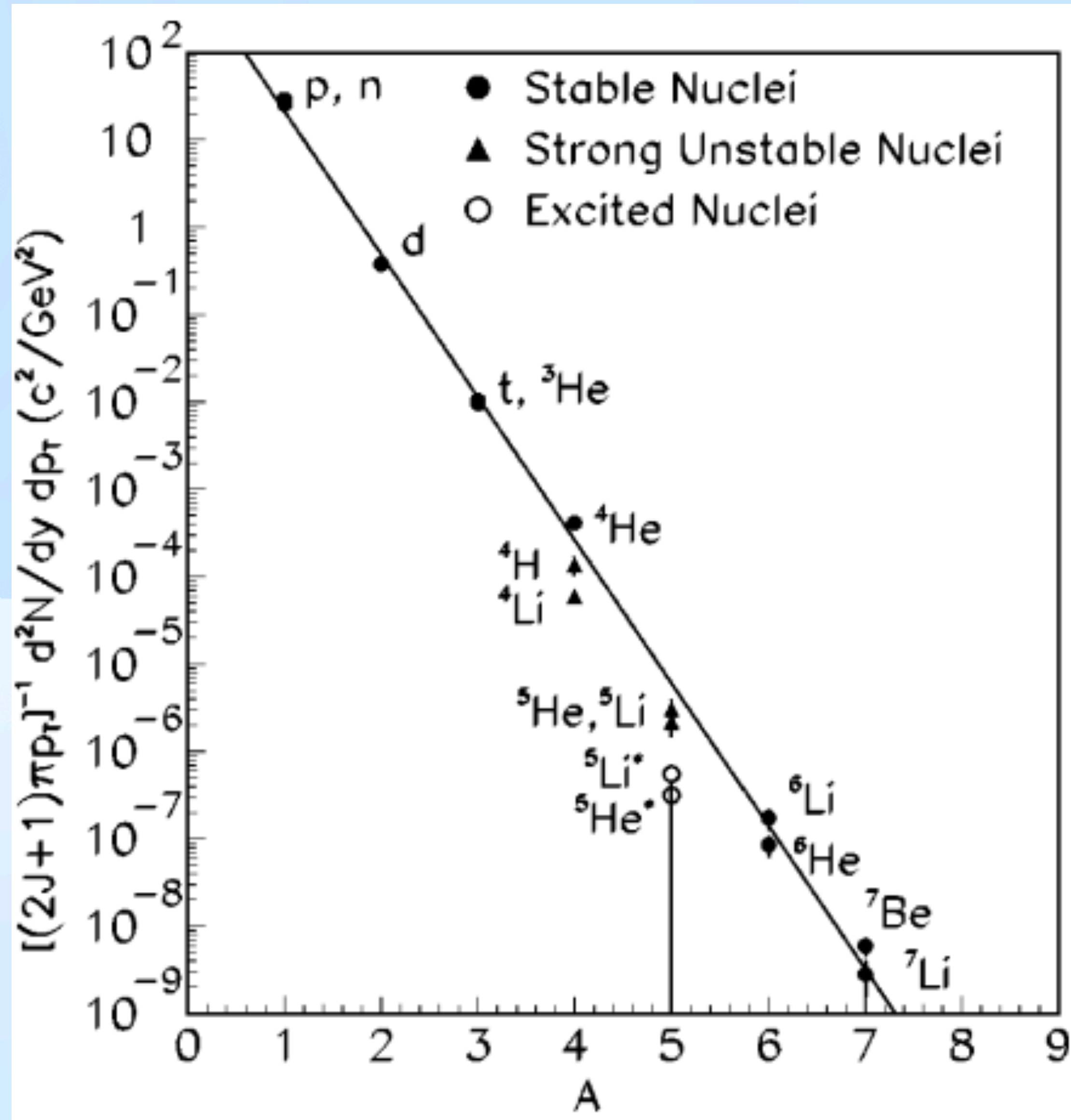
| $A = 4$         | $E_x$ (MeV) | $J^\pi$ | Decay channels               |
|-----------------|-------------|---------|------------------------------|
| ${}^4\text{H}$  | g.s.        | $2^-$   | n(100%)                      |
|                 | 0.31        | $1^-$   | n(100%)                      |
|                 | 2.08        | $0^-$   | n(100%)                      |
|                 | 2.83        | $1^-$   | n(100%)                      |
| ${}^4\text{He}$ | g.s.        | $0^+$   | stable                       |
|                 | 20.21       | $0^+$   | p(100%)                      |
|                 | 21.01       | $0^-$   | n(23.8%), p(76.2%)           |
|                 | 21.84       | $2^-$   | n(37.3%), p(62.7%)           |
|                 | 23.33       | $2^-$   | n(47.3%), p(52.7%)           |
|                 | 23.64       | $1^-$   | n(44.5%), p(55.5%)           |
|                 | 24.25       | $1^-$   | n(47.0%), p(50.5%), d(2.5%)  |
|                 | 25.28       | $0^-$   | n(48.3%), p(51.7%)           |
|                 | 25.95       | $1^-$   | n(48.5%), p(51.5%)           |
|                 | 27.42       | $2^+$   | n(3%), p(3%), d(94%)         |
|                 | 28.31       | $1^+$   | n(47%), p(48%), d(5%)        |
|                 | 28.37       | $1^-$   | n(2%), p(2%), d(96%)         |
|                 | 28.39       | $2^-$   | n(0.25%), p(0.25%), d(99.5%) |
|                 | 28.64       | $0^-$   | d(100%)                      |
|                 | 28.67       | $2^+$   | d(100%)                      |
|                 | 29.89       | $2^+$   | n(0.4%), p(0.4%), d(99.2%)   |
| ${}^4\text{Li}$ | g.s.        | $2^-$   | p(100%)                      |
|                 | 0.32        | $1^-$   | p(100%)                      |
|                 | 2.08        | $0^-$   | p(100%)                      |
|                 | 2.85        | $1^-$   | p(100%)                      |

branching ratios in the right column of Table I. For the  $A = 5$  states the channels  ${}^5\text{H} \rightarrow \text{t} + \text{n} + \text{n}$ ,  ${}^5\text{He} \rightarrow {}^4\text{H} + \text{p}$ ,  ${}^5\text{He} \rightarrow \text{t} + \text{d}$ ,  ${}^5\text{Li} \rightarrow {}^4\text{He} + \text{p}$ , and  ${}^5\text{Li} \rightarrow {}^4\text{Li} + \text{n}$  are taken into account. The excited nuclei feeding will thus affect the yields of nucleons and stable light nuclei, but not other particles.

| $A = 5$         | $E_x$ (MeV) | $J^\pi$         | Decay channels |
|-----------------|-------------|-----------------|----------------|
| ${}^5\text{H}$  | g.s.        | $\frac{1}{2}^+$ | 2n(100%)       |
| ${}^5\text{He}$ | g.s.        | $\frac{3}{2}^-$ | n(100%)        |
|                 | 1.27        | $\frac{1}{2}^-$ | n(100%)        |
|                 | 16.84       | $\frac{3}{2}^+$ | n(60%), d(40%) |
| ${}^5\text{Li}$ | g.s.        | $\frac{3}{2}^-$ | p(100%)        |
|                 | 1.49        | $\frac{1}{2}^-$ | p(100%)        |
|                 | 16.87       | $\frac{3}{2}^+$ | p(70%), n(30%) |



# Feed-down from Unstable Nuclei

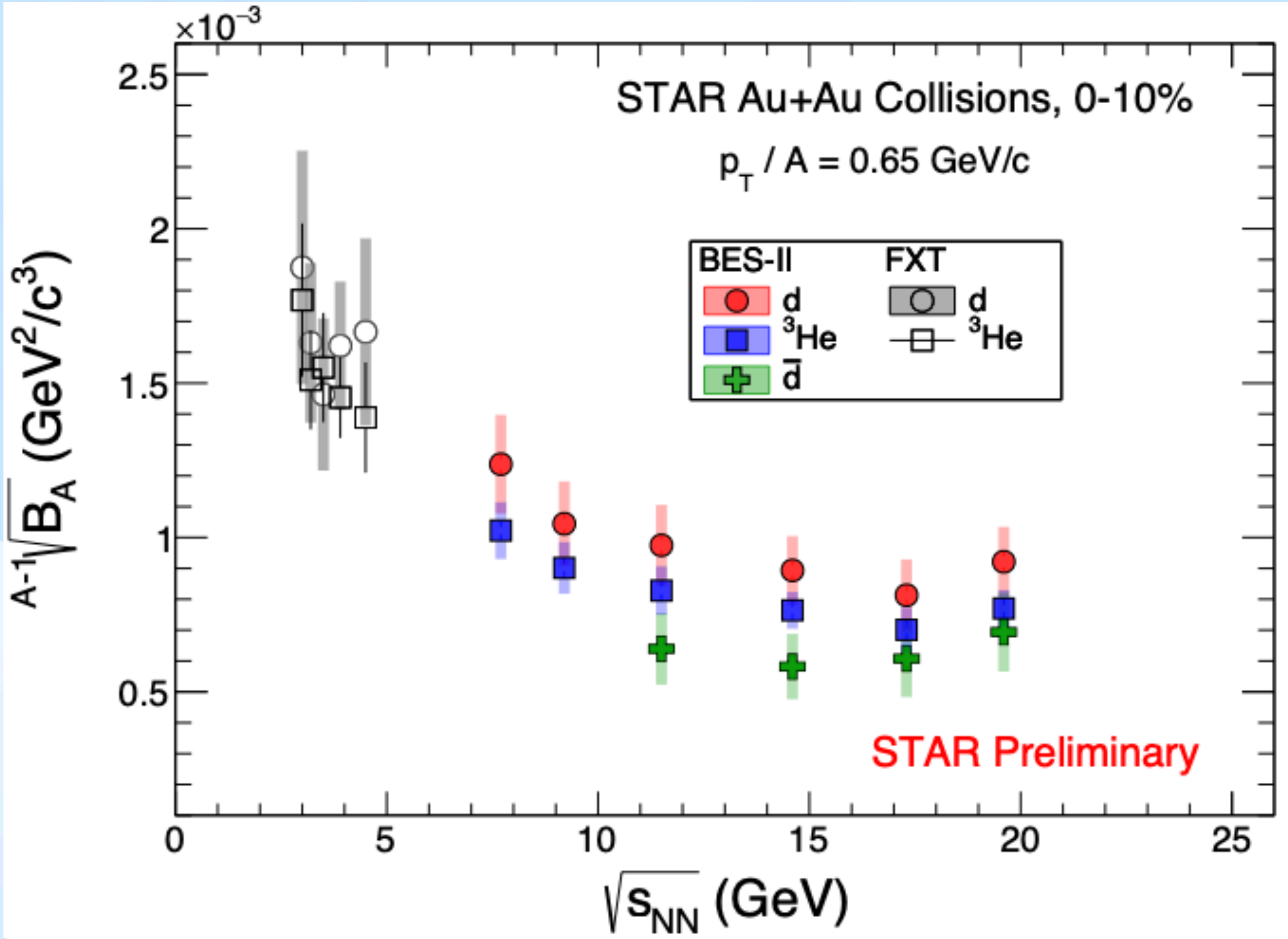


- Suppression of A=4 unstable states compared to <sup>4</sup>He ground state observed at E864
- Indicates unstable nuclei yield likely overestimated by thermal model



# Coalescence Parameters

See poster by Yixuan Jin ( xx/xx)

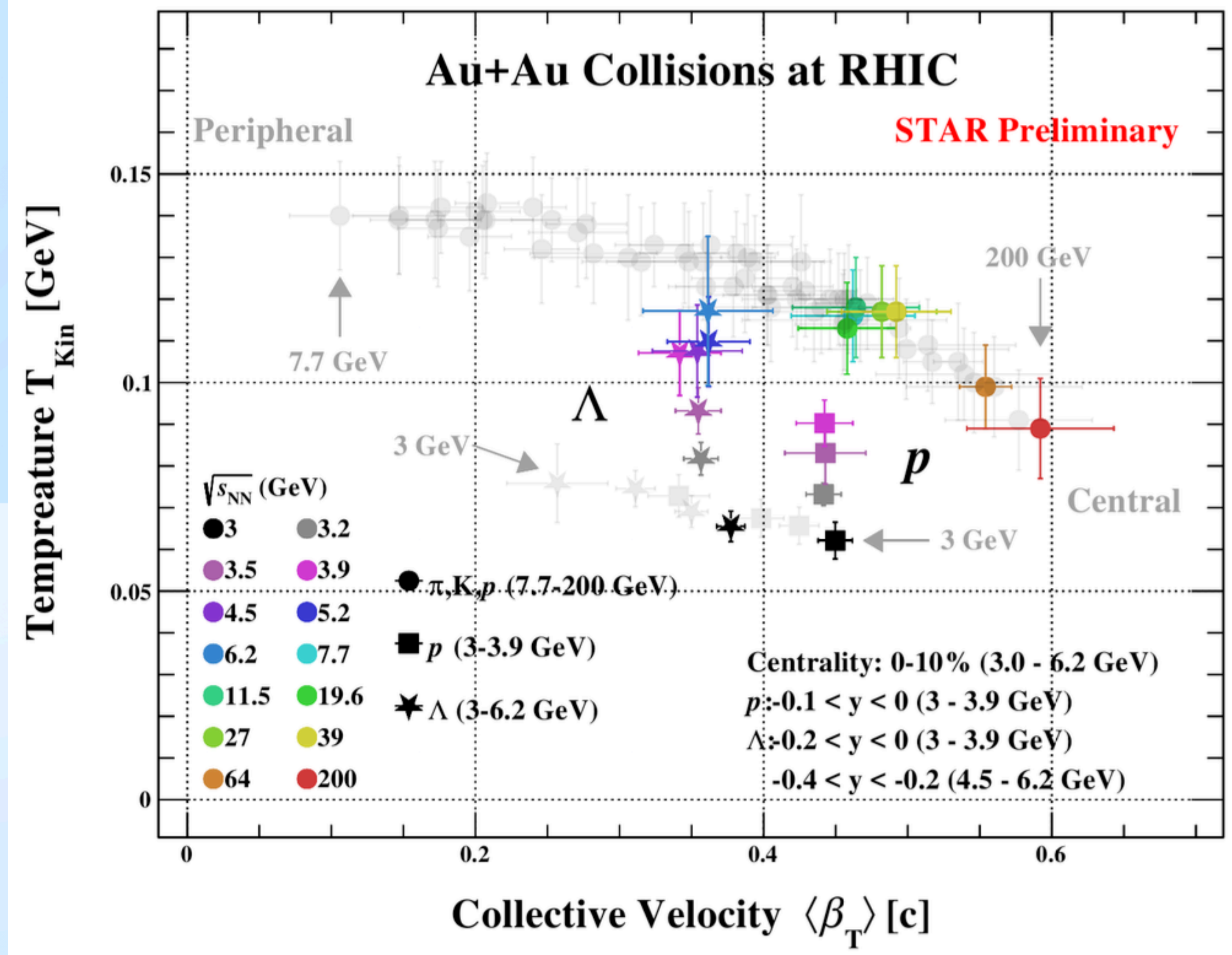


- $B_A$  decrease with increasing energy due to the larger effective volume

STAR, PRC 99, 064905 (2019)  
K. J. Sun et al., PLB 774, 103 (2017); PLB 781, 499 (2018)

# Kinetic Freeze-out Parameters

See talk by Hongcan Li ( xx/xx)



- Energy dependence at the most central collisions
- Different freeze-out parameter due to different production mechanism between proton and  $\Lambda$
- Hadronic transport model UrQMD model qualitatively reproduces the trend at 3 - 6.2 GeV



# Dalitz Plots

$${}^5_{\Lambda}\text{He} \rightarrow {}^4\text{He} + \text{p} + \pi^{-}$$

I. Kisel, EPJ Web Conf. 271, 08001 (2022)

